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Investigating the structural, electronic, and optical properties of the novel double perovskite K_2AgBiI_6 using DFT

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In this groundbreaking study, we unveil the remarkable structural, electronic, and optical Properties of the newly discovered double perovskite material, K_2AgBiI_6 , presenting a paradigm shift in materials science. The unique crystal structure and diverse atomic interactions inherent in this double perovskite make it an up-and-coming candidate for various technological applications, particularly in photovoltaics; owing to its stability and resistance to heat and humidity, we aim to shed light on the extraordinary potential of K_2AgBiI_6 . Our study provides valuable insights for researchers engaged in tailored material design. We anticipate that the exceptional electronic properties of K_2AgBiI_6 will not only redefine the boundaries of materials engineering but also catalyze unprecedented advances in sustainable technology. Employing the powerful computational tool CASTEP, we conducted detailed electronic structure calculations within the framework of Density Functional Theory (DFT) to unravel the electronic properties of the double perovskite K_2AgBiI_6 . Our investigation thoroughly explored structural properties, band structure, total density of states (DOS), and partial density of states (PDOS). Furthermore, we systematically examined the influence of different exchange-correlation functionals, including LDA, GGA, and m-GGA, on the electronic and optical features of the material by presenting a comparative analysis of these approximations.

KEYWORDS

double perovskite (K_2AgBiI_6), solar cell, electronic properties, optical properties, structural properties, DFT

1 Introduction

Perovskite materials have garnered immense attention in recent years due to their exceptional electronic and optoelectronic properties. They are promising candidates

for various technological applications such as solar cells, light-emitting diodes, and photodetectors. Among the diverse range of perovskite compositions, the emerging class of double perovskites, characterized by their $A_2BB'X_6$ stoichiometry, has piqued significant interest owing to their intriguing structural features and potential for novel electronic behaviors. (Ali et al., 2021; Mir and Gupta, 2021; Dar et al., 2020). One such compound, K_2AgBiI_6 , presents a captivating platform for exploring the rich landscape of perovskite properties. Comprising potassium (K), silver (Ag), bismuth (Bi), and iodine (I) atoms, K_2AgBiI_6 exhibits a unique crystal lattice with remarkable electronic interactions that can be tailored to achieve specific functional characteristics. Given the importance of accurate theoretical predictions in understanding and optimizing material properties, we conducted a comprehensive investigation of K_2AgBiI_6 using the cutting-edge CASTEP software (Dar et al., 2020; Benmhidi et al., 2017). Within $A_2BB'X_6$, the Ag-containing halide double perovskite had excellent optoelectronic and thermoelectric properties (Bourachid et al., 2020; Saeed et al., 2022). Investigated Cs_2AgCrZ_6 (Z = Cl, Br, and I) through the density functional theory and recommended these materials for thermoelectric applications (Meyer et al., 2018; Lozhkina et al., 2018). Density Functional Theory (DFT) has proven to be an invaluable tool for elucidating the electronic properties of materials, offering insights into band structures, electronic densities, and optical behaviors. To delve into the electronic intricacies of K_2AgBiI_6 , we harnessed the power of three distinct DFT approximations: Local Density Approximation (LDA), Generalized Gradient Approximation (GGA), and meta-generalized Gradient Approximation (m-GGA). We employed these approximations to unravel the diverse aspects of electronic structure, bandgap behavior, and charge distribution within K_2AgBiI_6 . In this study, we present a comprehensive analysis of the electronic and optical properties of K_2AgBiI_6 , unveiling the subtle nuances that emerge from each approximation. Through a systematic comparison of the electronic structures obtained from LDA, GGA, and m-GGA calculations, we aim to provide a deeper understanding of each approach's unique features and challenges. Our investigation offers valuable insights into the nature of electronic interactions within K_2AgBiI_6 and lays the foundation for informed design strategies in harnessing its potential for next-generation optoelectronic devices (Bourachid et al., 2020; Meyer et al., 2018). In the following sections, we will detail the methodology employed, present our computational results, and discuss the implications of our findings in the context of perovskite materials and their applications (Zhao et al., 2018; Wang et al., 2020). This study serves as a testament to the symbiotic relationship between theoretical exploration and experimental advancement as we strive to unlock the captivating properties of K_2AgBiI_6 and further contribute to the exciting field of perovskite-based electronics.

2 Method of calculations

Applying the first-principles method based on plane wave pseudo-potentials and within the framework of Density Functional Theory (DFT), we used the Cambridge Sequential Total Energy Package (CASTEP) module to conduct our computational analysis. The electronic exchange-correlation interactions were managed

utilizing the Perdew-Burke-Ernzerhof (PBE) approach, which includes several approximations such as the Local Density Approximation (LDA), Generalized Gradient Approximation (GGA), and meta-generalized Gradient Approximation (m-GGA) projects. To describe the interactions between electrons and ions, we employed the OTFG Ultrasoft pseudo-potential. In the investigation of solid materials, it is paramount to attain fully relaxed structures with accurately converged total energies, forces, and stresses (Cs et al.; Perdew et al., 1996). This process is crucial for the exploration of the electronic and optical properties of materials. Consequently, achieving stable states of relaxed structures requires meticulous attention to convergence criteria for atomic positions and lattice parameters. To perform calculations within the Brillouin zone integration, we adopted the monkhorst-pack scheme, utilizing a k grid point resolution set at $4 \times 4 \times 4$.

To ensure the convergence of energy and configuration at a nearly complete plane wave basis level, we established a plane wave cut-off energy of 500 eV for the system. During the process of self-consistent field iteration, we employed the Pulay density mixing technique, setting the self-consistent field threshold at 2×10^{-5} eV/atom. We used the Broyden-Fletcher-Goldfarb-Shanno (BFGS) algorithm to optimize the model's structure, where specific optimization parameters were utilized. These parameters encompassed a convergence criterion of 0.05 eV/Å for interatomic interactions, a monatomic energy convergence standard of 2×10^{-5} eV, a threshold of 10 GPa for internal stress convergence within the crystal, and a maximum criterion of 0.002 Å for atomic displacement convergence. The program effectively achieved the desired convergence by optimizing these four parameters.

3 Structural properties

3.1 Tolerance factor

Figure 1 shows the crystal structure of the K_2AgBiI_6 double perovskites compounds. To ascertain these structural configurations, the unit cells underwent an initial optimization process employing energy minimization, followed by the relaxation of internal parameters. This meticulous exploration was undertaken to delve into the nuances of these architectures. Subsequently, Murnaghan's equation of state was employed to deduce the lattice constants of the double perovskite structures in their ground state (Del Campo et al., 2012; Al-Hattab et al., 2020). The lattice parameter's computed value is 11.85762 \AA . The tolerance factor (T_f) (Del Campo et al., 2012; Al-Hattab et al., 2020). K_2AgBiI_6 for this compound, as determined by Equation 1, is approximately 0.95. This value signifies the compound's stability within the cubic phase; the allowable range for (T_f) to maintain stability and form a cubic structure is typically between 0.9 and 1.1.

$$T_f = \frac{(R_K + R_I)}{\sqrt{2} \left(\left(\frac{R_{Ag} + R_{Bi}}{2} \right) + R_I \right)} \quad (1)$$

The double perovskite compound K_2AgBiI_6 exhibits a crystalline structure that aligns with cubic symmetry, characterized by the Fm3m (#225) space group. The constituent atoms occupy specific functional coordinates in this arrangement, contributing to the

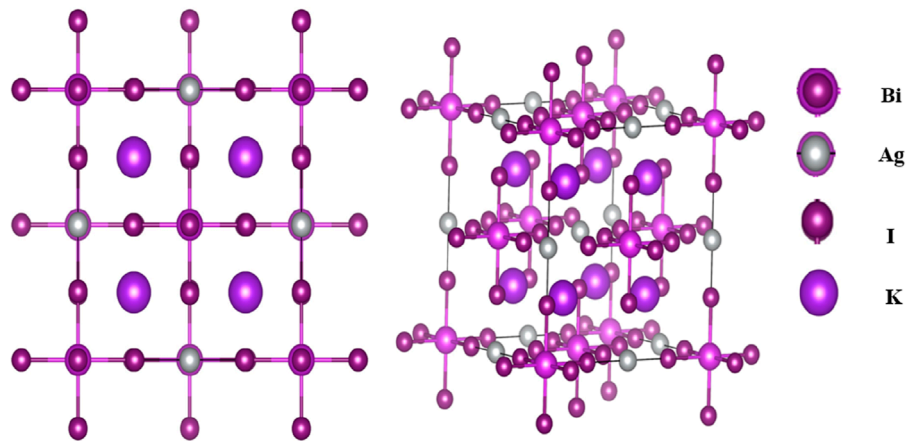


FIGURE 1
 K_2AgBiI_6 double perovskite structure.

overall crystalline framework. The potassium (K) atoms are positioned at fractional coordinates (0.25, 0.25, 0.75), occupying the 8c Wyckoff site. Similarly, the silver (Ag) atoms are situated at fractional coordinates (0.25, 0, 0), residing within the 4a Wyckoff site. On the other hand, the bismuth (Bi) atoms are coordinated at the origin (0, 0, 0), aligning with the 4b Wyckoff position. Completing this composition, the iodine (I) anions are located at fractional coordinates (0.253643, 0, 0), corresponding to the 24e Wyckoff site.

This arrangement of atoms within the crystal lattice highlights the precise organization and distribution of each element, thereby defining the unique structural properties of the double perovskite K_2AgBiI_6 . Such insight into the positional specifics of these constituent atoms enhances our understanding of the material's behavior and its potential applications in various scientific and technological domains (Meyer et al., 2018). We have determined the effective masses of electrons in the conduction band and holes in the valence band to analyze these particles' characteristics. The diagonal elements of influential mass tensor m_c are classically dined according to the following formula (2) (Zhao et al., 2018):

$$\frac{1}{m_c(k)} = \frac{1}{\hbar^2} \frac{\partial^2 E(k)}{\partial k^2} \quad (2)$$

The effectiveness of electron (m_e)-hole (m_h) pairs are ($m_e = 0.094, m_h = 2.13$), ($m_e = 0.087, m_h = 2.42$) and ($m_e = 0.074, m_h = 2.07$) for LDA, GGA, and mGGA, respectively.

3.2 Material thermodynamic stability

To assess the thermodynamic stability of this double perovskite, we calculated the formation energy of this compound through chemical Equation 3:

$$\Delta H_f = E_{tot}(K_2AgBiI_6) - [2E(K) + E(Ag) + E(Bi) + 6E(I)] \quad (3)$$

The symbol E_{tot} (K_2AgBiI_6) denotes the comprehensive energy aggregate inherent in the examined halides. Conversely, $E(K)$,

$E(Ag)$, $E(Bi)$, and $E(I)$ correspond to the energies attributed to the separate K, Ag, Bi, and I atoms, respectively. The negative magnitude of ΔH_f signifies the quantified energy dissipation accompanying the progression of material formation. This value encapsulates the reduction of essential energy during this process. Recognizing that the ensuing material stability is intrinsically tethered to this energy alteration is imperative. As a result, the computed numerical values of K_2AgBiI_6 , amounting to - 3.592 eV, substantiate the foundational thermodynamic resilience of the resultant synthesized material (Wang et al., 2020). Given that solar panels operate within elevated temperature conditions, the necessity for thermal resilience is paramount. The established thermodynamic stability of the $K_2AgSbBr_6$ material renders it particularly well-suited for application in photovoltaic solar cells. The valence electrons considered in the calculations were $K-4s^1$, $Ag-4d^{10}5s^1$, $Bi-4f^{14}5d^{10}6s^26p^3$, and $I-5s^25p^5$, respectively. These electron configurations were crucial in accurately describing the electronic properties of the system under investigation.

The compound K_2AgBiI_6 exhibits a cubic crystal structure derived from the perovskite arrangement, characterized by its placement within the cubic $Fm\bar{3}m$ space group. Within this structure, the potassium cations (K^{1+}) are intricately linked to twelve identical iodine anions (I^{1-}), resulting in the formation of KI_{12} cuboctahedra. These cuboctahedra are interconnected through shared corners with twelve other equivalent KI_{12} cuboctahedra, shared faces with six equivalent KI_{12} cuboctahedra, and further shared faces with four equivalent AgI_6 octahedra as well as four equivalent BiI_6 octahedra. The lengths of all potassium-iodine (K-I) bonds measure 4.30 Å. The silver cations (Ag^{1+}) create AgI_6 octahedra in this structure by bonding with six identical iodine anions (I^{1-}). These octahedra are characterized by their shared corners with six equivalent BiI_6 octahedra and their shared faces with eight identical KI_{12} cuboctahedra. Importantly, these octahedra maintain a non-tilted configuration. The distances between silver and iodine (Ag-I) bonds in these octahedra measure 2.99 Å.

Likewise, the bismuth cations (Bi^{3+}) form BiI_6 octahedra by bonding with six identical iodine anions (I^{1-}). These octahedra share corners with six equivalent AgI_6 octahedra and faces with

eight identical KI_{12} cuboctahedra. Similar to the silver-containing octahedra, these bismuth-based octahedra also maintain a non-tilted orientation. The lengths of bismuth-iodine (Bi-I) bonds in these octahedra measure 3.08 Å. The iodine anions (I^{1-}) themselves are linearly bonded to four equivalent potassium cations (K^{1+}), one silver cation (Ag^{1+}), and one bismuth cation (Bi^{3+}). This intricate network of bonding interactions and the precise spatial arrangement of these constituent elements define the unique structural characteristics of K_2AgBiI_6 , shedding light on its potential properties and applications in diverse scientific and technological contexts. Based on databases from materials project (Ali et al., 2023; Jain et al., 2013).

4 Results and discussion

4.1 Electronic propriety

4.1.1 Band structure

Band structure is a critical concept in condensed matter physics and materials science that describes the distribution of energy levels, or bands, in a solid material. It provides information about electrons' electronic properties and behavior in a crystalline solid. Carried out research on $Cs_2BiAgBr_6$ and reported it as an indirect bandgap ($E_g = 1.728$ eV) compound (Ali et al., 2021). Their experiment and theory also confirmed the indirect band semiconducting nature of $Cs_2BiAgBr_6$ along with $Cs_2BiAgCl_6$ (Dong et al., 1979). Suggested Rb_2AgBiI_6 for solar cell applications because of its good optical properties (Mathew et al., 2019), recommend employing $Cs_2AgInCl_6$ ($E_g = 1.1$ eV) for its favorable optoelectronic attributes. In a separate study (Nazir et al., 2023; Madsen and Singh, 2006).

The negative side of energy describes the valence band (VB), while the positive energy side shows the conduction band (CB), separated by the Fermi level E_f . Figure 2 gives us the value of the optical gap of K_2AgBiI_6 in the three approximations; the values of these bandgaps decrease by changing approximation. It has 0.439 eV, 0.653 eV, and 1.078 eV for LDA, GGA, and m-GGA approximations, respectively. The band gap results are similar to this work when it has a value of 0.857 eV using the GGA approximation (Bouich et al., 2022). The band gap of the double perovskite K_2AgBiI_6 is experimentally determined to be approximately 1.95 eV (Chen et al., 2024). This value aligns with similar lead-free perovskites, making it suitable for photovoltaic applications. Other related materials, like $Cs_2AgBiBr_6$, show slightly higher band gaps around 2.2 eV (Dar et al., 2019). The width of this bandgap is deemed sufficiently small to manifest characteristics akin to those of semiconductors. In the context of the three approximations, we have identified an indirect gap [X, L]. Observably, both the upper and lower regions of the conduction band demonstrate pronounced dispersion in proximity to W, L, and X points. This observation implies substantial mobility, suggesting favorable charge transport properties.

Local Density Approximation (LDA): LDA is the simplest approximation, where the exchange-correlation function is based only on the electron density at each point in the crystal. It assumes that the exchange-correlation energy depends only on the local electron density and is the same for all crystal regions. Although LDA can work well for some materials, it often underestimates the

bandgap energy in most cases. Indeed, it neglects some nonlocal correlation effects.

Generalized Gradient Approximation (GGA): GGA is a more sophisticated approximation that considers the electron density and the gradient of the electron density. By accounting for the gradient, GGA can capture some of the nonlocal correlation effects that LDA lacks, improving the accuracy of various properties, including bandgap energy. As a result, GGA tends to provide more considerable bandgap energies than LDA.

Meta-Generalized Gradient Approximation (mGGA): mGGA is even more advanced and accurate than GGA. It includes higher-order derivatives of electron density, allowing it to capture more complex nonlocal correlation effects. The increased accuracy of mGGA often leads to more considerable bandgap energies than both LDA and GGA (Parrey et al., 2018; Sharma et al., 2016).

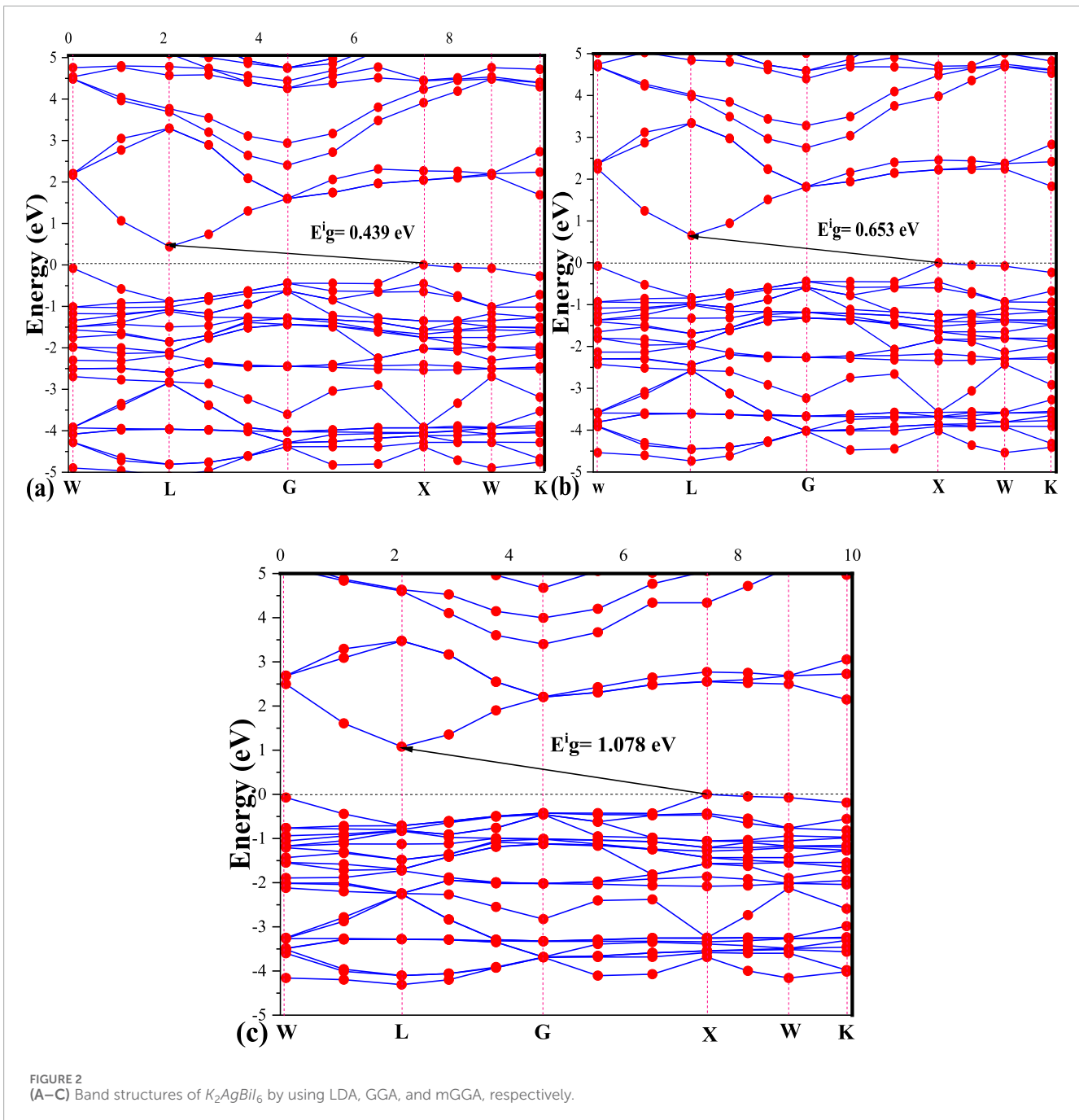
The increase in bandgap energy from LDA to GGA to mGGA can be explained by the progressive improvement of the approximations used to describe the exchange-correlation function. As you move from LDA to GGA and then to mGGA, the model becomes more sophisticated and incorporates more accurate descriptions of electronic interactions, resulting in better predictions of bandgap energies. It is important to note that the choice of the exchange-correlation function is crucial in electronic structure calculations, and the most appropriate level of approximation depends on the specific material and properties under study. In some cases, mGGA might be the most accurate choice, while in others, GGA might suffice. It is always good practice to compare the results with experimental data, if available, to validate the chosen approximation and methodology (Saeed et al., 2022; Karouchi et al., 2023).

4.1.2 Density of state

Density of States (DOS) is a fundamental concept in condensed matter physics that provides a statistical description of the distribution of energy states available to electrons in a material. It represents the number of electronic states per unit energy interval at each energy level (Perdew et al., 1996). Figure 3 of the DOS shows that in the valance band, the sharper peaks at around -0.7 eV decreased from 17, 21, and 24 states/eV, respectively, for LDA, GGA, and mGGA approximations. The DOS peaks in the valance region in the energy range increased with the LDA and GGA, compared with the m-GGA method. In the conduction band, every graph of approximation begins in our band gap; these three graphs are, in general, close to each other, with a phase shift equal to the difference between values of the gap.

The plotted Partial Density of States (PDOS) graphs in Figure 4 also provide insight into the electronic contributions. Notably, the p-orbitals associated with the I-atom exhibit significant contributions to the overall Density of States (DOS) within the valence region spanning from -2 to 0. In this range, the contributions from the other atoms are notably marginal. In contrast, within the conduction region, it becomes apparent that the p-states originating from the Bi-atom exert the most substantial influence on the DOS. It is worth mentioning that the DOS peaks, attributable to both the I and Bi atoms, indicate the pronounced presence of halogen and post-transition compounds in this material.

Despite the presence of peaks in the DOS, it is essential to note that due to the discernible band gap and the



distinct peaks just slightly above the Fermi energy (E_f) (Moulaoui et al., 2023), we draw a preliminary inference that this compound exhibits properties consistent with a semiconducting (Moulaoui et al., 2022a; Moulaoui et al., 2022b). This assessment is based on the approaches of local density approximation (LDA), generalized gradient approximation (GGA), and meta-generalized gradient approximation (m-GGA). The PDOS plots for the K_2AgBiI_6 double perovskite plot the contribution of different atomic orbitals to the electronic density of states (DOS). The x -axis (eV) of the graph corresponds to the energy levels, while the y -axis represents the density of states (Electrons/eV) or intensity of the PDOS. The PDOS is calculated by projecting the electronic states obtained

from the DOS calculation onto the localized atomic orbitals of each atom in the K_2AgBiI_6 unit cell. The PDOS plot shows distinct peaks and features that provide valuable insight into the electronic structure and bonding behavior of the K_2AgBiI_6 double perovskite. Each peak is associated with specific atomic orbitals and reflects the contribution of different electron energy levels to the total DOS. Peaks in the valence band region of the PDOS plot correspond to occupied electronic states below the Fermi level. These states represent the electrons present in the valence band of the material. The peaks can be attributed to the contribution of specific atomic orbitals, such as the p and d orbitals of K-site cations, the p and d orbitals of Ag-site

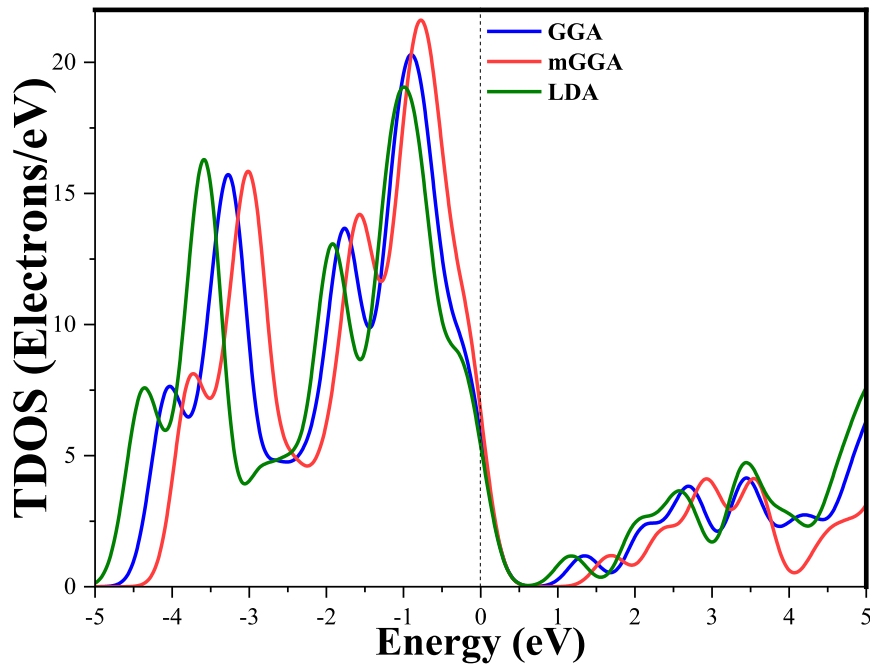


FIGURE 3
Total density of state of K_2AgBiI_6 .

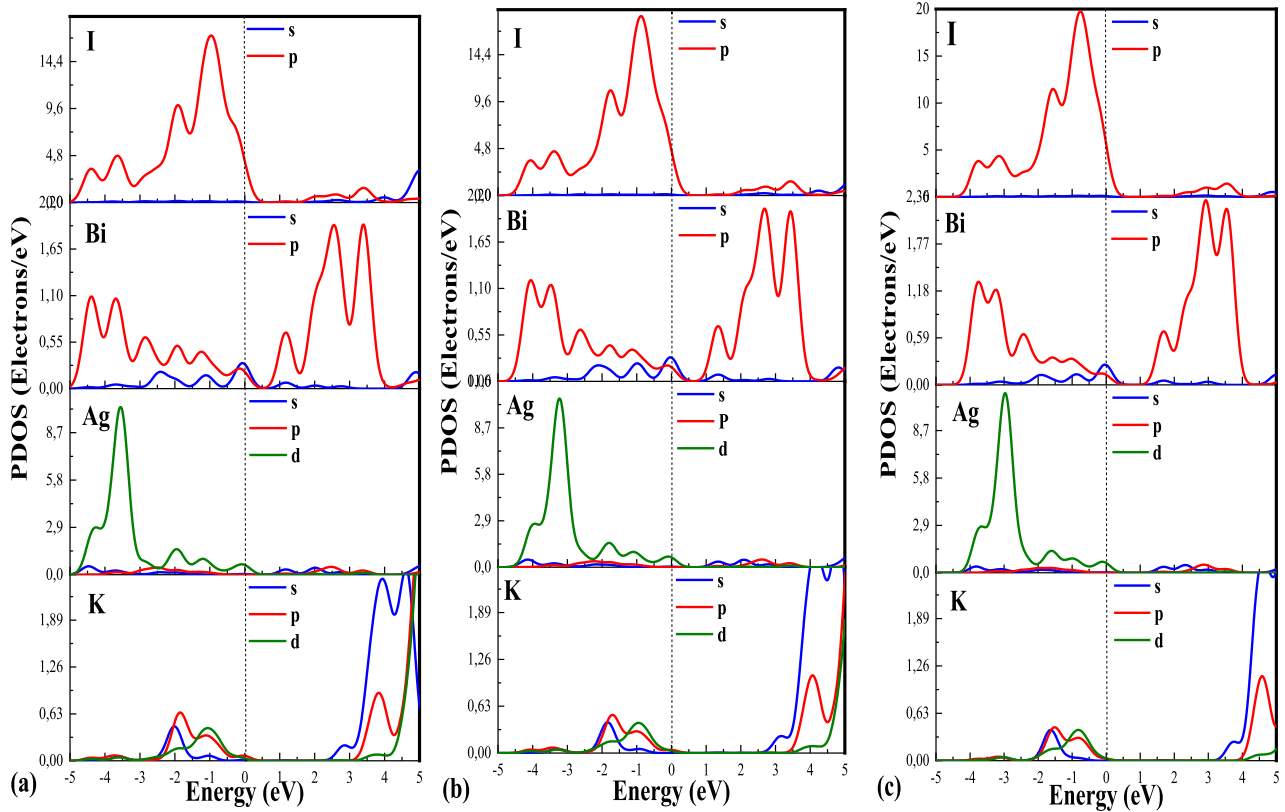
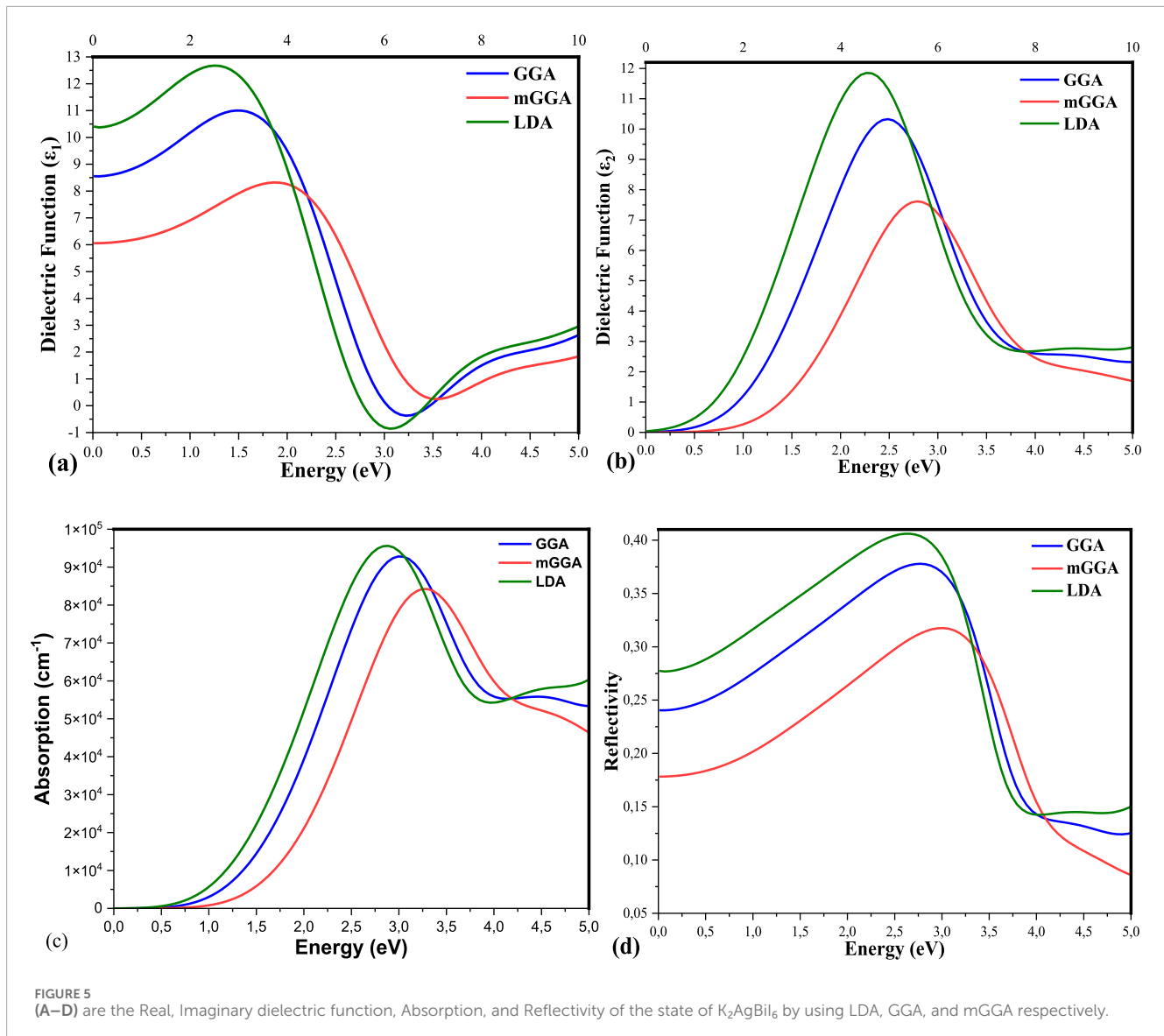


FIGURE 4
(A–C) are the partial densities of the state of K_2AgBiI_6 by using LDA, GGA, and mGGA, respectively.



cations, the s and p orbitals of Bi-site cations, and the p orbital of I-site anions.

The peaks in the conduction band region of the PDOS plot represent the unoccupied electronic states above the Fermi level (Al-Hattab et al., 2022; Archi et al., 2023). These states are responsible for the material's electronic conductivity and charge transport properties (Najim et al., 2022; Bajjou et al., 2022). The conduction band peaks are mainly attributed to the contribution of the p orbitals of the Bi-site anions. The intensity and position of these peaks can provide information about the bandgap energy of the material and its suitability for specific electronic applications. The features observed in the PDOS plot provide insight into the bonding and electronic properties of the K_2AgBiI_6 double perovskite. The PDOS plot can also reveal the presence of impurity states or defect levels that influence material properties. By describing the peaks and curves of the PDOS plot of the double perovskite K_2AgBiI_6 using this approach, researchers can

gain a comprehensive understanding of its electronic structure, bonding behavior, and potential applications in various fields of materials science.

5 Optical properties

5.1 Dielectric function, absorption and reflectivity

When investigating a material's optical characteristics, it is critical to calculate and interpret the complicated dielectric function. The function is essential since it allows us to calculate the other optical qualities. Its optical characteristics. Furthermore, the imaginary portion, $\epsilon_2(\omega)$, this complex dielectric function offers information about absorption in the material in the form of peaks associated with the transitions of interband electrons

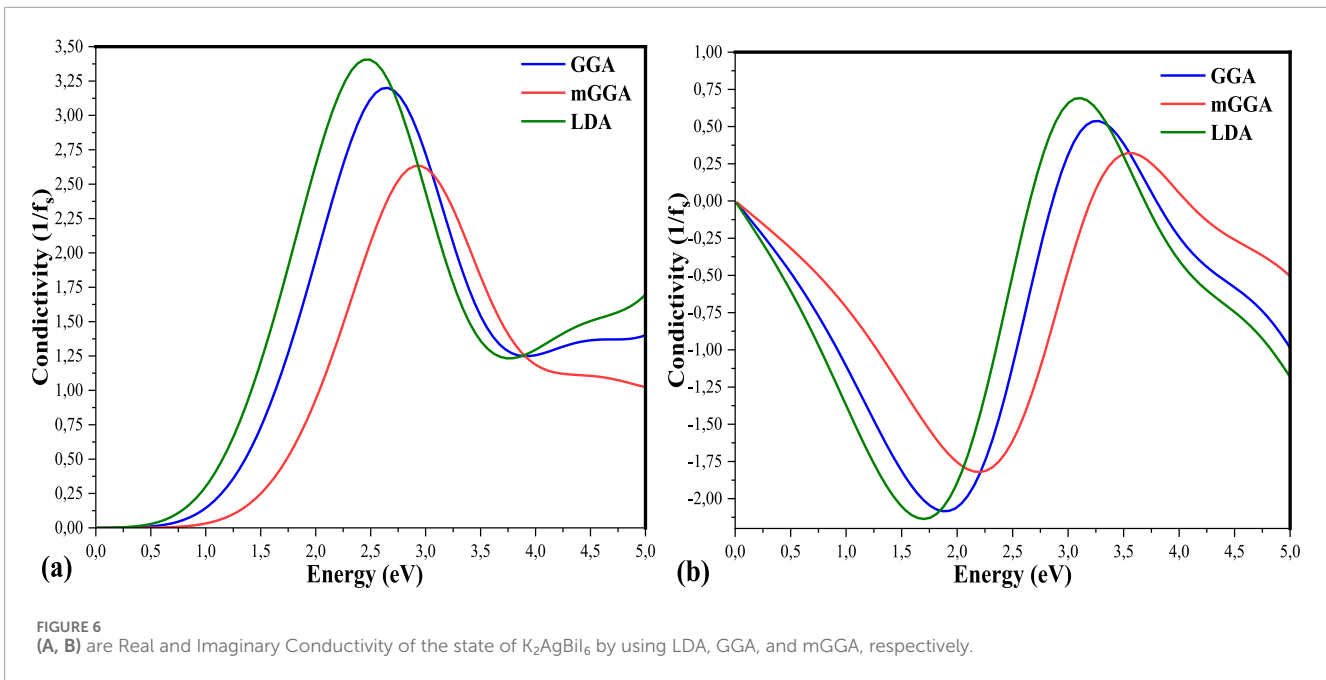


TABLE 1 Maximum values calculated for K_2AgBiI_6 , Dielectric constant (real $\epsilon_{1,max}$ and imaginary $\epsilon_{2,max}$ parts), Absorption coefficient (α_{max}), Reflectivity (R_{max}), Optical conductivity (real $\sigma_{1,max}$ and imaginary $\sigma_{2,max}$ parts), Refractive index (n_{max}) and extinction coefficient k_{max} , Energy loss function (L_{max}).

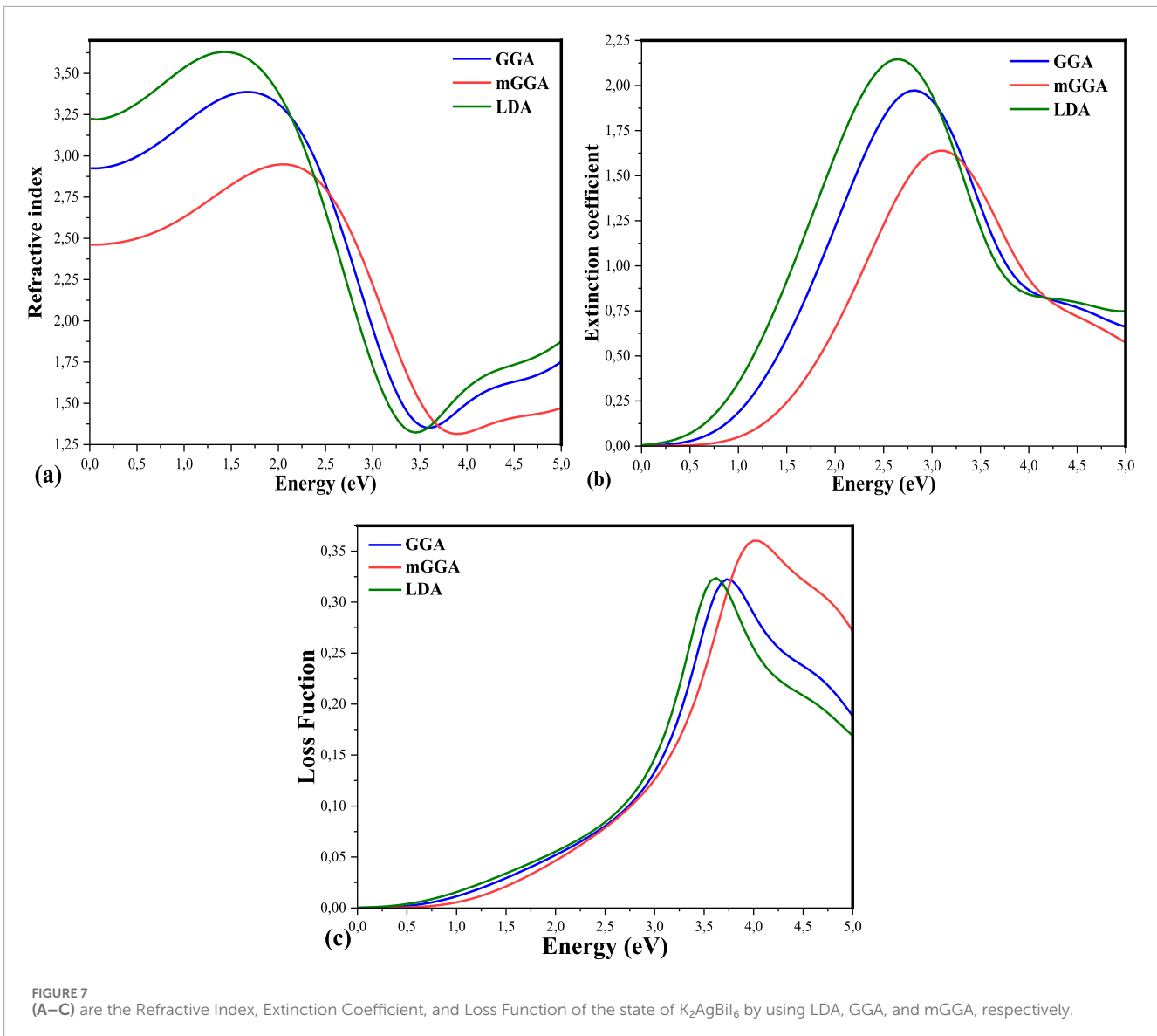
Approximation	LDA	GGA	m - GGA
$\epsilon_{1,max}$	(E = 1.25) 12.68	(E = 1.53) 11	(E = 1.91) 8.314
$\epsilon_{2,max}$	(E = 2.268) 11.85	(E = 2.48) 11.32	(E = 2.78) 7.61
α_{max}	(E = 2.88) 9.56 * 10 ⁴	(E = 3) 9.28 * 10 ⁴	(E = 3.30) 8.43 * 10 ⁴
R_{max}	(E = 2.66) 0.406	(E = 2.82) 0.377	(E = 3.07) 0.316
$\sigma_{1,max}$	(E = 2.494) 3.405	(E = 2.657) 3.2	(E = 2.95) 2.632
$\sigma_{2,max}$	(E = 3.11) 0.69	(E = 3.276) 0.537	(E = 3.53) 0.321
n_{max}	(E = 1.477) 3.62	(E = 1.75) 3.38	(E = 2.09) 2.94
k_{max}	(E = 2.66) 2.145	(E = 2.88) 1.96	(E = 3.07) 1.638
L_{max}	(E = 3.62) 0.323	(E = 3.727) 0.322	(E = 3.99) 0.36

from the valence band to the conduction band (Bhorde et al., 2021; Mathew et al., 2019). Figures 5A,B illustrates the variations of the dielectric function's imaginary part $\epsilon_2(\omega)$ and real part $\epsilon_1(\omega)$ in an energy interval ranging from 0 eV to 5 eV for K_2AgBiI_6 compounds using various approximations based on DFT (LDA, GGA and mGGA). For our compounds, $\epsilon_1(\omega)$ starts from zero frequency and increases to its maximum value, then decreases. The value of $\epsilon_1(0)$, then increases from mGGA to LDA. An important observation that can be made from the graphs of the real part of the

dielectric function is that $\epsilon_1(\omega)$ decreases to negative values. The negative values of $\epsilon_1(\omega)$ show that, in this energy region, incident electromagnetic waves are mainly reflected by the material. Hence, the material exhibits metallic behavior and can be used in radiation protection layers in this energy range. The variation of the imaginary part of the dielectric function as a function of energy is also shown in Figure 5B for the compound K_2AgBiI_6 . We note that the critical points (appearance) in the spectra of the different structures are shifted to lower energy when moving from mGGA to LDA. The imaginary dielectric function spectra for the compound K_2AgBiI_6 are composed of the peaks. These peaks correspond to the transition of electrons from the valence band to the conduction band. The absorption coefficient α of K_2AgBiI_6 is shown in Figure 5C. The absorption maximum increases with the change of approximation from mGGA to LDA. In the same Figure 5D, we plot the reflectivity R of K_2AgBiI_6 . Note that the value also increases with the change of approximation from mGGA to LDA. The highest reflectivity values are recorded for three approximations in the visible range.

5.2 Optical conductivity

Electron conduction is due to an applied electromagnetic field characterized by optical conductivity (Nazir et al., 2023; Ejjabli et al., 2024). Figures 6A, B shows that the real and imaginary optical conductivity $\sigma(\omega)$ for our compounds starts around the values corresponding to the material gap. After these points, it increases and reaches its maximum value, then decreases with minor variations. The variation in approximations shifts the maximum values of real conductivity to lower energies. The maximum values of the real and imaginary parts of optical conductivity are shown in Table 1.



5.3 Refractive index, extinction coefficient, and loss function

Knowledge of the refractive index dispersion n in the transparency region is crucial for designing various optoelectronic devices (Ali et al., 2023; Madsen and Singh, 2006). Figures 7A, B illustrates the variation of refractive indices for different approximation methods. It can be observed that the refractive indices exhibit an increase–decrease trend across all approximations. Regarding the extinction coefficient k , we observed an initial increase followed by a decrease and eventual stabilization for all three approximation methods.

Maximum refractive index values and extinction coefficients for K_2AgBiI_6 materials are shown in Table 1. To better understand the optical properties, we calculated the electron loss function L , an essential factor in describing the energy loss of a fast electron moving through a material (Perdew et al., 1996; Del Campo et al., 2012). Figure 7C shows the energy loss function of a fast electron

moving through K_2AgBiI_6 material. All approximations recorded similar peaks, and the maximum values of the energy loss function are shown in Table 1. The intensity and position of these peaks can provide information about the bandgap energy of the material and its suitability for specific electronic applications (Bouich, 2020; Bouich et al., 2022).

6 Conclusion

In this study, we delved into the intricate electronic properties of the perovskite material using a comprehensive approach by employing three distinct approximations: LDA, GGA, and mGGA. Our investigation aimed to unravel the multifaceted interplay between theoretical models and the electronic behavior of the perovskite structure. Through rigorous computations, we uncovered compelling insights into the nature of the electronic structure within the perovskite framework. We find indirect band gap equal

0.439 eV, 0.653 eV, and 1.078 eV for LDA, GGA, and mGGA approximations, respectively. With its simplicity and computational efficiency, the LDA approximation offered a valuable initial glimpse into the material's properties. The GGA approach refined our understanding by incorporating gradient corrections, accurately capturing electronic interactions, and shedding light on crucial bandgap characteristics. However, the most captivating revelations emerged from the mGGA approximation. By considering meta-generalized gradient corrections, we witnessed a remarkable convergence between theoretical predictions and experimental observations. This advanced approach discerned intricate electronic nuances and revealed the material's sensitivity to subtle electronic rearrangements. Our findings underscore the importance of methodological choices in modeling perovskite materials, with each approximation offering a distinct lens through which to perceive their electronic intricacies. As the pursuit of efficient and high-performance photovoltaic materials intensifies, our study serves as a testament to the significance of comprehensive theoretical investigations in guiding the design and optimization of materials for cutting-edge applications.

Data availability statement

The original contributions presented in the study are included in the article/Supplementary Material, further inquiries can be directed to the corresponding author.

Author contributions

MK: Conceptualization, Data curation, Formal Analysis, Investigation, Methodology, Writing—original draft. AE: Conceptualization, Formal Analysis, Investigation, Validation, Visualization, Writing—review and editing. OB: Investigation, Methodology, Resources, Writing—original draft. JG: Data curation, Investigation, Methodology, Visualization, Writing—review

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Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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