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A MEMS Coriolis-Based Mass-Flow-to-Digital Converter for Low Flow Rate Sensing

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***Abstract*— This article presents a microelectromechanical system (MEMS) Coriolis-based mass-flow-to-digital converter (ΦDC) that can be used with both liquids and gases. It consists of a micromachined Coriolis mass flow sensor and a CMOS interface circuit that drives it into oscillation and digitizes the resulting mass flow information. A phase-locked loop (PLL) drives the sensor at its resonance frequency (f_D), while a low $1/f$ noise switched-capacitor (SC) proportional–integral (PI) controller maintains a constant drive amplitude. Mass flow through the sensor causes Coriolis-force-induced displacements, which are detected by co-integrated sense capacitors. In-phase (I) and quadrature (Q) components of these displacements are then digitized by two continuous-time delta–sigma modulators (CT-ΔΣMs) with finite impulse response (FIR)-DACs and passive mixers. Their outputs are used to accurately estimate and cancel sense path delay, thus improving sensor stability. To ensure constant sensitivity over a wide range of fluid densities, a background sensitivity tuning (BST) scheme adjusts the sense capacitors' bias voltage as a function of f_D, which is a good proxy for fluid density. Implemented in a standard 0.18-μm CMOS technology, the interface circuit consumes 13 mW from a 1.8-V supply. The proposed MEMS Coriolis ΦDC achieves a state-of-the-art noise floor of 80 μg/h/$\sqrt{\text{Hz}}$ and a zero stability (ZS) of ±0.31 mg/h, which is at par with MEMS thermal flow sensors.**

***Index Terms*— Capacitive interface, CMOS readout interface, Coriolis mass flow sensor, delta–sigma modulator (ΔΣM), flow rate, flowmeter, demodulator, low noise, microelectromechanical systems (MEMS), microfluidics.**

I. Introduction

FLOW sensors are required in many industrial and microfluidic applications, such as flow cytometry [1], biological/chemical assays [2], and volatile organic compound detection [3], [4]. Their primary function is to translate fluid flow into analog or digital information, thus allowing it to be measured and controlled. However, detecting the low flow rates, often less than 1 g/h, associated with microfluidic applications requires sensors with high resolution (<200 μg/h/$\sqrt{\text{Hz}}$) and low drift (<±0.4 mg/h). They should also be able to maintain good accuracy in the presence of variations in fluid viscosity and density. Last but not least, they should be as small as possible to facilitate their use in a wide range of microfluidic systems, which are often volume constrained.

Flow sensors based on microelectromechanical systems (MEMS) are well suited for microfluidic and low flow industrial applications [5]. Compared with conventional flow sensors, they can be manufactured in volume using wafer-level batch fabrication techniques, which, in turn, reduces process variability and costs. In addition, they can be co-integrated and packaged with CMOS circuitry, which confers ease of use as well as a small system size.

Several MEMS/CMOS flow sensors suited for low flow applications have been reported [6], [7], [8], [9], [10], [11], [12], [13], [14], [15], [16], [17], [18]. Thermal flow sensors [6], [7], [8], [9], [10] are among the most popular. Such sensors typically employ a few temperature sensors located around a heater to measure the temperature asymmetry caused by fluid flow. They have no moving parts and, thus, can achieve high stability and resolution (<100 μg/h/$\sqrt{\text{Hz}}$). However, since they sense mass flow indirectly via fluid heat capacity and temperature differences, thermal flow sensors have to be calibrated for use with specific liquids [8] or gases [9], [10] and, consequently, cannot be used for unknown mixtures of fluids. Furthermore, their output tends to saturate at higher flow rates, limiting their dynamic range [8], [10].

Coriolis flow sensors [11], [12], [13], [14], [15], [16], [17], [18], [19], [20], [21], on the other hand, directly measure mass flow by sensing the Coriolis force produced by a flowing mass in a vibrating channel. This force is independent of fluid properties, and so, the same sensor can be used to accurately (<±0.5%) measure the mass flow rate of both liquids and gases. In addition, fluid density can be estimated from the resonance frequency of the vibrating channel [22], [23]. However, commercially available Coriolis flow sensors are made of stainless steel (SS) [21] and so are bulky and expensive. They also suffer from poor stability (±10 mg/h) at low flow rates, since the associated Coriolis forces are quite small and, thus, difficult to detect [24].

In contrast, MEMS Coriolis flow sensors [11], [12], [13], [14], [15], [16], [17], [18], [19] can be much more sensitive. In [12], a sensor intended for low flow rates is presented, but its use of a monocrystalline silicon flow channel limits its mass and stiffness, resulting in low sensitivity and the need

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for vacuum packaging. Furthermore, it also suffered from a large (±200 mg/h) zero stability (ZS). Bulk-micromachining-based surface channel technology (SCT) [25], [26] enables the fabrication of silicon-rich silicon nitride (SiRN) flow channels with much lower mass and stiffness, resulting in high sensitivity at low flows [16]. In recent work, it has been shown that such sensors can be combined with CMOS circuitry to achieve compact drive and sensing functionality while maintaining high resolution (300 μg/h/$\sqrt{\text{Hz}}$) and low ZS (±0.8 mg/h) [14]. However, sub-100μg/h/$\sqrt{\text{Hz}}$ noise floor and further reduction of the ZS are required to make MEMS Coriolis a reliable alternative to MEMS thermal flow sensors in low flow and microfluidic applications.

In this work, an extended version of [15], a mass-flow-to-digital converter (ΦDC) based on a MEMS Coriolis mass flow sensor is presented. CMOS circuitry is used to implement both its drive loop and readout electronics, while also providing a compact digital-output system. To achieve low ZS, the circuitry is designed to have low $1/f$ noise, while residual quadrature (Q)-induced drift is cancelled in the digital domain. The proposed ΦDC achieves a noise floor of about 80 μg/h/$\sqrt{\text{Hz}}$ and a ZS of ±0.31 mg/h in a 4-Hz bandwidth, which corresponds to a 3.75× and 2.6× improvement, respectively, on the state of the art [14].

The rest of this article is organized as follows. Section II briefly introduces the MEMS Coriolis mass flow sensor's working principle and design considerations. Section III presents the proposed sensor interface architecture and some key circuit-level details. In Section IV, measurement results are discussed and compared with other state-of-the-art flow sensors. Finally, Section V concludes this article.

II. Micromachined Coriolis Mass Flow Sensor

A. Operating Principle

The operation of a MEMS Coriolis mass flow sensor is illustrated in Fig. 1. The sensor consists of a suspended vibrating flow channel with two main resonance modes: the drive (twist) and the sense (swing) modes. In the presence of a static magnetic field (B), the drive mode is actuated by passing an alternating current (i_D) through a metal track at the top of the suspended channel. This generates a Lorentz force ($F_L = i_D \times B \times L$) at the sections [$L$, in Fig. 2(a)] of the channel, where i_D is orthogonal to B [27], resulting in a drive displacement (x_D) at the outer corners of the channel, as shown in Fig. 1(a).

The drive mode can be modeled as a lumped-element mass-spring system with resonance frequency [23]

$$f_D = \frac{1}{2\pi}\sqrt{\frac{k_D}{m_{\text{ch}} + m_f}} = \frac{1}{2\pi}\sqrt{\frac{k_D}{m_{\text{ch}} + \rho V_{\text{ch}}}} \tag{1}$$

where k_D is the drive mode spring constant, m_{ch} is the mass of the empty channel, m_f is the fluid mass inside the channel, V_{ch} is the channel's volume, and ρ is the fluid density. From (1), it can be seen that the fluid density directly impacts the resonance frequency in proportion to the added fluid mass, allowing ρ to be estimated by measuring changes in f_D.

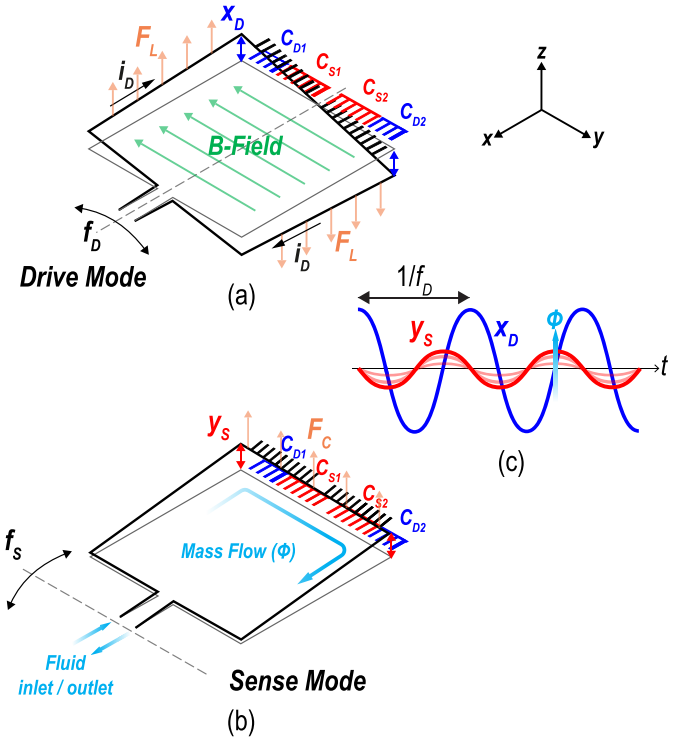

Fig. 1. Operation principle of the MEMS Coriolis mass flow sensor: movement when (a) driven at its drive mode resonance frequency, (b) sense mode is excited by the Coriolis force, and (c) drive (x_D) and sense (y_S) displacements in the time domain.

Fluid flowing through the channel with a mass flow rate Φ will experience a Coriolis force (F_C) with amplitude

$$F_C \approx 4\omega_D x_D \Phi. \tag{2}$$

This induces an out-of-plane sense displacement (y_S) that is proportional to the drive velocity, thus orthogonal to the drive mode displacement x_D, and modulated by f_D, as shown in Fig. 1(b) and (c). The Coriolis force is proportional to Φ and is independent of temperature, pressure, flow profile, and fluid properties. From a lumped model of the sense mode, assuming that it is harmonically driven at f_D, the Coriolis force-induced displacement amplitude $|y_S|$ can be expressed as follows:

$$|y_S(\omega_D)| \approx \frac{F_C}{k_S\sqrt{\left(1-\left(\frac{\omega_D}{\omega_S}\right)^2\right)^2 + \left(\frac{\omega_D}{\omega_S Q_S}\right)^2}} \tag{3}$$

where k_S, ω_S, and Q_S are the sense mode spring constant, resonance frequency ($f_S = \omega_S/2\pi$), and quality factor, respectively. Since the ratio (ω_D/ω_S) is almost independent of fluid density [18], and the amplitude of x_D can be regulated, the Coriolis force-induced y_S can be used as a frequency-dependent measure of mass flow.

From (1) and (3), to accurately sense mass flow and density, the sensor must be operated at f_D, while the amplitudes of y_S and x_D must be sensed and regulated, respectively.

B. Suspended Channel Design

High sensitivity is desired to measure mass flows below 1 g/h with high resolution. To achieve this, the mass of the

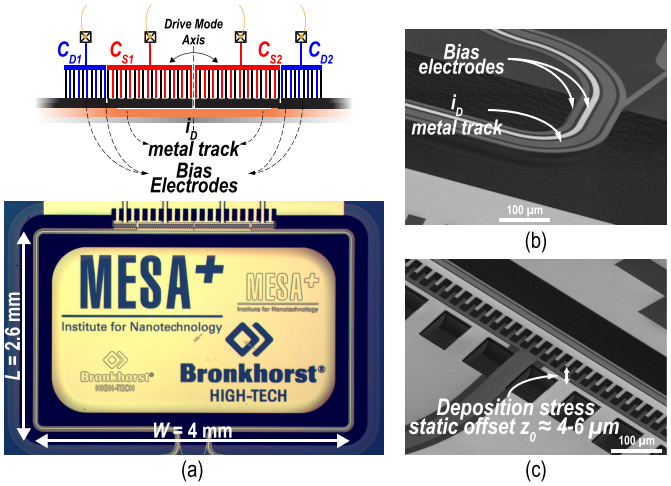

Fig. 2. (a) Micrograph of the SCT-fabricated MEMS Coriolis mass flow sensor and integrated readout capacitors illustration, (b) SEM images showing the access electrodes above the suspended channel, and (c) static offset z_0 between the readout electrodes.

TABLE I
TYPICAL PARAMETERS OF THE MEMS CORIOLIS MASS FLOW SENSOR

Parameter	Value	Unit
Suspended Channel Size (*WxL*)	4 x 2.6	mm^2
Channel Mass (m_{ch})	~15	µg
Channel Diameter (d_{in})	<60	µm
Wall Thickness (t_W)	1.2	µm
Drive Resonance Frequency (f_D)	1.78 - 3	kHz
Sense Resonance Frequency (f_S)	1.01 - 1.71	kHz
Sense Mode Q-Factor (Q_S)	40	-
Drive Displacement (x_D)	10	µm
Drive Current (i_{DRIVE})	5	mA
Readout Comb Offset (z_0)	4 - 6	µm
Sense DC Bias (V_{BS})	20	V
Mass Flow Sensitivity* ($\Delta C_S/\Phi$)	[300 - 400]	aF/(g/h)
Brownian Noise**	~1	µg/h/√Hz

*Estimate from N_2 to H_2O. **[19]

channel itself must be small compared with the mass of the fluid moving through it. Previous MEMS Coriolis designs used silicon to fabricate the sensor channel, leading to heavy, stiff, and low sensitivity sensors [11], [12]. Here, the sensor is fabricated using SiRN with a bulk-micromachining-based SCT process, as explained in detail in [25] and [26].

Fig. 2 shows the micrograph and scanning electron microscope (SEM) images of the fabricated rectangular shape device [15]. It consists of a suspended channel with a wall thickness (t_W) of 1.2 μm and an inner diameter (d_{in}) of about 60 μm [16]. The size of the channel [$W \times L$, in Fig. 2(a)] is limited by the need to keep m_{ch} small. In this work, $W = 4$ mm and $L = 2.6$ mm. This results in a (ω_D/ω_S) ratio of about 1.75, which ensures that y_S is quite insensitive to variations in Q_S, since, from (3), y_S will change by less than 0.1% in response to a ±50% variation in Q_S.

C. Capacitive Readout

To read out the x_D and y_S displacements, capacitive electrodes with 2.7-μm-thick comb fingers are integrated in the sensor, as shown in Figs. 1 and 2. Metal tracks are deposited on the top of the channel to access the electrodes and to provide a path for i_D, as shown in Fig. 2(b). The mechanical stress caused by these tracks creates a static offset (z_0) between the comb electrodes [see Fig. 2(c)], which ensures that the resulting capacitance change (ΔC) is a monotonic function of displacement. Furthermore, this design simplifies the fabrication process and reduces squeezed film damping, which results in a drive mode quality factor $Q_D > 60$ without vacuum encapsulation, thus allowing the sensor to operate at atmospheric pressure, unlike [12] and [22].

For large displacements, when the overlap between the electrodes tends to zero, the resulting ΔC will become a nonlinear function of displacement. To mitigate this, the detection combs ($C_{D1,2}$ and $C_{S1,2}$) are symmetrically placed around the drive mode rotation axis. This placement limits the maximum displacement experienced by the combs to less than z_0 (≈4 μm), thus ensuring adequate linearity in ΔC.

For x_D detection, the differential motion around the drive mode should be sensed. Compared with the comb pair $C_{S1,2}$, the pair $C_{D1,2}$ is more sensitive to x_D, since they are located further away from the drive rotation axis. So, $C_{D1,2}$ are used to detect drive motion, while $C_{S1,2}$ are used to detect sense motion. The $C_{D1,2}$ and $C_{S1,2}$ standoff capacitances are in the order of 25 and 45 fF and typically change by ±3 and ±1 fF each, respectively. Table I summarizes the typical values for the key mechanical parameters of the sensor used in this work.

D. Mass Flow Sensing Readout

Although combs $C_{S1,2}$ are intended to detect the sense motion, they still experience large changes due to x_D (ΔC_X) and only small changes due to y_S (ΔC_Y), as shown in Fig. 1. In the time domain, these capacitance changes are orthogonal to each other [see Fig. 1(c)], and so, C_{S1} and C_{S2} can be expressed as follows:

$$C_{S1} = \Delta C_Y \cos(\omega_D t) + \Delta C_X \sin(\omega_D t) \quad (4)$$

$$C_{S2} = \Delta C_Y \cos(\omega_D t) - \Delta C_X \sin(\omega_D t). \quad (5)$$

Consequently, the phase difference ($\Delta\varphi$) between C_{S1} and C_{S2} due to y_S can be expressed as follows [17]:

$$\Delta\varphi = 2 \arctan\left(\frac{\Delta C_Y}{\Delta C_X}\right) \propto \Phi \quad (6)$$

where $\tan(\Delta\varphi) \approx \Delta\varphi$ for small phase shifts. Therefore, mass flow (Φ) is directly proportional to the ratio $\Delta C_Y/\Delta C_X$ and so is robust to variations in the amplitude of x_D.

The resulting readout architecture is shown in Fig. 3(a), where each comb $C_{S1,2}$ is biased by a DC voltage (V_{BS}) and then readout by a capacitance-to-voltage (C/V) converter and

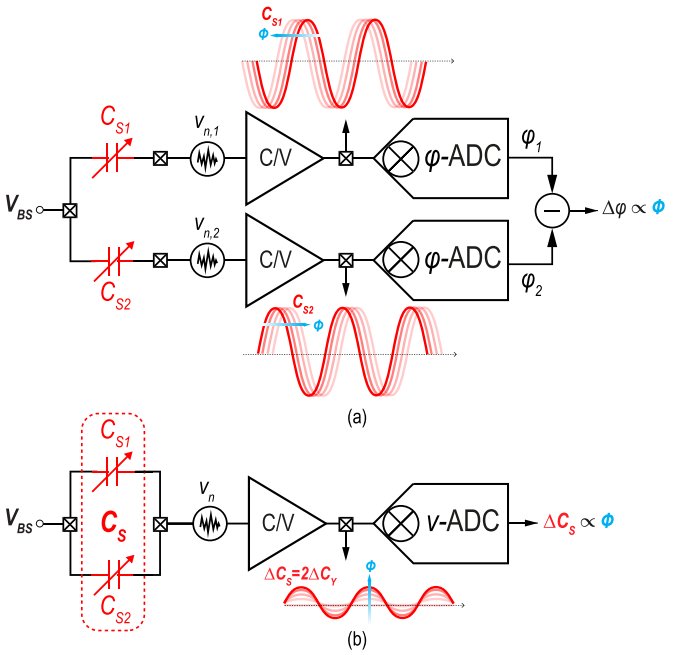

Fig. 3. (a) Phase-domain readout [14] and (b) single C/V voltage-domain readout.

a phase digitizer [13], [14], [16], [17], [18]. With this architecture, good performance has been achieved: a noise floor of $<300\ \mu$g/h/$\sqrt{\text{Hz}}$ and a ZS of $<\pm 1$ mg/h [14]. However, the need for two C/V converters increases the overall noise (v_n) by 3 dB, while the dynamic range of the analog-to-digital converters (ADCs) must be large enough to accommodate the large ΔC_X component present in $C_{S1,2}$.

Since the mass flow information is only contained in ΔC_Y, it can be extracted by simply connecting C_{S1} and C_{S2} in parallel. From (4) and (5), this cancels ΔC_X, resulting in a net sense capacitance change $\Delta C_S = 2\Delta C_Y$, as shown in Fig. 3(b). As a result, only a single C/V converter is required, which simplifies the front end, improves the noise floor by 3 dB, and relaxes the dynamic range requirements on the ADC. From (2), accurate control of x_D is now required, since variations in its amplitude will directly affect ΔC_S.

The minimum detectable mass flow for the front end in Fig. 3(b) is then determined by the sensor's Brownian noise [19], the equivalent current noise of the single C/V converter, and the sensor's mass flow sensitivity ($\Delta C_S/\Phi$).

III. Proposed Architecture and Circuit Design

A. System Overview

Fig. 4 shows the system overview of the proposed ΦDC. It consists of a drive loop, which drives the sensor at its resonance frequency f_D and regulates the amplitude of the drive displacement x_D, and a sense readout channel, which derives the mass flow output from the sense capacitance changes ΔC_S.

The drive loop consists of a C/V converter with a fixed drive bias voltage (V_{BD}), which reads out the drive capacitance changes ΔC_D. The output of the C/V converter is applied to a phase-locked loop (PLL), which locks to f_D, and an amplitude regulator, which maintains the ΔC_D amplitude

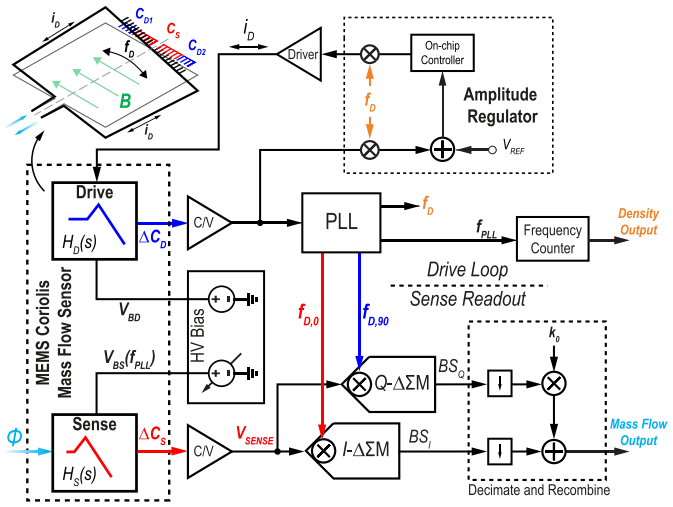

Fig. 4. System overview of the MEMS Coriolis ΦDC.

constant. Since f_D decreases with fluid density, the sensor's sensitivity with liquids is smaller than with gases, as expected from (3). To maintain constant sensitivity for different fluids, an external frequency counter measures f_D, providing the density information, and is then used to adjust the sense capacitors' bias voltage (V_{BS}) accordingly.

The changes in sense capacitance ΔC_S are converted to a sense voltage V_{SENSE} by another C/V converter. In principle, the mass flow component of V_{SENSE} could then be extracted by a single in-phase (I) demodulator, since the sense capacitor C_S should only detect Coriolis-force-induced motion. However, due to manufacturing tolerances, a portion of the drive motion will also couple to the sense mode, thus causing mass-flow-independent variations in C_S. Since this mechanical crosstalk is 90° phase-shifted relative to the Coriolis signal, it is commonly referred to as quadrature (Q) error and has been shown to be a significant source of drift in MEMS gyroscopes [28], [29], [30], [31], [32], [33], [34].

In this work, the symmetrical location of C_S around the drive mode axis provides the first-order attenuation of the Q error. However, any phase error (φ_{err}) between V_{SENSE} and the I demodulating reference ($f_{D,0}$) will also cause Q-induced drift. This is because the Q error can be expressed as an input-referred mass flow offset (Φ_Q) that is directly coupled to the input mass flow (Φ_I). The measured mass flow rate (Φ_O) can then be expressed as follows [29]:

$$\begin{aligned}\Phi_O &= \Phi_I \cos(\varphi_{\text{err}}) + \Phi_Q \sin(\varphi_{\text{err}}) \\ &\approx \Phi_I + \Phi_Q \varphi_{\text{err}} \end{aligned} \tag{7}$$

which shows that the drift in the measured mass flow rate Φ_O may be caused either by φ_{err} and Φ_Q. Furthermore, φ_{err} includes phase noise, which can further increase the noise floor in proportion to Φ_Q amplitude. For the MEMS Coriolis sensor used in this work, Φ_Q is often in the same order (or larger) as Φ_I at full scale.

The phase error φ_{err} will drift with both the sense-mode transfer function and delays in the C/V converter, while Φ_Q amplitude drift will come from manufacturing tolerances,

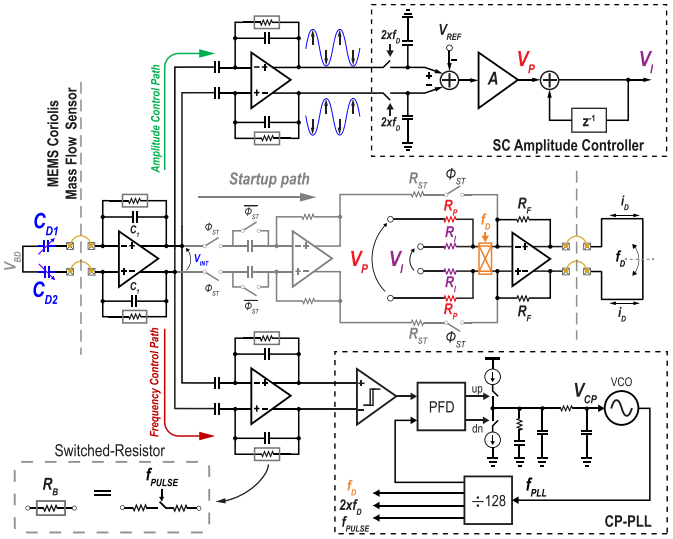

Fig. 5. Simplified circuit implementation of the drive loop showing the startup, frequency, and amplitude control paths.

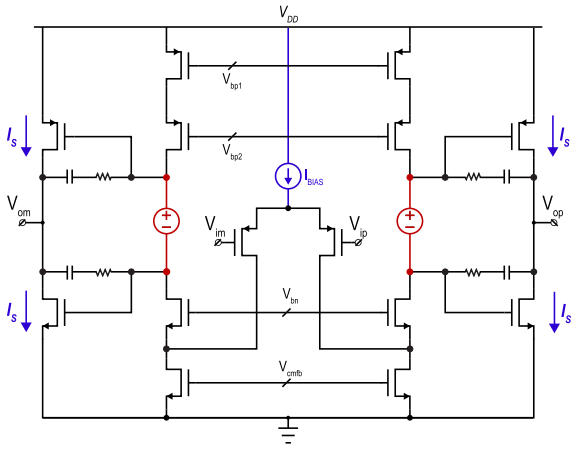

Fig. 6. Schematic of the two-stage op-amp used to implement the driver.

packaging stress, and aging-related effects. To estimate φ_{err} and compensate for the Φ_Q residual drift, an additional demodulator is used to monitor the Q signal.

V_{SENSE} is digitized with high resolution (>14 bits) by I and Q delta–sigma modulators ($\Delta\Sigma$Ms), I-$\Delta\Sigma$M and Q-$\Delta\Sigma$M, respectively. Their respective bitstream outputs, BS_I and BS_Q, are then decimated and combined with appropriate weights (k_0) in the digital domain to cancel any residual Q error [28]. The estimation of φ_{err} [$\approx 90° \pm \arctan(Q/I)$] is performed at zero flow in a temperature stable environment. The resulting $k_0 \approx \pm\varphi_{\text{err}}$ is then fixed and used for all subsequent measurements. It is important to note that any in-phase and quadrature parasitic couplings that appear at Φ_I and Φ_Q [29] may affect the accuracy of φ_{err} estimation, thus leading to residual errors in the cancellation of Φ_Q in (7).

B. Drive Loop Circuit Implementation

The simplified circuit implementation of the drive loop is shown in Fig. 5. Differential changes in $C_{D1,2}$ are converted into a voltage (V_{INT}) by a low noise capacitive transimpedance amplifier (C-TIA). Its gain is set by a feedback capacitor $C_1 = 1$ pF and varies from 154 to 158 dBΩ, as f_D varies from 1.8 kHz (liquids) to 3 kHz (gases). For the drive loop, the sensor can be simply seen as a resonator. Thus, as given by (7), the phase noise of the demodulating reference (f_D) should be minimized to avoid increasing the mass flow noise floor. For a given drive mode Q-factor (Q_D), this can be achieved by increasing V_{BD} and x_D and reducing the C-TIA's noise [35]. The former is set to 20 V and ~10 μm, respectively, to prevent the pull-in of the $C_{D1,2}$ combs and to ensure that the resonator operates in its linear region. With this, the C-TIA's input referred current noise is designed to be <12 fA/$\sqrt{\text{Hz}}$ over the f_D operating range.

To guarantee proper startup at f_D, the integrating C-TIA is followed by a differentiator during the startup phase (ϕ_{ST}). This adds an extra 90° phase shift to V_{INT} to ensure oscillation [36], while a driver with a gain R_F/R_{ST} provides the drive current i_D. Fig. 6 shows the two-stage op-amp used to implement the driver. It employs a class-AB output stage to efficiently provide the large currents (≈10 mA) needed to drive the 175-Ω metal track with rail-to-rail ($1.8V_p$) signals. A static current (I_S) of 1 mA is enough to provide the peak i_D.

The clock signals used in the drive loop are provided by a charge-pump (CP) PLL. This is driven by a comparator with hysteresis, which converts V_{INT} into a well-defined square wave. Once the PLL locks to f_D, which takes place within 25 ms, the startup path is disabled, and the drive loop is controlled by the PLL's output. The comparator's hysteresis is also reduced to avoid phase errors [30]. Compared with conventional sustaining circuits [14], this PLL-based frequency control ensures robustness to the sensor's parasitic resonant modes and simplifies the generation of the drive loop and sense readout clocks [31]. During normal operation, the output of the C-TIA is boosted by two capacitively coupled amplifiers (CCAs), which drive the PLL and the amplitude regulator, respectively. They block the output offset of the C-TIA and isolate it from the kickback of the PLL's comparator and the amplitude controller's switching transients.

DC feedback around the C-TIA and CCAs is provided by GΩ bias resistors (R_B) to ensure that the associated high-pass corner frequencies are at least two orders of magnitude lower than f_D (1.8–3 kHz). As a result, the associated phase shift has a negligible effect on drive loop operation. Although pseudo-resistors offer an area-efficient way to realize GΩ resistors, they require process/temperature compensation to ensure reliable operation [37]. In this work, switched resistors [38], [39] are used to boost 1-MΩ poly resistors to >5 GΩ. As shown in Fig. 5, the PLL provides the duty-cycled clock signal f_{PULSE} ($=16 \times f_D$) to realize R_B.

In the amplitude control path (see Fig. 5), a switched-capacitor (SC) network provides amplitude regulation by sampling the peak of V_{INT} synchronously [30], [32] and comparing it with a reference voltage (V_{REF}). The Q-factor of the sensor's drive mode (Q_D) is quite limited (~60–~100). Thus, a proportional–integral (PI) controller is used to achieve accurate control of the drive amplitude. To decouple the controller parameters, the proportional and

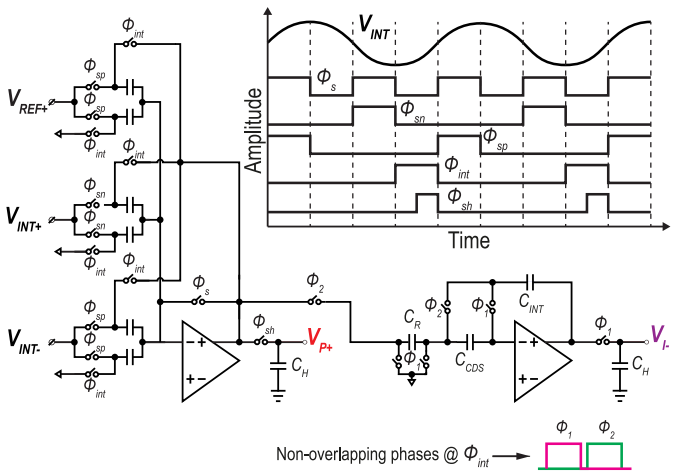

Fig. 7. Half circuit of the SC amplitude controller.

integral gains are set by resistors R_P and R_I, respectively. These convert the proportional (V_P) and integral (V_I) voltages into currents, which are then up-modulated to f_D and summed at the driver's virtual ground, thus providing the amplitude-controlled i_D at a steady state.

Since ΔC_S is modulated by f_D and is proportional to x_D, amplitude control signals with low $1/f$ noise are required for good long-term stability [14], [33], [40]. Fig. 7 shows the simplified half circuit of the SC amplitude controller and its timing diagram. It consists of a flip-around amplitude detector [34], which generates the controller's error voltage by detecting the amplitude of V_{INT} and comparing it with V_{REF}. It is followed by a correlated-double-sampling (CDS) SC integrator, which samples and integrates the error voltage. V_P and V_I are sampled and held by large C_H ($\approx$10 pF) capacitors, which are buffered before being fed to R_P and R_I. Since both positive and negative peaks of V_{INT} are compared, the offset of the preceding stages is cancelled. This leads to a low $1/f$ noise, stable, and offset-free control of the drive amplitude.

The sensor is used for liquids and gases, and therefore, the drive loop must be able to operate with a wide range of resonance frequencies while at the same time accommodating changes in Q_D. Due to the SC implementation, most of the controller's time constants are locked to f_D. Therefore, the drive loop remains stable with both liquids and gases.

C. Sense Front-End Implementation

Fig. 8(a) shows the circuit implementation of the sense C/V front-end. It consists of a low-noise C-TIA that converts ΔC_S into a voltage V_{SENSE}, which contains mass flow information at f_D. Around f_D, its input-referred current noise can be expressed as follows [36]:

$$i_{n,i}^2 = v_{n,\mathrm{OTA}}^2\left[(C_S + C_F + C_p + C_{\mathrm{GG}})s\right]^2 + i_{n,B}^2 \qquad (8)$$

where $v_{n,\mathrm{OTA}}^2$ is the input-referred voltage noise of the OTA used to implement the C-TIA, $i_{n,B}^2 = 4kT/R_B$ is the feedback resistor equivalent current noise, k is Boltzmann's constant, T is the temperature, and C_F, C_P, and C_{GG} are the feedback,

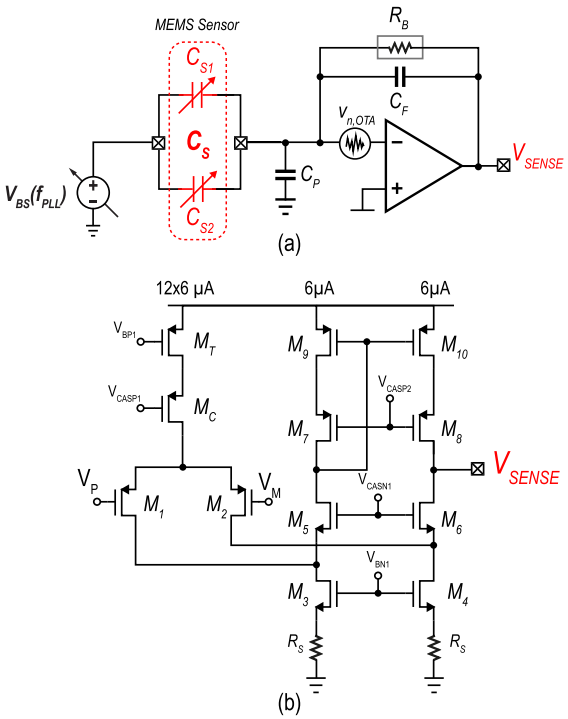

Fig. 8. (a) Sense readout capacitors C-TIA front end and (b) its OTA implementation.

the parasitic, and the OTA's input capacitances, respectively. As in the drive loop C-TIA, the large DC feedback resistor R_B ($>$5 GΩ) is realized by a switched resistor. As a result, the sense C-TIA input-referred noise is mainly determined by the OTA's voltage noise in combination with the total input capacitance.

The schematic of the OTA used in the sense C-TIA is shown in Fig. 8(b). When operated with H_2O, $f_D \approx 1.8$ kHz, and so, the OTA's $1/f$ corner frequency must be lower than this. A folded-cascode topology provides $>$85-dB DC gain, which is enough to establish a well-defined virtual ground for the biasing of the sense combs. To achieve low $1/f$ noise, the input pair (M_1–M_2) consist of large (1000/0.6 μm) transistors. They are sized, such that their input capacitance (C_{GG}) is less than 1/3 of the total parasitic capacitance ($C_P > 10$ pF) at the virtual ground, thus limiting their contribution to the total noise in (8). The OTA's current sources ($M_3 - M_4$) are degenerated by resistors R_S, which reduces their $1/f$ noise at the expense of thermal noise. To mitigate this, a 1/12 current ratio is used between the input-pair tail (M_T) and the current sources (M_9–M_{10}) in the folding branch, thus reducing their transconductance and, hence, their $1/f$ noise. This leads to a $1/f$ corner frequency of $\sim$1.5 kHz.

D. IQ-$\Delta\Sigma$Ms Implementation

Fig. 9 shows the simplified circuit diagram of the I-$\Delta\Sigma$M used in the sense readout channel. It is based on the second-order CT-$\Delta\Sigma$M. The sensed V_{SENSE} containing mass flow information at f_D (node I) is converted into a current by the first integrator's input resistor (R_{IN}). The resulting current is then demodulated to DC by a passive mixer (chopper)

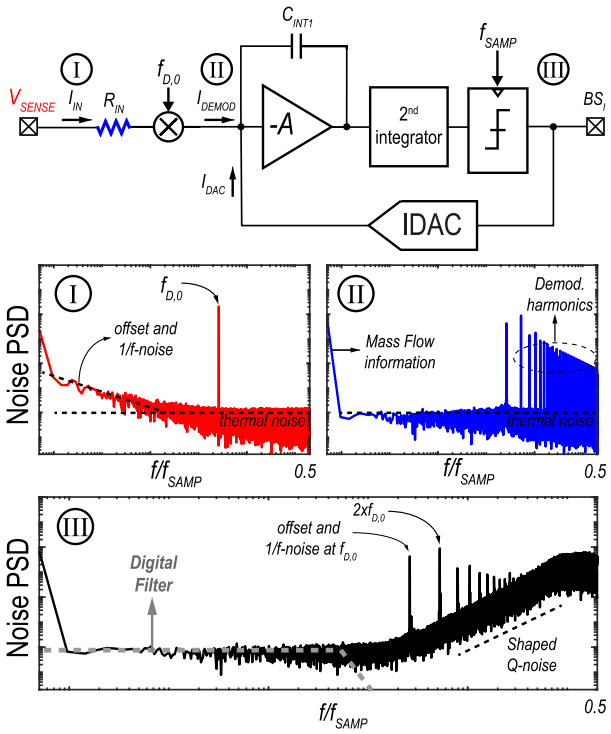

Fig. 9. Simplified circuit diagram of the sense readout I-$\Delta\Sigma$M and its noise PSD.

with a reference $f_{D,0}$, which is in-phase with the mass flow component in V_{SENSE}. Then, the downconverted current (I_{DEMOD}) is the input of $\Delta\Sigma$M (node II). As a result, the modulator's bitstream (BS_I) average is, thus, a measure of mass flow (Φ). The offset and $1/f$ noise components in V_{SENSE} will be modulated to the harmonics of $f_{D,0}$, where they are removed by the modulator's digital decimation filter (node III).

Since the symmetrical placement of C_S around the drive mode axis inherently attenuates the amount of Q error in V_{SENSE}, the extra harmonics caused by the residual Q error can also be removed by the decimation filter. As discussed in Section III-A, an additional Q-$\Delta\Sigma$M is used to compensate for the residual drift in BS_I due to Q as well as to estimate φ_{err}.

Fig. 10 shows the schematics of the sense readout used in the ΦDC. The sense C-TIA output is boosted by a CCA, which functions as a single-to-differential stage that drives both I- and Q-$\Delta\Sigma$Ms [see Fig. 10(a)]. The I/Q passive mixers are implemented by NMOS choppers placed between the $\Delta\Sigma$M first integrator's virtual ground and the input resistor (R_{IN}), as shown in Fig. 10(b) [41]. Thus, high linearity can be achieved, since the voltage swing across the chopper's nonlinear switches is only a fraction of the input signal swing. Since a square wave is used for demodulation, odd harmonics will be demodulated back to DC with the wanted signal, with the third-order harmonic being the most dominant among them. The third-harmonic distortion (HD_3) of the mixers can be approximated as follows [42]:

$$\mathrm{HD}_3 = \frac{3}{32}\left(\frac{V_{\mathrm{IN}}}{V_{\mathrm{GT}}}\right)^2 \cdot \left(\frac{r_{\mathrm{ON}}}{R_{\mathrm{IN}}}\right)^3 \tag{9}$$

where V_{GT} is the gate overdrive voltage, and V_{IN} is the maximum demodulator input. To keep $\mathrm{HD}_3 < -130$ dB, the switch resistance r_{ON} (≈1.5 kΩ) is much smaller than R_{IN} (180 kΩ).

The downconverted I/Q currents, and their harmonics, provided by the CCA are digitized by two second-order CT-$\Delta\Sigma$Ms using feedforward (FF) loop filters with finite impulse response (FIR) DAC [43], as shown in Fig. 10(a). The use of four-tap FIR-DAC reduces the swing of current error to be processed by the first integrator, thus improving the passive mixer's linearity. Furthermore, for the sampling frequency (f_{SAMP}) of 500 kHz, the use of a FIR-DAC allows the first integrator to be chopped at a fraction of f_{SAMP} (1/8), mitigating its offset and $1/f$ noise without incurring quantization noise folding [44]. The first stage amplifier is a two-stage miller-compensated op-amp based on current-reuse stages [see Fig. 10(c)]. Its energy-efficient design allows the two $\Delta\Sigma$Ms to consume about 200 μW, thus adding little to the overall system power consumption. The second integrator is implemented in an FF fashion, which also facilitates the use of a capacitive DAC (CDAC) to compensate for the delay added by the main FIR-DAC.

The resulting bitstreams (BS_I and BS_Q) are decimated off-chip and then recombined to realize the low-drift mass flow digital output of the ΦDC. Moreover, since the offset of the stage preceding the $\Delta\Sigma$Ms is up-modulated to f_D, and the first integrator is chopped, the ratio between BS_I and BS_Q averages (μ_Q/μ_I) at zero flow can be used as a measure of $\varphi_{\mathrm{err}} \approx 90° \pm \arctan(\mu_Q/\mu_I)$ with negligible error.

IV. Measurement Results

The integrated CMOS sensor interface [see Fig. 11(a)] was implemented in a standard 0.18-μm process and co-integrated with the SCT-fabricated MEMS Coriolis mass flow sensor in a chip-on-board assembly for prototype testing [see Fig. 11(b)]. The interface circuitry and the MEMS chip occupy 2.2 × 2.2 and 15 × 7.5 mm, respectively. For flexibility and testability, the required digital logic and HV bias generation were implemented off-chip. The complete system draws about 7.2 mA from a 1.8-V supply during steady-state operation.

A daughter printed-circuit board (PCB) containing the ΦDC is mounted over a small custom 3D printed fluid connector. To prevent leakage, O-rings are used below the PCB. At the same time, a metal clamp containing the magnets to generate $B \approx 50$ mT tightens the fluidic connection. This is possible, since the MEMS samples include glass caps that are integrated into the sensor at the wafer level.

The fluid connector and the setup used for mass flow measurements are shown in Fig. 11(c). Several fluids, from liquids to gases, are used for the characterization of the proposed system: nitrogen (N_2), argon (Ar), carbon dioxide (CO_2), water (H_2O), isopropanol (IPA), and mixtures of H_2O + IPA. Pressure regulated (6 bar) tanks are used for gases and pressurized (2 bar) Falcon tubes for liquids. Bronkhorst mini Cori-FLOW mass flow controller and sensor are used at the flow input and output, respectively, and serve as a reference [21]. This way, the mass flow through the device under test can be guaranteed with negligible leakage. The daughter PCB is connected below a motherboard, and a flat cable is used to communicate with

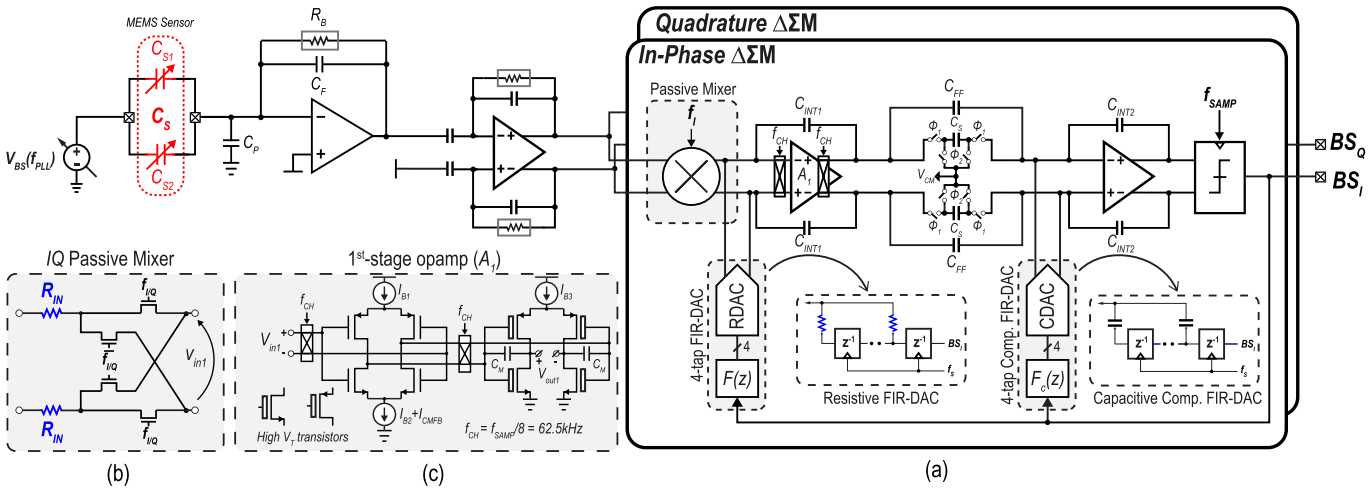

Fig. 10. Detailed schematic of the (a) proposed ΦDC, (b) passive mixer, and (c) first integrator opamp implementation.

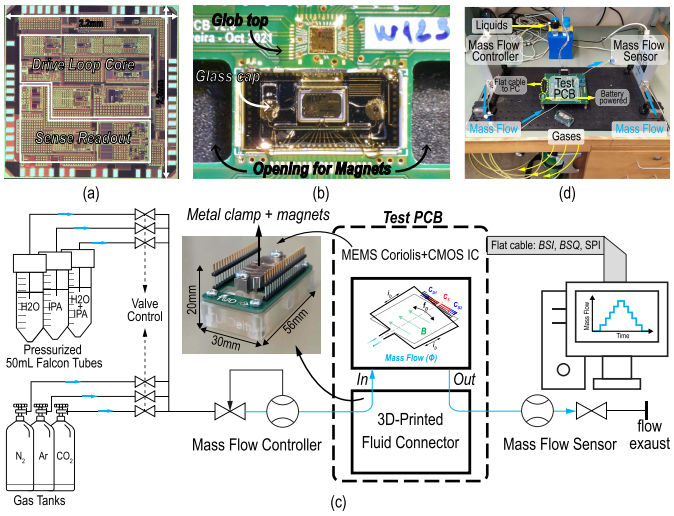

Fig. 11. (a) Micrograph of the fabricated 0.18-μm CMOS sensor interface. (b) Chip-on-board co-integration with the SCT-fabricated MEMS Coriolis mass flow sensor. (c) Mass flow measurement setup illustration, 3D printed fluid connector close-up, and (d) photograph of the actual setup.

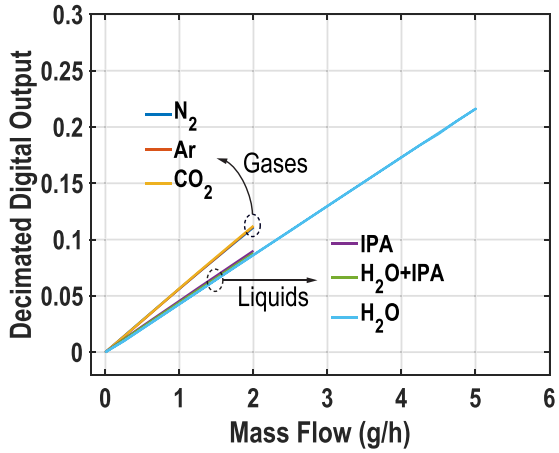

Fig. 12. Measured ΦDC decimated output versus mass flow for liquids and gases.

a PC to acquire the ADCs bitstreams and define the system's setting via SPI. The complete setup is shown in Fig. 11(d).

Fig. 12 shows the ΦDC decimated output for mass flows of liquids and gases up to 5 and 2 g/h, respectively, with 250 mg/h, for a fixed $V_{\mathrm{BS}} = 20$ V. The outputs of both increasing and decreasing mass flow are overlapped, thus showing that the sensor has no hysteresis. Since the ΔΣMs have to process the mass flow signal at DC, its harmonics, and the up-modulated front-end offset, the flow range is limited to 5 g/h to avoid saturating the modulators.

The sensor's mass flow sensitivity has a linear relationship with f_D, as shown by the red dashed line in Fig. 13(a), which is expected from (3). To achieve a single response from liquids to gases, a resonance frequency-dependent gain calibration can be performed in the digital domain [14]. However, this leads to fluid-dependent resolution, since both noise and signal are scaled simultaneously. Here, a relatively straightforward background sensitivity tuning (BST) scheme performs a frequency-dependent gain calibration in the analog domain. This calibration scheme is described as follows.

1) Sensitivities at N_2 and H_2O are matched: Having the highest sensitivity as a reference (N_2 with $V_{\mathrm{BS}} = 20$ V), V_{BS} for H_2O is increased until its sensitivity matches that of N_2 at 2 g/h [see Fig. 13(a)].
2) Relationship Between f_D and V_{BS}: From step 1), a linear fit between N_2 and H_2O is obtained to map the required f_D to V_{BS} relationship to match their sensitivities [see Fig. 13(a) (black solid line)].
3) Based on 1) and 2), V_{BS}, when operating with all fluids between N_2 and H_2O, is tuned continuously in the background proportionally to f_D.

The normalized sensitivities with and without the BST scheme are shown in Fig. 13(a). The ΦDC mass flow response after the calibration is performed, and V_{BS} is tuned externally in the background, is shown in Fig. 13(b), with a between fluids sensitivity error of <2.6%. Although most of the sensitivity drop is compensated by the increase in V_{BS}, the relatively large residual error indicates that changes in (ω_D/ω_S) due to pressure drop or compressibility [18], and variations in drive amplitude and capacitive readout nonlinearity related

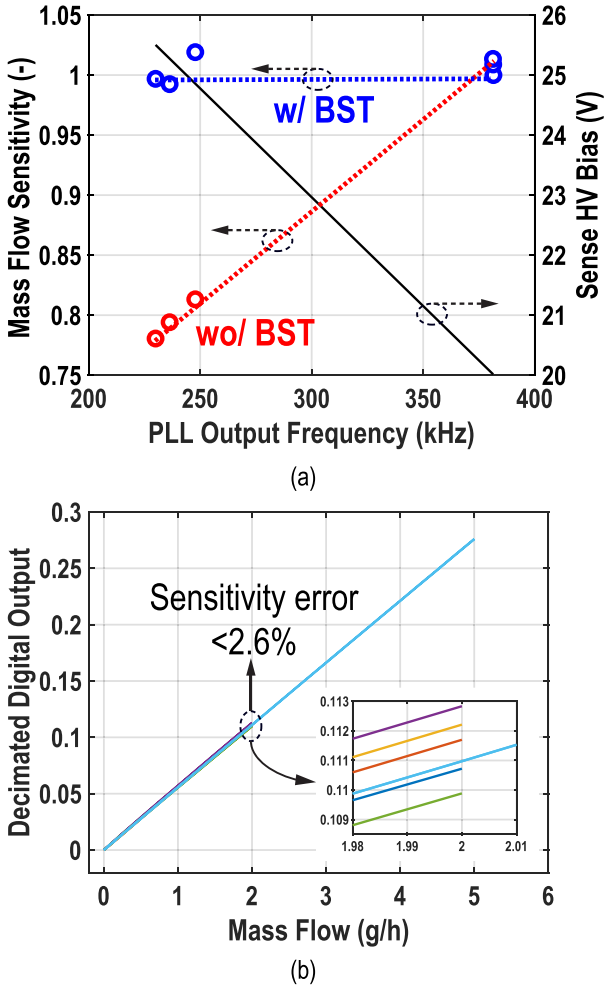

Fig. 13. (a) Normalized mass flow sensitivity versus resonance frequency. (b) Mass flow response after BST.

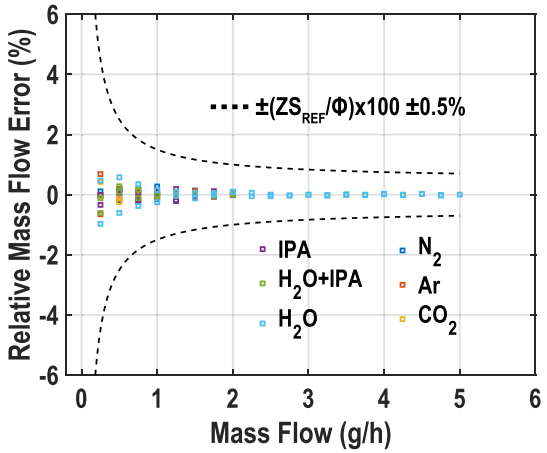

Fig. 14. Relative mass flow error for all tested fluids after the fifth-order polynomial fit.

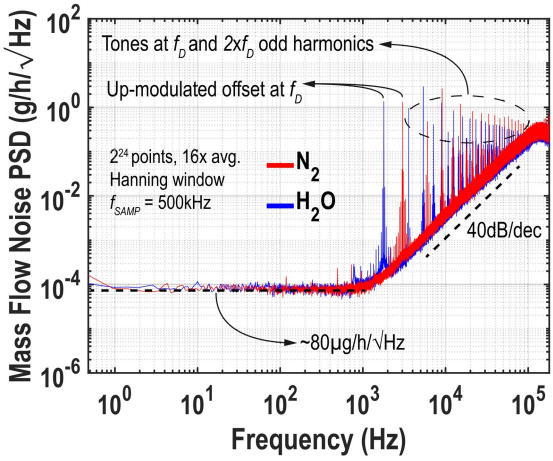

Fig. 15. Mass flow noise power spectral densities for N_2 and H_2O.

effects, which are not proportional to f_D, might also have to be compensated to obtain lower sensitivity errors. Moreover, the expected flow error at 2 g/h for the used reference flow sensor [21] is expected to be about ±1%.

Coriolis sensors are often required to achieve mass flow errors in the order of ±0.5% [24]. To achieve this, up to ten mass flow points are often used to fit the sensor's response. The ΦDC relative mass flow error is shown in Fig. 14. Its output is fit by the fifth-order polynomial and compared with the actual mass flow. Its errors are well within the "trumpet curve" boundaries, representing the ZS of the reference sensor (ZS_{REF}) plus its ±0.5% calibrated error. The error is relatively similar for liquids and gases and stays below ±1% at 250 mg/h.

The bitstream spectra (∼33.5-s interval, Hanning window, 16× averaging) of ΦDC output under zero flow condition for when the sensor is filled with N_2 and H_2O are shown in Fig. 15. It shows a noise floor of about 80 μg/h/$\sqrt{\text{Hz}}$ after compensating for the drop in sensitivity when operating with H_2O. The up-modulated offset and the demodulation harmonics are filtered by a sinc^2 decimation filter. With the parameters in Table I, this indicates that the equivalent capacitance resolution is below 24 zF/$\sqrt{\text{Hz}}$, which is mainly limited by thermal noise of the sense C/V converter.

To obtain the ZS of the sensor, its decimated output in a 4-Hz bandwidth was observed at zero flow, when filled with H_2O, and room temperature. As shown in Fig. 16, it achieves a ±0.31 mg/h standard deviation for a measurement time of about 4 h and is relatively constant with both N_2 and H_2O [15].

The combined phase error (φ_{err}) due to the sense mode limited Q-factor and C/V front-end causes the output to drift, since Q leaks into the wanted mass flow signal (I). Fig. 17 shows the Allan deviation for the data in Fig. 16. Its minima (∼8 μg/h), which represents the lowest resolution achievable by averaging, is limited by $1/f$ noise. The minima is improved by about 1.6× when the compensation is used. This indicates that Q, and thus its induced drift, is already substantially reduced by design, since the sense combs experience much smaller displacements than the actual x_D. Moreover, the minima is 19× smaller than that obtained with conventional phase readout [14], thus showing how the use of a single C/V converter to readout $C_S = C_{S1} + C_{S2}$ minimizes the drift due to the large Q signal while at the same time reducing the noise floor by more than 3 dB.

TABLE II
PERFORMANCE SUMMARY AND COMPARISON

	[8]	[9]	[10]	[21]	[12]	[13]	[14]	This Work
Sensor Type	MEMS Thermal	MEMS Thermal	MEMS Thermal	SS Coriolis	MEMS Coriolis	MEMS Coriolis	MEMS Coriolis	**MEMS Coriolis**
Readout Technology	CMOS	0.18µm CMOS	COTS	COTS	COTS	COTS	0.18µm CMOS	**0.18µm CMOS**
ASIC Output	Analog	Analog	Analog	-	-	Analog	Analog	**Digital**
Fluids	H_2O	N_2	N_2	Liquids & Gases	Liquids (H_2O/IPA)	N_2	Liquids & Gases	**Liquids & Gases**
Full Scale (g/h)	0.09	0.36[c]	0.09[c]	5	500	0.75	5	**5**
Noise Floor (µg/h/√Hz)	<600[a]	<200	<100	<4000[a]	N/A	NA	300	**80**
Bandwidth (Hz)	4	25	25	3	1.6	0.63	3	**4**
Dynamic Range[b] (dB)	43.5	65.1	59	61.9	N/A	N/A	84.4	**95.9**
Zero Stability (mg/h)	N/A	N/A	N/A	±10	±223	2.6	±0.8	**±0.31**
Power Consumption (mW)	21.5	8.9	13.1	2500	400	1250	14.6	**13**

[a]Estimated from repeatability. [b]1Hz bandwidth. [c]Limited to the linear range. N/A = Specification not available

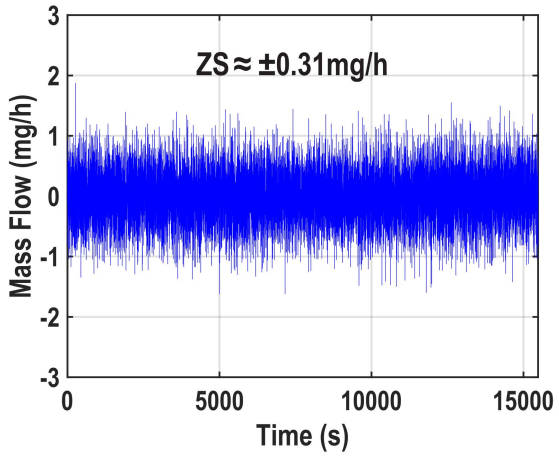

Fig. 16. Measured zero stability when the sensor is filled with H_2O.

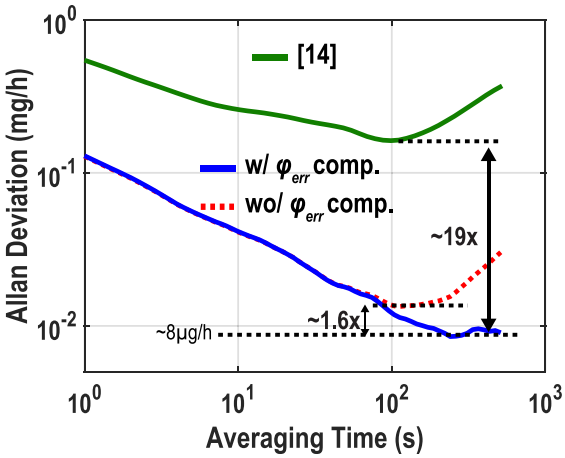

Fig. 17. Allan deviation of the mass flow output with and without phase-error compensation and its comparison with [14].

An advantage of MEMS Coriolis mass flow sensor over other sensing principles is its ability to estimate fluid density [23]. To measure this, the sensor is filled with different

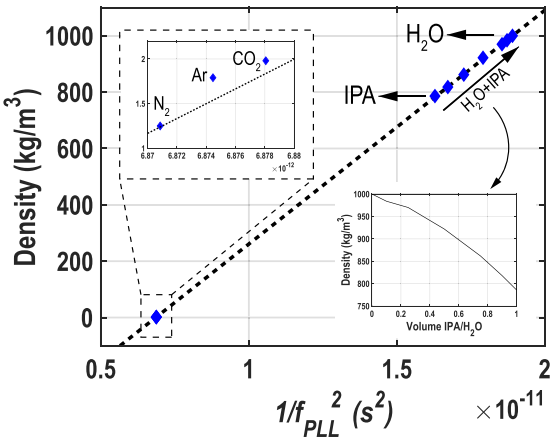

Fig. 18. Measured density to resonance frequency relationship.

fluids, and its resonance frequency change is observed with a frequency counter (Keysight 53230A). The relationship between fluid density and resonance frequency is measured under ambient pressure and zero flow. The sensor's measured density to f_D response is shown in Fig. 18. To cover the wide density range between H_2O and IPA, the response for various mixtures of the two with volume ratios of 10/90, 25/75, 50/50, 75/25, and 90/10 is also included. As shown by the first-order fit dashed line, the response is a linear function of $f_D{}^{-2}$, where $f_D = f_{\mathrm{PLL}}/128$, thus matching the expected response from (1). This relationship is more closely followed by liquids, since with gases, the added fluid mass (m_f) is smaller than that of the suspended channel (m_{ch}).

Table II summarizes the performance of the proposed ΦDC and compares it with state-of-the-art MEMS Coriolis and MEMS thermal flow sensors. Compared with the lowest flow range miniaturized SS Coriolis sensor in the market [21], it achieves a >30× improvement in resolution and ZS while dissipating >190× less power.

Against MEMS Coriolis's prior works [13], [14], the proposed system provides a digital output, and its noise floor

and ZS are 3.75× and 2.6× lower, respectively, while remaining the same for all tested fluids. Furthermore, its performance is at par with state-of-the-art MEMS thermal flow sensors and, in some cases, even outperforms them [8].

V. Conclusion

A ΦDC capable of sensing low flow rates of liquids and gases, as well as their densities, has been presented. The system consists of a micromachined Coriolis flow sensor and a 0.18-μm CMOS readout interface, which drives the sensor into resonance and digitizes the mass flow information. With the help of dynamic offset cancellation techniques applied to both drive and sense readout circuitry, low noise floor and low offset drift could be achieved. Similar performance for both liquids and gases is achieved using a resonance frequency-dependent sensitivity tuning scheme. Furthermore, the phase error in the sense readout path is cancelled in the digital domain, reducing quadrature-induced drift and improving its long-term stability. These advances lead to a MEMS Coriolis digital-output system with a state-of-the-art noise floor of about 80 μg/h/$\sqrt{\text{Hz}}$ and a ZS of $\pm$0.31 mg/h. The proposed sensor is a promising alternative to MEMS thermal flow sensors in microfluidic and low-flow industrial applications, such as high-pressure liquid chromatography, drug infusion, and etching machines.

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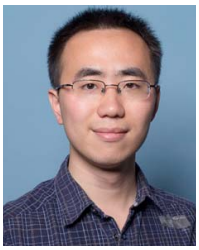

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