

## Energy-Efficient Smart Temperature Sensors in CMOS Technology

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# **Energy-Efficient Smart Temperature Sensors in CMOS Technology**

Kamran Souri



# Energy-Efficient Smart Temperature Sensors in CMOS Technology

Proefschrift

ter verkrijging van de graad van doctor  
aan de Technische Universiteit Delft,  
op gezag van de Rector Magnificus prof. ir. K.C.A.M. Luyben,  
voorzitter van het College voor Promoties,  
in het openbaar te verdedigen

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*To my lovely wife, Sara*



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# Chapter 1

## Introduction

Temperature is the most often-measured environmental quantity [1.1]. This is because nearly all physical, chemical, mechanical and biological systems exhibit some sort of temperature dependence. Temperature measurement and control are therefore critical tasks in many applications. Traditionally, temperature sensors have been implemented with discrete components such as resistance temperature detectors (RTDs), thermistors, or thermocouples. In the last three decades, integrated temperature sensors, particularly in CMOS technology, have become a promising alternative. A sustained research effort has been devoted to the development of compact, low-cost temperature sensors with co-integrated readout circuitry, thus providing temperature information in a digital format. Such *smart* temperature sensors (see Figure 1.1), are conventional products nowadays [1.3 - 1.7].

There are several advantages associated with smart sensors; firstly, since a digital output is almost mandatory in modern systems, no external analog-to-digital converter (ADC) is required. This higher level of integration reduces component count, and therefore size and, typically, cost. Secondly, in contrast to digital signals, analog signals are prone to interference and thus are not well suited for accurately transmitting data to other blocks in a system. Lastly, by integrating the readout circuit and the sensor on the same chip, on-chip digital post-processing becomes possible, which usually results in simpler systems.

### 1.1 Motivation

Smart temperature sensors have been around for many years. However, with the recent development of low-power radio systems, wireless temperature sensing has become very

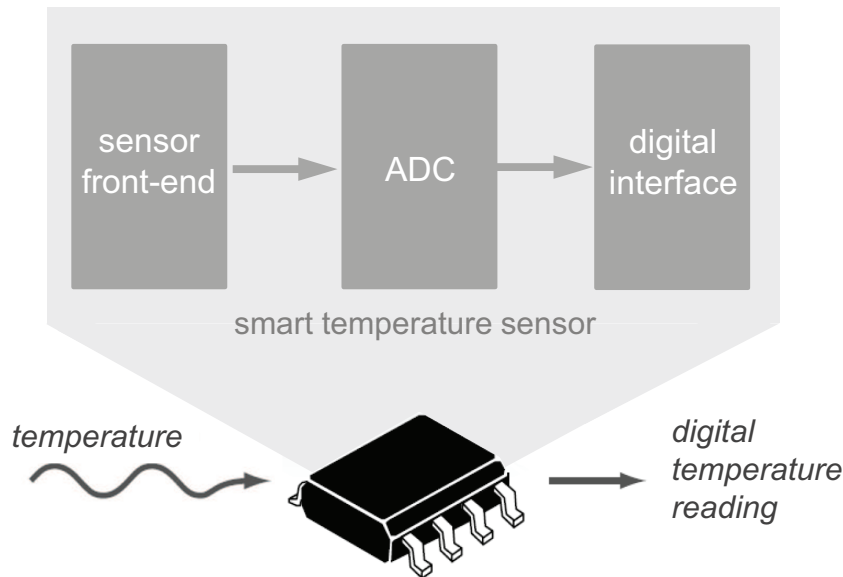


Figure 1.1: Block diagram of an integrated smart temperature sensor [1.2].

attractive, as it opens up a wide variety of new applications. One can think of applications in cold supply chains, monitoring of perishable goods, animal husbandry and agriculture, automotive, building automation, and healthcare. Wireless sensor networks (WSNs), which consist of spatially distributed sensor nodes with a wireless communication infrastructure, were introduced in the 2000s [1.8]. Various physical or environmental quantities such as temperature, sound, humidity, motion, pressure, etc. can be sensed and digitized by the sensor nodes. The digitized signals are then passed through the communication network towards a centralized or distributed control unit for further processing, as shown in Figure 1.2. As the name WSN suggests, and mainly due to cost reasons and ease of integration, wireless operation is a key feature, which at the same time makes powering the sensor nodes a challenging task. Most WSNs have used battery-powered sensors nodes, while quite recently, nodes based on energy harvesting or scavenging have also been introduced [1.9].

Another opportunity for wireless sensing has recently emerged through the introduction of radio frequency identification (RFID) technology as a versatile wireless communication platform. RFID has been around for years now, and has become a billion dollar market over the last few years and it is still growing. With an estimated \$5.6 billion market in 2010, and an average 15% year-on-year growth rate (see Figure 1.3), the forecasted market in 2020 will exceed \$21.9 billion [1.10]. This shows that RFID technology has achieved solid penetration throughout worldwide commerce, boosted by dynamic growth in the retail apparel sector. The freedom provided by small size and easy positioning, non-line-of-sight wireless operation and powering, and extended read ranges are key features

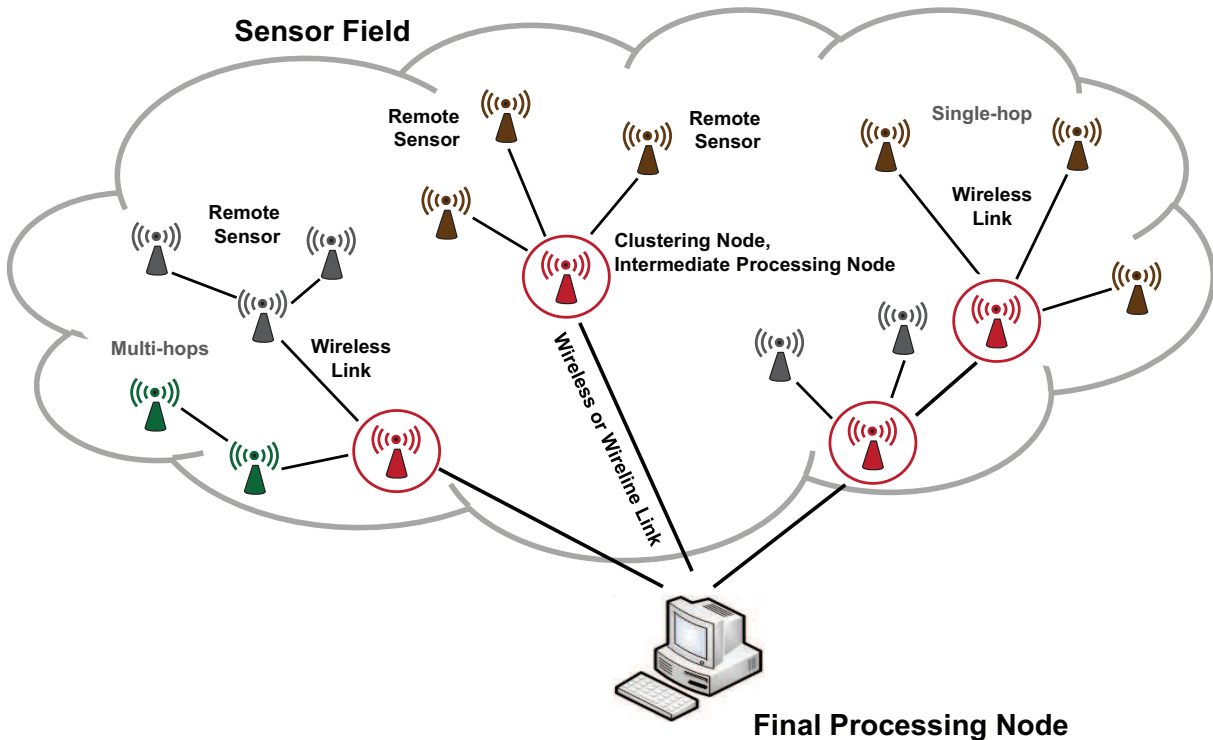


Figure 1.2: A typical wireless sensor network (WSN) arrangement [1.8].

that have made RFID technology so promising.

Apart from its primary application in identification and tracking, RFID has become a pragmatic building block for the internet of things (IoT), thus creating a flood of new applications in numerous industries [1.11]. According to an IC Market Drivers report in 2016 [1.12], 30.0 billion Internet connections are expected to be in place worldwide in 2020, with 85% of them being to web-enabled *things*, meaning a wide range of commercial, industrial, and consumer systems, distributed sensors, vehicles, and other connected objects. As reported, IoT applications will fuel strong sales growth in optoelectronics, sensors/actuators, and discrete semiconductors, which are projected to rise by a compound annual growth rate (CAGR) of 26.0% between 2015 and 2019, thus offering a forecasted market of \$11.6 billion in 2019.

Most RFID tags consist of two main parts (see Figure 1.4). The first part is an integrated circuit (IC) to implement the target functionality, e.g. the storing and processing of information, as well as the RF transceiver. This part usually occupies only a small portion of the total area of the tag. The second part, which takes up the bulk of the area, is the antenna, which is required for receiving and transmitting the RF signal. Depending on their source of energy, RFID tags can be classified into *passive* and *active* tags. Active RFID tags include a battery to power the IC, which makes autonomous operation possible.

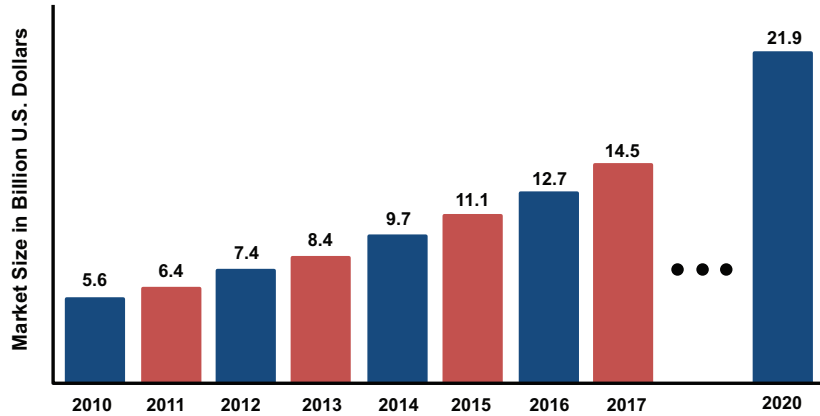


Figure 1.3: Projected size of the global market for RFID tags from 2010 to 2020 (in billion U.S. dollars) [1.10].

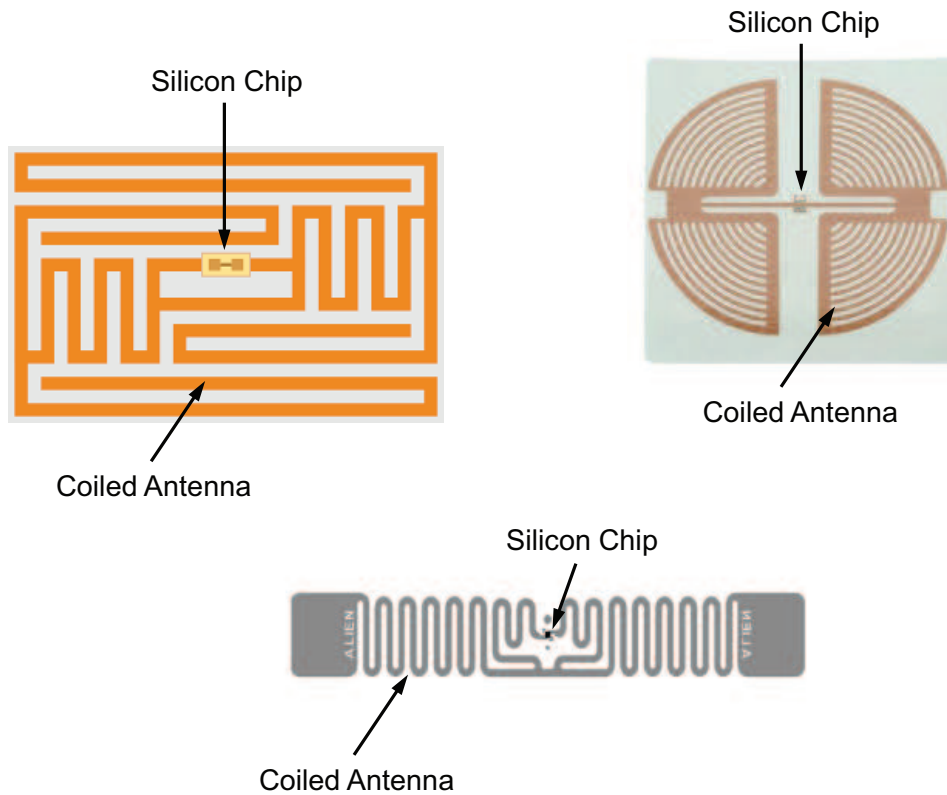


Figure 1.4: Various samples of RFID tags; each tag is composed of a large antenna and a silicon integrated circuit (IC).

In consequence, low power designs, along with brief operating periods, are desirable in order to maximize battery lifetime. Passive RFID tags, in contrast, are not equipped with a battery and consequently, autonomous operation is not possible. Instead, the power required to operate the tag is scavenged from an external magnetic/electromagnetic field, transmitted by a reader. The energy absorbed via an antenna from the field is used to power the tag, thus, enabling data transmission and other functionalities. In other words, the antenna of a passive RFID tag is used to transfer information as well as to receive power.

The choice of RFID tag type depends on the target application. Battery-equipped or active RFID tags can communicate over long distances, up to 100 meters or more. Furthermore, they can operate continuously. However, they have limited lifetime (typically 1-4 years), significantly higher production costs e.g. few dollars and larger package size, all due to the use of a dedicated battery. The major advantage of passive RFID tags is that they can operate without a battery, thus offering much lower production cost (usually a few pennies), longer lifetime (20 years or more) and much smaller package size. For many years, the main drawback of passive tags was known to be their limited operating range, e.g. 3 - 5 meters. Recent tags with operating range up to 100 meters have been developed [1.13], thus making them the tags of choice for most RFID applications.

## 1.2 Challenges in Wireless Sensing

Although wireless temperature sensors seem very promising, there are many challenges associated with their implementation. To be cost-effective, such sensors must be fully compatible with CMOS technology. Fortunately, various temperature sensing elements are available in standard CMOS technology. However, due to the process spread of various elements, CMOS sensors often require sophisticated and/or time-consuming *calibration* and *trimming* processes (e.g. two-temperature calibration and trimming) to obtain sufficient accuracy. The calibration process is usually performed by comparing the sensor's output with that of a reference sensor at a number of known temperatures. Since both sensors need to reach thermal equilibrium, such thermal calibration can take several tens of seconds. The extra time required to perform calibration and trimming, however, increases the production cost, and thus sensors with no calibration or a minimum number of calibration points are desired. Alternatively, calibration techniques based on electrical measurements, can be developed to simultaneously achieve low-cost and good accuracy [1.14].

The required accuracy of a temperature-sensing node depends on the target application, ranging from  $\pm 0.1^\circ\text{C}$  for medical [1.15], [1.16] to  $\pm 1^\circ\text{C}$  for food and environmental



monitoring applications [1.17]. The operating temperature range also depends on the target application, e.g. from 35°C to 45°C in medical applications, from -40°C to 85°C in environmental monitoring, and from -40°C to 150°C in automotive applications. The actual number of required calibration points then depends on the type of sensing element, the target accuracy, and the sensor's operating temperature range. Clearly, there is a trade-off between the number of required calibration points (and therefore cost) and the target accuracy for a given application.

Furthermore, in the design of temperature-sensing wireless nodes, the power- and energy-efficiency of the co-integrated temperature sensor are key parameters. Typical CMOS smart sensors suffer from relatively high power consumption, e.g. 500 $\mu$ A in [1.3] and 2.2mA in [1.5], and/or long conversion time ( $T_{conv}$ ), e.g. 300 ms in [1.3] and 1.5s in [1.4], which results in high "energy consumption". Such sensors are ill-suited for use in battery-powered WSNs or active RFID tags as they would dramatically decrease the battery's lifetime, and thus are not cost-effective. They are also not suitable for use in passive RFID tags or WSNs operating based on energy harvesting or scavenging. This is due to the restricted amount of energy available in such systems, which either limits the maximum communication range or requires a larger antenna or energy storing element e.g. a capacitor, or calls for using energy harvesters. Moreover, the power received at a passive RFID tag falls off as the square of the distance. Therefore, there is a trade-off between the sensor's energy consumption on the one hand, and the operating range, size and cost of the sensor node on the other hand. This implies that energy-efficient sensors i.e. low-power (e.g. a few  $\mu$ W) sensors with fast conversion times are essential for wireless temperature sensing applications.

Temperature sensors for wireless sensing were introduced prior to the start of this research [1.17] - [1.18]. The design in [1.17] presents a temperature sensor, which is embedded into a passive-RFID tag. The tag dissipates 10 $\mu$ A to operate, and requires a conversion time of 510ms. It achieves an inaccuracy of  $\pm 2.5^\circ\text{C}$  (4 samples) from 0°C to 100°C, after a one-point calibration. The read range is limited to 10cm - 25cm, depending on the size of the antenna used. The sensor in [1.18] is quite power-efficient, dissipating 220nW from a 1V supply. However, it requires a conversion time of 100ms to obtain a resolution of 0.1°C. Furthermore, it requires a two-point calibration to achieve an inaccuracy of -1.6°C/+3°C (5 samples) from 0°C to 100°C. In 2010, a sensor was presented which dissipates 100nW, and achieves a resolution of 35mK in a conversion time of 100ms [1.16]. It also achieves an inaccuracy of  $\pm 0.1^\circ\text{C}$  (3 samples), over a range from 35°C to 45°C, but only after a two-point calibration. Recently, another temperature sensor embedded into a passive RFID tag has been presented [1.19]. The sensor dissipates 350nA from a 1V supply. After a one-point calibration, it achieves an inaccuracy of  $\pm 1.5^\circ\text{C}$  ( $3\sigma$ ) from -30°C to 60°C. In a conversion time of 14.5ms, it obtains a resolution of 0.3°C. As can

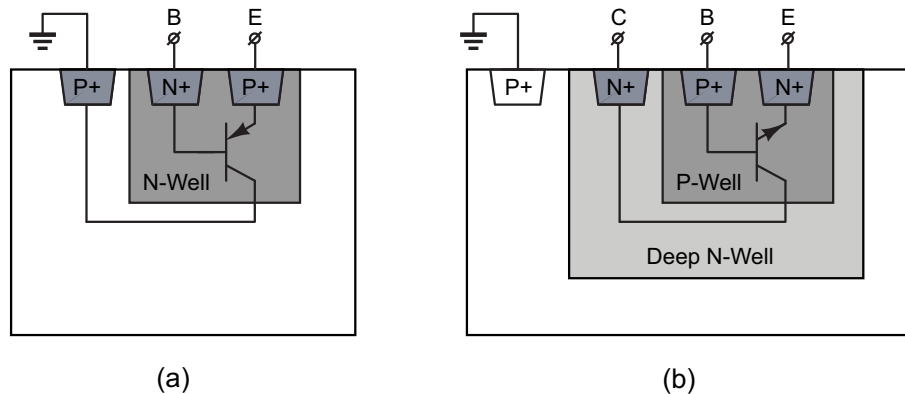


Figure 1.5: (a) Cross section of vertical PNP transistors in standard CMOS; (b) cross section of vertical NPN transistors in modern CMOS technology supporting deep N-well.

be seen, most of these low power/energy sensors suffer from poor accuracy, even after calibration.

In this thesis, we will focus on the design of low-cost, accurate, and energy-efficient CMOS temperature sensors. To understand the existing design trade-offs, we will first review various CMOS-compatible sensing elements from the perspectives of accuracy and energy-efficiency, which will be presented in the following section. A general figure-of-merit (FoM) will then be presented, which will facilitate comparisons between the energy-efficiency of different types of sensors. Lastly, a short survey of the state-of-the-art in 2009 will be provided, which enables us to evaluate the state of the art at the start of this research.

### 1.3 CMOS-Compatible Sensing Elements

In CMOS technology, the temperature dependence of several different circuit elements can be used for temperature sensing. The correct choice of sensing element, however, is not trivial and depends on the requirements of the target application, such as: accuracy, resolution, power consumption, conversion time, operating supply voltage range, operating temperature range and power supply rejection ratio (PSRR), etc. In the following, various CMOS-compatible sensing elements are briefly introduced and then investigated based on some of the aforementioned requirements.

### 1.3.1 Bipolar Junction Transistors (BJTs)

In CMOS technology, the same diffusions normally used to realize MOSFETs can be used to realize *parasitic* vertical bipolar junction transistors (BJTs). While smart temperature sensors based on lateral PNP transistors have been realized [1.20], [1.21], nowadays vertical PNP transistors are preferred due to their lower sensitivity to process spread and packaging stress [1.22], [1.23]. Such parasitic vertical PNPs, however, usually offer limited implementation flexibility, collector is formed inside the P-substrate, and thus, is not directly accessible (see Figure 1.5.a). In modern CMOS technologies with twin well or deep N-Well options, vertical NPN transistors are also available as shown in Figure 1.5.b. They exhibit significantly larger current gain than PNPs, e.g.  $\beta_F = 24$  (NPNs) versus  $\beta_F = 4$  (PNPs) in a TSMC 0.18 $\mu\text{m}$  CMOS technology. NPNs also offer more circuit design flexibility, since their collector terminals are accessible.

The base-emitter voltage  $V_{BE}$  of a BJT can be expressed as follows:

$$V_{BE} \approx \frac{kT}{q} \ln \left( \frac{I_C}{I_S} + 1 \right), \quad (1.1)$$

where  $k$ ,  $T$  and  $q$  denote the Boltzmann constant ( $1.38 \times 10^{-23}$  J/K), the temperature in Kelvin and the electron charge ( $1.6 \times 10^{-19}$  C) respectively. The parameter  $I_S$  denotes the saturation current of the bipolar transistor. It can be shown that  $V_{BE}$  exhibits complementary-to-absolute temperature (CTAT) behavior with a slope of  $\approx -2\text{mV}/^\circ\text{C}$  [1.2]. However, if two BJTs are biased at different collector current densities with a ratio  $p$ , the difference  $\Delta V_{BE} = V_{BE2} - V_{BE1}$  will be a proportional-to-absolute temperature (PTAT) voltage with a temperature coefficient that depends on the constants  $k/q$  and the ratio  $p$  [1.2]. The well-defined temperature dependency of  $V_{BE}$  and  $\Delta V_{BE}$  makes BJTs attractive for use in CMOS temperature sensors and bandgap voltage references. In fact, BJT-based temperature sensors have been widely used in the industry for decades [1.3] - [1.7]. The reasons for this are as follows: for a properly designed sensor, the dominant source of inaccuracy is the process spread in  $V_{BE}$ , which has been shown to have a PTAT profile [1.2], and thus, can be corrected by means of a cost-effective one-point PTAT trim, e.g.  $\pm 0.5^\circ\text{C}$  ( $3\sigma$ ) from  $-50^\circ\text{C}$  to  $120^\circ\text{C}$  in [1.24] and  $\pm 0.1^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$  in [1.25]. Another advantage is that the necessary temperature dependent and reference voltages are both generated from the same circuit, which significantly simplifies the implementation. They require bias currents in the range of  $\mu\text{A}$  or even sub- $\mu\text{A}$  to operate, and exhibit low supply dependency, usually a few tenths of degrees Celsius per Volt, e.g.  $0.5^\circ\text{C}/\text{V}$  in [1.24] and  $0.1^\circ\text{C}/\text{V}$  in [1.25].

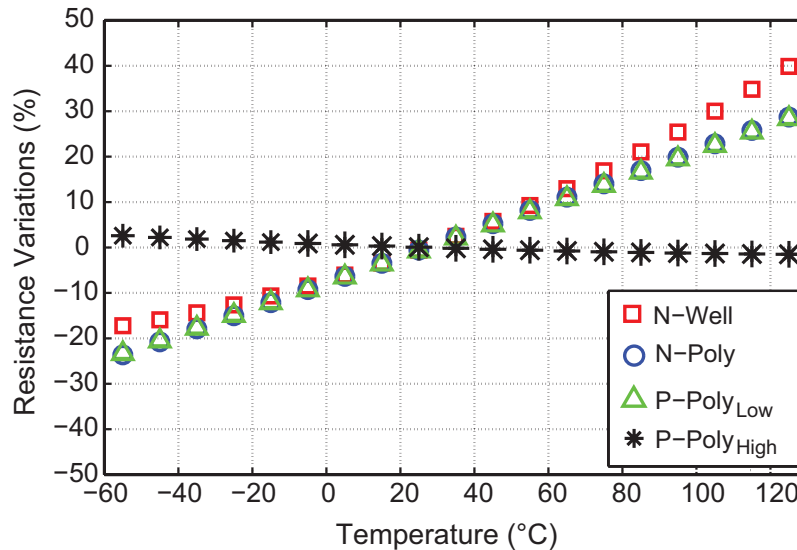


Figure 1.6: Temperature dependency of some types of resistors available in a TSMC  $0.18\mu\text{m}$  process. Resistance variations are normalized to the value at  $25^\circ\text{C}$ .

### 1.3.2 Resistors

Resistance temperature detectors (RTDs) have been widely used as stand-alone temperature sensing elements. Temperature information is obtained by reading out resistance variations as a function of temperature, implying that a large temperature-coefficient is often desired. As it turns out, most CMOS-compatible resistors exhibit significant temperature coefficients, with 1<sup>st</sup>-order coefficients ranging between  $0.1\%/^\circ\text{C}$  and  $0.4\%/^\circ\text{C}$ , depending on the resistor type. Figure 1.6 shows the simulated temperature-dependency of some of the resistors available in the TSMC  $0.18\mu\text{m}$  CMOS process. The variations are normalized to the resistance at  $25^\circ\text{C}$ . The temperature coefficient of  $+0.4\%/^\circ\text{C}$  exhibited by a typical N-Well or N-Poly resistor, means that its resistance will increase by about 72% over the temperature range from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ , which is reasonably large sensitivity. In such resistor-based sensors, the minimum supply voltage is usually limited by the readout circuit, thus enabling low supply voltages. The value of the bias current is defined by thermal-noise and area constraints.

A drawback of resistors as temperature sensing elements is the fact that the spread of most resistances in CMOS is in the range of 15 - 20% across the process corners. Their temperature coefficients also suffer from process spread and higher order non-linear terms, as can be noticed from Figure 1.6. As a result, resistors usually require a costly multiple-temperature calibration to achieve decent accuracy, where the number of calibration points could range between 3 to 5, depending on the target accuracy. The work presented in [1.26] and [1.27], for example, both achieve an inaccuracy of  $\pm 0.15^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$

to 85°C, but only after a costly three-temperature trim. Employing a single-temperature trim, the work in [1.28] achieves an inaccuracy of  $\pm 1^\circ\text{C}$  ( $3\sigma$ ) from  $-45^\circ\text{C}$  to  $125^\circ\text{C}$ , which is among the best reported for similarly-trimmed resistor-based sensors.

### 1.3.3 Electro-Thermal Filters (ETFs)

The thermal diffusivity of silicon  $D$  is defined as the rate at which heat diffuses through a silicon substrate. Recent research has shown that  $D$  is a well-defined parameter, as the silicon used for IC fabrication is highly pure [1.29]. Furthermore,  $D$  is strongly temperature dependent, and can be approximated by a power law:  $D \propto 1/T^{1.8}$  [1.30] - [1.32]. This well-defined temperature dependency can thus be exploited to realize temperature sensors. Figure 1.7 shows the structure of an electro-thermal filter (ETF), which uses a heater to generate heat pulses, and a (relative) temperature sensor (thermopile), fabricated at a distance  $s$  from the heater, which converts the received temperature variations into a small voltage signal. In the thermal domain, an ETF behaves like a low-pass filter. Driving such a filter at a given excitation frequency results in a temperature-dependent phase-shift [1.32], [1.33]:

$$\phi_{ETF} \propto (s\sqrt{f_{ref}})T^{n/2}, \quad (1.2)$$

where  $n \approx 1.8$ . A phase-domain ADC can then be used to digitize  $\phi_{ETF}$  and obtain temperature in digital format [1.33]. Figure 1.8 shows the phase-shift  $\phi_{ETF}$  versus temperature for a typical ETF. As shown, and is also clear from the above expression,  $\phi_{ETF}$  is slightly non-linear with temperature, which calls for linearization in the digital domain.

Since an ETF requires heat pulses to operate, it is naturally ill-suited to low-power applications, e.g. the ETF-based sensor presented in [1.33], requires 5mW to operate. However, decent accuracies can be obtained without trimming, and only based on batch-

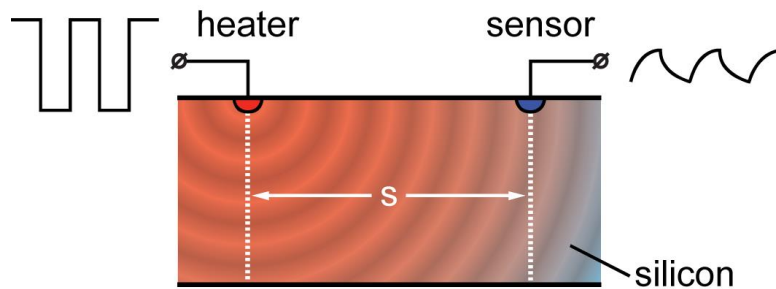


Figure 1.7: Cross section of an electro-thermal filter (ETF) consisting of a heater and a temperature sensor (thermopile) at a distance  $s$  formed in the silicon substrate.

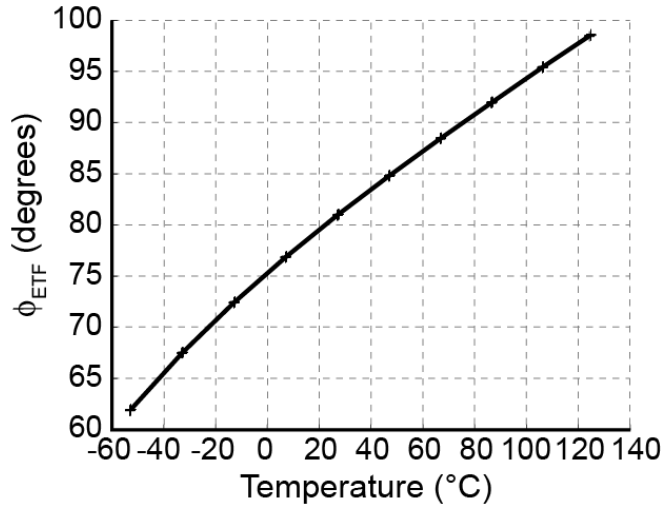


Figure 1.8: Phase shift of an electro-thermal filter (ETF) as a function of temperature [1.33].

calibration of sensors, e.g.  $\pm 0.5^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$  in  $0.7\mu\text{m}$  CMOS process [1.33], and even  $\pm 0.2^\circ\text{C}$  ( $3\sigma$ ) in  $0.18\mu\text{m}$  CMOS [1.34]. This is due to the fact that, the accuracy of ETF-based sensors depends on that of the lithography that realizes the distance  $s$ , and is thus expected to scale with every CMOS process node. This makes such sensors quite promising in applications where uncalibrated accuracy is critical, while their relatively large power consumption can be tolerated, e.g. in the thermal management of microprocessors.

### 1.3.4 MOSFETs

When biased in the sub-threshold region, the drain current  $I_D$  and the gate-source voltage  $V_{GS}$  of a MOSFET exhibit a temperature-dependent exponential relationship, similar to that between the collector current  $I_C$  and  $V_{BE}$  of a BJT [1.35]:

$$I_D^{bulk} \propto \frac{W}{L} \exp \left[ \frac{q}{mkT} (V_{GS} - V_T^{bulk}) \right], \quad (1.3)$$

where  $k$  is the Boltzmann's constant,  $T$  is the absolute temperature, and  $q$  is the electron charge, and  $W$ ,  $L$  represent the width and length of the device, respectively. The parameter  $m = 1 + C_D/C_{OX}$ , is the body effect coefficient, where  $C_D$  and  $C_{OX}$  are the depletion-layer and gate-oxide capacitances, respectively [1.35]. Similar exponential relationships between equations (1.1) and (1.3) suggest that MOSFETs can replace BJTs as

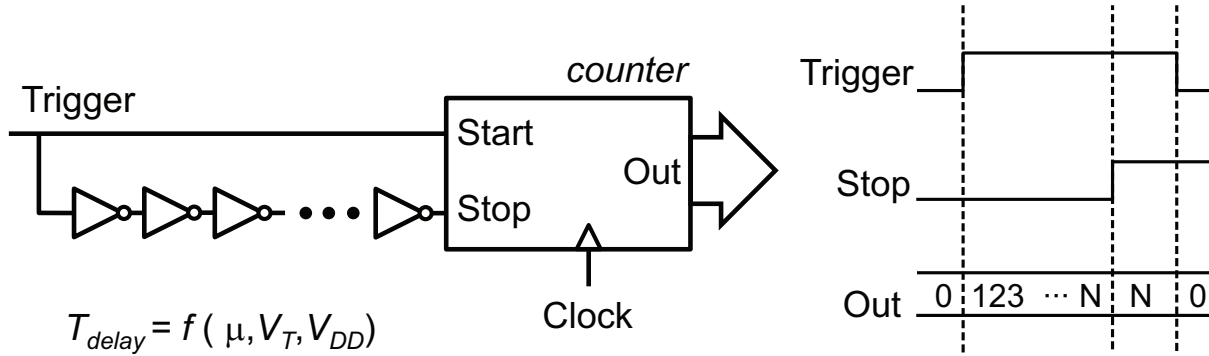


Figure 1.9: Block diagram of a MOSFET-based temperature sensor based on inverter delay.

temperature sensing elements [1.36]. Compared to BJTs, however, the gate-source voltage  $V_{GS}$  of a MOSFET biased in sub-threshold is substantially smaller and can be controlled by sizing  $W$  and/or  $L$ . This, in turn, offers a potential advantage for low supply voltage operation. However, the oxide capacitance  $C_{OX}$  suffers from process spread, while the threshold voltage  $V_T^{bulk}$  also varies due to the body-effect and suffers from the process spread as well. In consequence, MOSFET-based sensors suffer from the process spread of two different parameters, which, in turn, results in greater inaccuracies when compared to equally one-point calibrated BJTs. Therefore, MOSFET-based sensors often require two-temperature calibration to meet the accuracy requirements of most of applications.

The propagation delay of a CMOS inverter chain, or alternatively, the frequency of a ring oscillator, can also be used as a measure of temperature [1.37]. Figure 1.9 shows the operating principle of such sensors, where a counter is used to measure the propagation delay through a chain of inverters. The average propagation delay  $T_P$  of an inverter composed of balanced PMOS and NMOS devices can be expressed as [1.37]:

$$T_P = \frac{(L/W)C_L}{\mu C_{OX}(V_{DD} - V_T)} \cdot \ln \left[ \frac{3V_{DD} - 4V_T}{V_{DD}} \right], \quad (1.4)$$

in which the mobility  $\mu$  and  $V_T$  are temperature-dependent parameters. Assuming  $V_{DD} \gg V_T$ , then  $T_P$  will depend on temperature mainly through  $\mu$ . This assumption, however, becomes less and less valid in the modern CMOS processes with reduced supply voltages. Besides,  $T_P$  suffers from the process spread in  $V_T$  and from the variations in  $V_{DD}$  as well. In consequence, such sensors usually require two-point calibration and suffer from a poor power supply sensitivity, usually in the range of several degrees Celsius per Volt, e.g.  $10^\circ\text{C}/\text{V}$  in [1.38]. This is about two orders of magnitude worse than typical BJT-based sensors, and is prohibitively large for most of applications. Therefore, in practice,

such sensors should be used with voltage regulators, which calls for extra area and power consumption.

### 1.3.5 Dynamic Threshold MOSFETs (DTMOSTs)

Consider a standard MOSFET biased in sub-threshold region, with the gate and bulk terminals tied together, as shown in Figure 1.10. This connection fixes the width of the depletion layer under the channel, thereby causing the threshold voltage to vary dynamically, hence the name dynamic-threshold MOST (DTMOST). As a result, the drain current  $I_D^{DT}$  of a DTMOS transistor operated in the sub-threshold region can be expressed as follows [1.39]:

$$I_D^{DT} \propto \frac{W}{L} \exp \left[ \frac{q}{kT} (V_{GS} - V_T^{DT}) \right], \quad (1.5)$$

The key observation is that the gate-body connection ensures that the threshold voltage  $V_T^{DT}$  of a DTMOS transistor is well-defined. As a result, a diode-connected DTMOST i.e. a DTMOS diode exhibits a near-ideal exponential relationship between  $I_D^{DT}$  and  $V_{GS}$ , which is less dependent on  $C_{OX}$  and  $C_D$  [1.35], [1.39]. Figure 1.11 compares the sub-threshold characteristics of a bulk PMOST operated in both *bulk* mode (gate and substrate electrically isolated) and DTMOST mode. As shown, a DTMOST configuration would result in a steeper sub-threshold slope, lowered threshold voltage, and thus higher  $I_D$ , when compared to the bulk configuration for the same device.

More importantly, unlike the bulk configuration, the sub-threshold slope in the DTMOST configuration is well-defined and is less dependent on device related parameters, as can be also seen from equation (1.5). In other words, the process spread of  $V_{GS}$  in

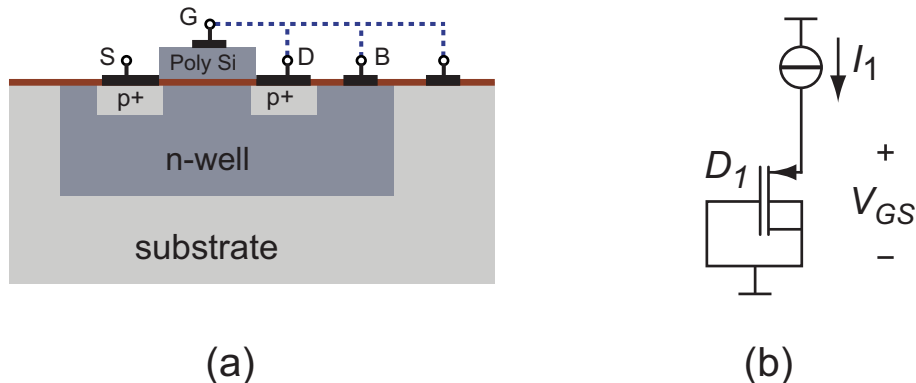


Figure 1.10: A P-type DTMOST diode; cross section view (a), symbol view (b).



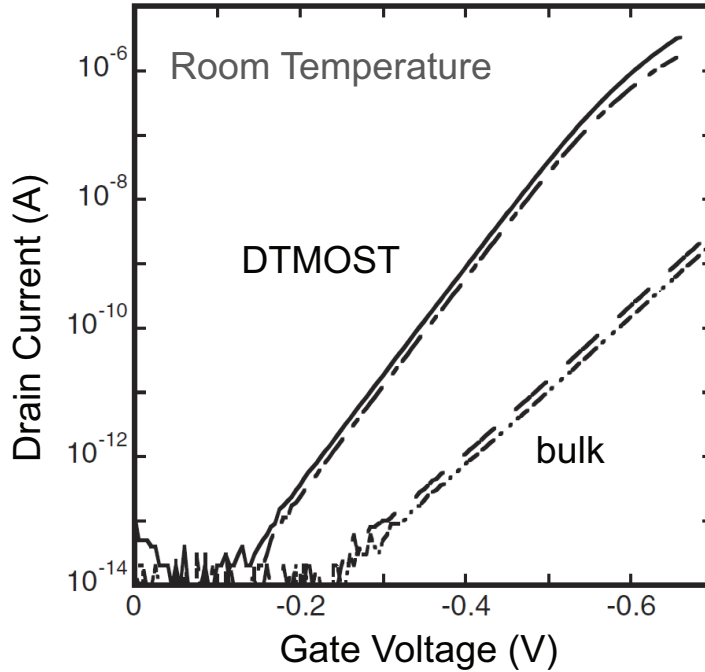


Figure 1.11: Subthreshold characteristics of a bulk PMOS device operated in both “bulk” and “DTMOST” modes, measured at room temperature [1.39].

the DTMOST configuration is less than that of the bulk configuration [1.39], [1.40]. This would suggest that similar to BJTs, DTMOSTs can be effectively calibrated at a single temperature, while offering the low-voltage capability of MOSFETs.

## 1.4 Energy-Efficiency and Resolution FoM

Given the variety of smart temperature sensors in standard CMOS, devising a single figure of merit (FoM) to assess their energy-efficiency performance would be very useful. As shown in Figure 1.1, a smart temperature sensor typically consists of an ADC that digitizes the sensor’s front-end output: usually a small signal contaminated by the thermal-noise. Moreover, the resolution of an optimally designed ADC is limited by thermal- rather than quantization-noise. This would suggest that as for general-purpose ADCs [1.41], a resolution figure-of-merit (FoM) involving the energy per conversion and resolution could be defined as follows [1.42]:

$$FoM = E_{conv} \cdot Resolution^2, \quad (1.6)$$

where  $E_{conv}$  is the amount of energy dissipated per conversion. It should be noted that

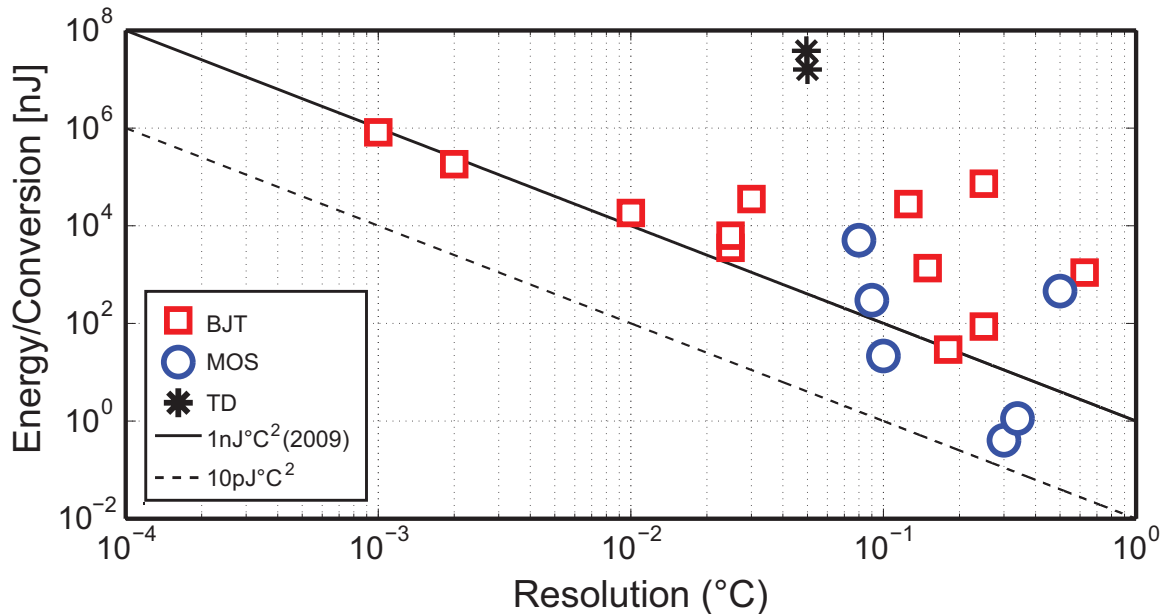


Figure 1.12: Energy per conversion versus the resolution of several smart temperature sensors, published prior to the start of this research (in 2009) [1.42]. Note that no resistor-based sensors were published before 2009.

in the context of smart temperature sensors, other figures of merit involving the sensor's accuracy might also be useful. However, this is complicated by the fact that various sensors employ different numbers of calibration points, making a fair comparison rather difficult.

Figure 1.12 shows the energy per conversion versus the resolution of several smart temperature sensors, published prior to the start of this research (in 2009). It can be seen that the resolution FoM defines a line that usefully bounds the state-of-the art, as would be expected for thermal-noise limited converters.

## 1.5 Prior-Art and Choice of Sensing Element

Table 1.1 lists some prior-art temperature sensors (back in 2009), employing different types of sensing elements. As shown, the work in [1.25] presents a PNP-based temperature sensor which draws  $75\mu\text{A}$  from a 2.5V supply. Thanks to the well-defined characteristics of PNPs in the used  $0.7\mu\text{m}$  CMOS, and combined with various dynamic correction techniques, it achieves an excellent inaccuracy of  $\pm 0.1^\circ\text{C}$  ( $3\sigma$ , 24 samples) over the military temperature range ( $-55^\circ\text{C}$  to  $125^\circ\text{C}$ ) and employing a single-temperature trim. In a conversion time of 100ms it achieves a thermal noise limited resolution of 10mK, resulting in

Parameter	JSSC'05 [1.25]	ISSCC'09 [1.43]	TCAS'09 [1.44]	JSSC'05 [1.37]	ISSCC'08 [1.33]
Sensor type	BJT	BJT	MOSFET	MOSFET	Thermal Diffusivity
CMOS technology	0.7 $\mu$ m	0.7 $\mu$ m	0.18 $\mu$ m	0.35 $\mu$ m	0.7 $\mu$ m
Chip area	4.5mm <sup>2</sup>	4.5mm <sup>2</sup>	0.032mm <sup>2</sup>	0.175mm <sup>2</sup>	2.3mm <sup>2</sup>
Supply current	75 $\mu$ A	25 $\mu$ A	0.4 $\mu$ A	3 $\mu$ A	1mA
Min supply voltage	2.5V	2.5V	1V	3.3V	5V
Supply sensitivity	0.03°C/V	0.05°C/V	8°C/V	0.1°C/V	-
Inaccuracy (calibration points)	$\pm 0.1^\circ\text{C}$ (3 $\sigma$ ) (1)	$\pm 0.1^\circ\text{C}$ (3 $\sigma$ ) (1)	-0.8 / +1°C (2)	-0.7 / +0.9°C (2)	$\pm 0.5^\circ\text{C}$ (3 $\sigma$ ) (batch-calibrated)
Temperature range	-55°C - 125°C	-55°C - 125°C	0°C - 100°C	0°C - 100°C	-55°C - 125°C
Resolution (T <sub>conv</sub> )	0.01°C (100msec)	0.025°C (100msec)	0.3°C (1msec)	0.08°C (500msec)	0.05°C (3030msec)
Resolution FoM	1.9nJ°C <sup>2</sup>	3.9nJ°C <sup>2</sup>	36pJ°C <sup>2</sup>	32nJ°C <sup>2</sup>	38 $\mu$ J°C <sup>2</sup>

Table 1.1: Performance summary of prior-art temperature sensors, back in 2009.

a FoM of 1.9nJ°C<sup>2</sup>. For the targeted energy-efficiency in this work, however, the current consumption of 75 $\mu$ A is prohibitively large and can not be afforded in wireless temperature sensing applications. Moreover, it uses a rather old process (0.7 $\mu$ m CMOS), resulting in a large die area of 4.5mm<sup>2</sup>. The work in [1.43] shows that similar accuracy can be achieved even after reducing the bias current of PNPs to 100nA, i.e. by 10x compared to [1.25]. This, suggests that PNPs are capable of achieving a decent accuracy over a wide range of bias current and using a single-temperature trim.

The work in [1.44] presents a sensor based on MOSFET transistors operating in the linear region. It is quite power-efficient, drawing only 400nA from a 1V supply. In a conversion time of 1ms, it achieves a resolution of 300mK, thus obtaining a decent energy-efficiency evidenced by a FoM of 36pJ°C<sup>2</sup>. It is also quite compact, occupying only 0.032mm<sup>2</sup> in a 0.18 $\mu$ m CMOS process. However, over the temperature range from 0°C to 100°C it achieves an inaccuracy of -0.8/+1°C after a costly two-end-point trim (4 samples). Besides, it suffers from a poor supply dependency of 8°/V. In [1.37], a sensor based on the temperature dependence of delay lines is presented. For the class of MOSFET-based designs, it achieves the best inaccuracy of -0.7/+0.9°C over the temperature range from 0°C to 100°C, but only after a two-end-point trim (25 samples). It also requires a long conversion time to operate; 500ms to achieve a resolution of 80mK, while drawing 3 $\mu$ A from a 3.3V supply. This results in a poor energy-efficiency of 32nJ°C<sup>2</sup>.

The work in [1.33] presents an ETF-based temperature sensor. It achieves an untrimmed (but batch-calibrated) inaccuracy of  $\pm 0.5^\circ\text{C}$  ( $3\sigma$ , 16 samples) over the military temperature range ( $-55^\circ\text{C}$  to  $125^\circ\text{C}$ ). The fact that a decent accuracy can be achieved without trimming makes ETF-based sensors very attractive, especially for low-cost applications, in which calibration cost needs to be minimized. Nevertheless, as previously mentioned, such sensors are rather power-hungry, e.g. drawing 1mA from a 5V supply in [1.33].

Studying the different sensors in Table 1.1 reveals the fact that BJTs are the devices of choice for wireless temperature sensing applications, in which energy-efficiency, low-cost, and precision are key requirements. Their bias current can be reduced to the tens of nA level with minimal impact on the accuracy. They also exhibit a well-defined process spread which can be effectively trimmed at a single-temperature. None of the MOSFET- or ETF-based sensors can offer such a combination. Back in 2009, no resistor- or DTMOST-based sensors existed, and thus these types are missing in the Table 1.1. It is, however, well understood that resistor-based sensors, although potentially suitable for use as energy-efficient sensors, require multiple-temperature calibration which rules out their use for low-cost applications. DTMOSTs, on the other hand, seem to be capable of offering high accuracy using a single-temperature trim. In order to operate in weak inversion, their bias current is naturally in the sub- $\mu\text{A}$  range, thus offering low-power operation as well. A distinct advantage of DTMOSTs compared to BJTs would be their capability for sub-1V supply operation. However, no experimental DTMOST-based sensor prototype had been realized prior to this research and thus no apples-to-apples comparison could be made.

## 1.6 Thesis Organization

In the next chapter, an architecture-level solution to realize energy-efficient BJT-based sensors is presented. This is done by analyzing existing sensor architectures and finding out the root causes of their lack of energy-efficiency. To bridge the energy-efficiency gap between the existing solutions and requirements of the target application, a novel *zoom*-ADC architecture will be presented which results in low-power and fast conversions, thereby significantly improving the energy-efficiency. This, in turn, results in a compact, and therefore, low-cost solution. In Chapter 3, two sensor prototypes based on 1<sup>st</sup>- and 2<sup>nd</sup>-order zoom-ADCs will be presented. Moreover, another prototype is presented for use in automotive applications, where sensing high temperatures ( $>150^\circ\text{C}$ ) is desired. In this research, DTMOST-based sensors have been also realized, which is the subject of Chapter 4. First, a DTMOST-based prototype based on a 1<sup>st</sup>-order zoom-ADC will be presented and an apples-to-apples comparison to BJT-based sensors will be made. It is shown

that DTMOSTs are in fact capable of achieving high accuracy after a single-temperature trim, and so can bridge the accuracy gap between the MOSFET- and BJT-based sensors. Based on the low-voltage operation capability of DTMOSTs, a second prototype is then realized which provides both sub-1V, and sub- $\mu$ A operations. Chapter 5 is devoted to the conclusions and a comparison between the different sensor prototypes realized in this research and state-of-the-art.

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## Chapter 2

# Readout Methods for BJT-Based Temperature Sensors

## 2.1 Introduction

As discussed in the previous chapter, BJT-based temperature sensors are promising candidates for use in wireless temperature sensing applications. In this chapter, we first describe the operating principle of BJT-based sensors, followed by an overview of various readout methods. The energy-efficiency of these methods is then discussed, and compared to the ultimate achievable efficiency of BJT-based sensors.

## 2.2 Operating Principle of BJT-Based Sensors

In general, a digital representation of temperature can be obtained by generating a temperature-dependent voltage and a temperature-independent or reference voltage  $V_{REF}$  and then digitizing the ratio with an analog-to-digital converter (ADC). In BJT-based temperature sensors, the bandgap voltage of silicon  $V_{bg} \approx 1.2\text{V}$  serves as the reference voltage, and hence they are often known as *bandgap temperature sensors*.

A smart temperature sensor therefore requires both a well behaved temperature dependent signal, preferably one that is linearly proportional-to-absolute temperature (PTAT), as well as a reference voltage  $V_{REF}$  for temperature readout. The main advantage of BJTs compared to other temperature sensing elements is that they can be used to generate both  $V_{PTAT}$  and  $V_{REF}$  simultaneously and with minimal circuit overhead. Moreover, the output of the resulting temperature sensor is a linear function of temperature, and thus no extra post processing is required. Lastly, but perhaps most importantly for

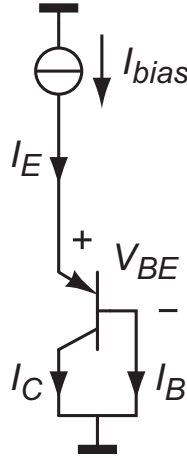


Figure 2.1: Diode-connected configuration of a PNP transistor [2.5].

industrial applications, single-temperature trimming is enough to achieve high precision, e.g. the  $\pm 0.1^\circ\text{C}$  inaccuracy achieved in [2.1], [2.2].

In the 160nm process (CMOS-14) used in this work, the most suitable BJT is a substrate-PNP. Compared to the lateral NPN which is also available, it is less sensitive to process spread and packaging stress [2.3], [2.4]. However, its main drawback is that its collector is grounded, which in turn, limits the number of practical circuit topologies. In fact, the best way of using the substrate PNP is in a diode-connected configuration, as shown in Figure 2.1, in which the base-collector junction of the device is short circuited ( $V_{BC} = 0$ ) [2.5]. This reduces leakage currents via the base-collector junction to negligible levels. It also reduces errors due to the transistor's finite Early voltage. Such leakage currents would otherwise disturb the collector current accuracy, which, in turn, would impact the sensor's ultimate accuracy.

### 2.2.1 Temperature Characteristics of BJTs

Figure 2.2.a illustrates two bipolar transistors  $Q_1$  and  $Q_2$  (with emitter area of  $r \cdot A_E$  and  $A_E$ , respectively), configured in a “diode-connected” fashion, biased by currents  $I_1$  and  $p \cdot I_1$ . For a diode-connected bipolar transistor, the collector current  $I_C$  can be expressed as [2.5]:

$$I_C = I_S \left( \exp \left( \frac{qV_{BE}}{kT} \right) - 1 \right), \quad (2.1)$$

where  $k$ ,  $T$  and  $q$  represent the Boltzmann constant ( $1.38 \times 10^{-23}$  J/K), temperature in Kelvin and electron charge ( $1.6 \times 10^{-19}$ C), respectively.  $V_{BE}$  is the base-emitter voltage

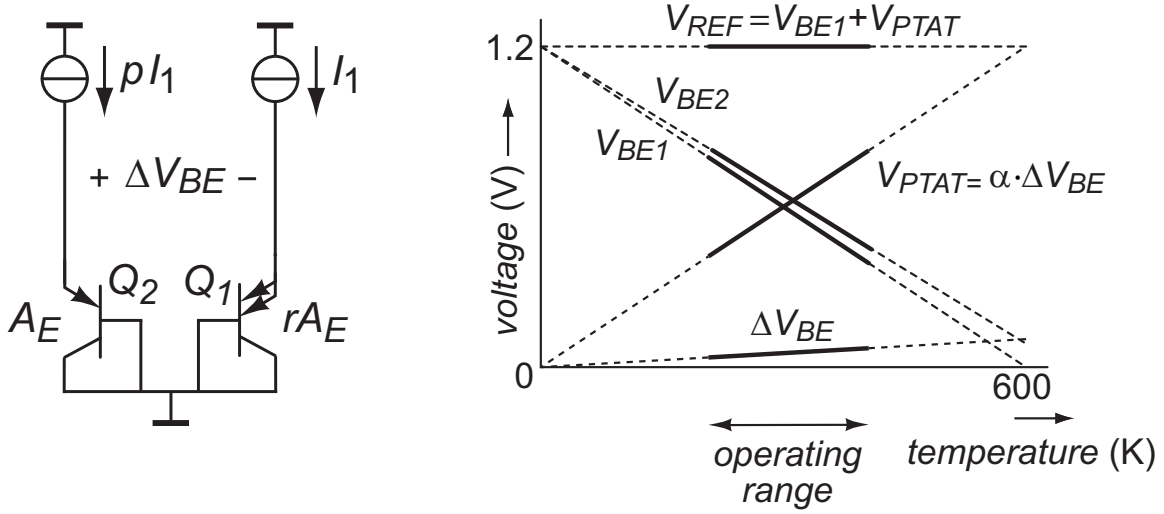


Figure 2.2: (a) Two diode-connected PNPs biased at a collector current density ratio of  $pr$ ; (b) these transistors are used to generate  $V_{PTAT}$  and  $V_{REF}$ , the ratio of which is a measure of temperature [2.5].

difference, and the parameter  $I_S$  denotes the saturation current of the diode-connected bipolar transistor, and is given by [2.5]:

$$I_S = CT^\eta \exp\left(\frac{-qV_{g0}}{kT}\right). \quad (2.2)$$

In this equation,  $C$  is a constant,  $V_{g0}$  is the extrapolated bandgap voltage of silicon at 0K, and  $\eta \approx 4$  is a constant for PNP transistors in CMOS technology [2.6]. We observe that the saturation current  $I_S$  is proportional to  $T^4$ , implying a strong sensitivity to temperature.

By solving equation (2.1), the following expression can be obtained for  $V_{BE}$ :

$$V_{BE} = \frac{kT}{q} \ln\left(\frac{I_C}{I_S} + 1\right). \quad (2.3)$$

At room temperature,  $kT/q$  is about 26mV. Although  $V_{BE}$  would seem to have a positive temperature coefficient, it actually has a negative temperature coefficient (about -2mV/°C) due to the strong temperature dependency of the saturation current  $I_S$ . In other words,  $V_{BE}$  is, to first order, complementary-to-absolute temperature (CTAT). Figure 2.2.b shows how  $V_{BE}$  rolls off from  $\approx 1.2V$  to zero over a temperature range of almost 600°

For two bipolar transistors, e.g.  $Q_1$  and  $Q_2$ , biased at different collector current

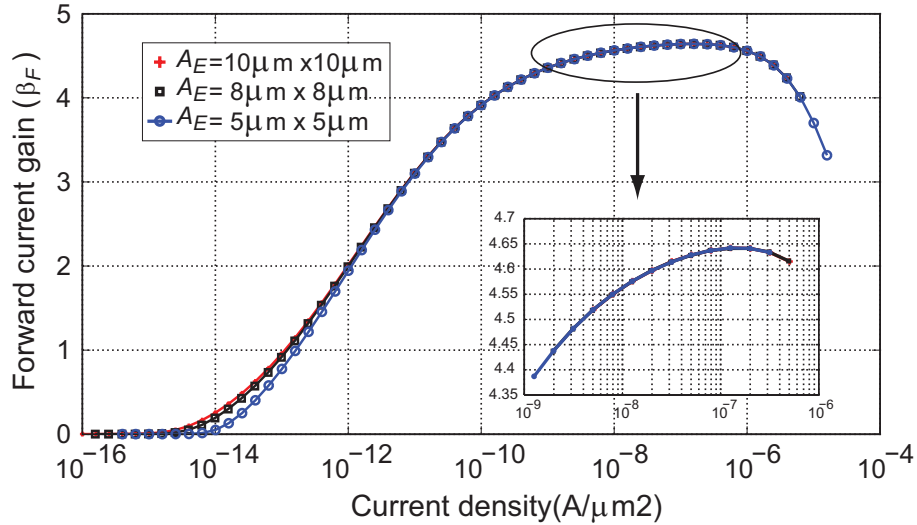


Figure 2.3: Forward current gain  $\beta_F$  (at 25°C) of substrate PNPs as a function of collector current density and for different emitter areas in the CMOS-14 process used.

densities, the difference between their base-emitter voltages is a PTAT voltage:

$$\Delta V_{BE} = V_{BE2} - V_{BE1} = \frac{kT}{q} \ln(pr), \quad (2.4)$$

in which  $r$  is the emitter area ratio between two bipolar transistors, and  $p$  is the collector current ratio (see Figure 2.2.a). As shown in Figure 2.2.b,  $\Delta V_{BE}$ , unlike  $V_{BE}$ , has a positive temperature coefficient which depends on the constants  $k/q$  and the product  $pr$ . In the rest of this chapter, we will assume the use of two identical bipolar transistors i.e.  $r = 1$ , for simplicity.

For substrate PNPs, however, it is their emitter current ratio, rather than their collector current ratio, that is under control. These ratios are equal if the PNP's forward current gain  $\beta_F$  can be assumed to be constant, i.e. independent of emitter current. This assumption, however, is only valid for older CMOS processes, e.g.  $0.7\mu\text{m}$  and  $0.5\mu\text{m}$ , where  $\beta_F$  is constant for at least two decades of emitter current [2.1], [2.7]. Figure 2.3 illustrates the forward current gain  $\beta_F$  as a function of emitter current for substrate PNPs in the target CMOS-14 process. We clearly observe that no flat region exists, thus implying that the collector current ratio is only approximately equal to  $p$  [2.8]. The impact of this is discussed in detail in [2.8], and is thus neglected in the rest of this chapter.

From a linear combination of  $V_{BE}$  (CTAT) and  $\Delta V_{BE}$  (PTAT), a reference voltage  $V_{REF}$ , with a nominally zero temperature-coefficient can be generated. As shown in Figure 2.2.b, this can be done by scaling  $\Delta V_{BE}$  with a constant  $\alpha$  such that the positive temperature-coefficient of  $\alpha \cdot \Delta V_{BE}$  compensates for the negative temperature-coefficient

of  $V_{BE}$ :

$$V_{REF} = V_{BE} + \alpha \cdot \Delta V_{BE}. \quad (2.5)$$

Since  $\Delta V_{BE} = 0$  at  $T = 0\text{K}$ ,  $V_{REF}$  is equal to the extrapolated value of  $V_{BE}$  at zero Kelvin, which is related to the bandgap voltage of silicon ( $V_{bg} \approx 1.2\text{V}$ ). Notably, both of the signals required for a ratiometric temperature measurement ( $V_{PTAT}$  and  $V_{REF}$ ) can be obtained by biasing BJTs at different collector current densities. The scaling factor  $\alpha$  can be found by solving equation (2.5) such that the temperature coefficient of  $V_{REF}$  is zero:

$$S_{V_{BE}}^T = \left| \frac{\partial V_{BE}}{\partial T} \right| \approx 2\text{mV}/^\circ\text{C} = \alpha \cdot \frac{k}{q} \cdot \ln(p), \quad (2.6)$$

indicating that:

$$\alpha \approx \frac{S_{V_{BE}}^T}{\frac{k}{q} \ln(p)}. \quad (2.7)$$

The larger the collector current density ratio  $p$  is, the smaller the scaling factor  $\alpha$  becomes. Practical values of  $p$  typically range between 2 and 20, which means that  $\alpha$  typically ranges between 7 and 33. It should be noted that since  $S_{V_{BE}}^T$  is a weak function of  $I_C$ , as can be seen from equation (2.3),  $\alpha$  can also be tuned by adjusting  $I_C$ .

## 2.3 Generic BJT Readout

### 2.3.1 Topology

A block diagram of a generic bandgap temperature sensor is shown in Figure 2.4. In the analog front-end, two diode-connected bipolar devices  $Q_1$  and  $Q_2$  are biased at a bias current density ratio of  $p$  to generate the PTAT signal  $\Delta V_{BE}$ . The reference voltage  $V_{REF}$  then can be obtained by scaling up  $\Delta V_{BE}$  and adding to  $V_{BE1}$ . Finally  $\alpha \cdot \Delta V_{BE}$  and  $V_{REF}$  voltages are fed to an analog-to-digital converter (ADC) for temperature reading. The ADC then outputs the ratio  $\mu$  [2.5]:

$$\mu = \frac{\alpha \cdot \Delta V_{BE}}{V_{BE} + \alpha \cdot \Delta V_{BE}} = \frac{V_{PTAT}}{V_{REF}}, \quad (2.8)$$

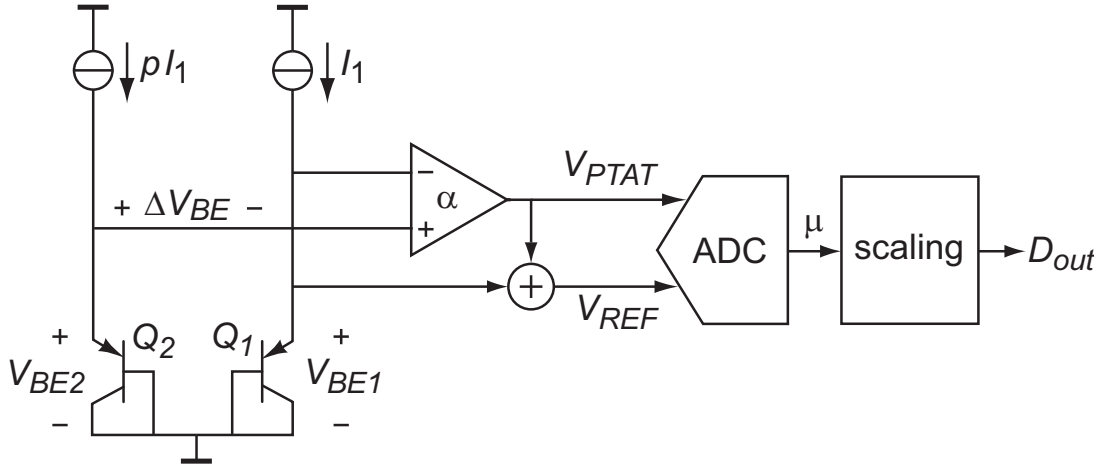


Figure 2.4: Operating principle of generic bandgap temperature sensors: two diode-connected PNPs generate  $\Delta V_{BE}$  and  $V_{BE}$ .  $\Delta V_{BE}$  is then amplified and combined with  $V_{BE1}$  to provide  $V_{PTAT}$  and  $V_{REF}$  at the input of an ADC for a ratio-metric measurement [2.5].

which is PTAT and varies between 0 and 1 over a temperature range of about 600K. The final step is to linearly scale  $\mu$  to obtain a digital output  $D_{out}$  in degrees Celsius [2.5]:

$$D_{out} = A \cdot \mu + B, \quad (2.9)$$

where  $A \approx 600\text{K}$  and  $B \approx -273\text{K}$ .

Combining  $V_{BE}$  and  $\Delta V_{BE}$  to generate  $V_{REF}$ , as shown in Figure 2.4, requires circuitry and thus power dissipation. As presented in [2.1], this can be avoided by an alternative scheme in which  $V_{BE}$  and  $\Delta V_{BE}$  are directly input to a  $\Delta\Sigma$ -ADC. The charge balancing in the ADC then ensures that its output is the desired  $\mu$ , i.e.  $V_{REF}$  is implicitly generated in a dynamic manner. The result is a very simple sensor front-end.

As the PTAT ratio  $\mu$  lies between 1/3 and 2/3 at the two extremes of the full military temperature range ( $-55^\circ\text{C}$  to  $125^\circ\text{C}$ ), only about 30% of the ADC's dynamic range is used in this configuration. This is also evident from Figure 2.2.b. In other words, the required resolution of the ADC has to be about three times larger than the desired temperature resolution, which means that about 70% of the ADC's dynamic range is wasted. In many other applications, the target temperature range is even smaller, e.g. from  $-40^\circ\text{C}$  to  $85^\circ\text{C}$  and from  $25^\circ\text{C}$  to  $45^\circ\text{C}$  for industrial and clinical applications, respectively, which further exacerbates the situation.

Another combination of  $V_{BE}$  and  $\Delta V_{BE}$  can be employed to efficiently utilize the ADC's dynamic range and avoid the wasted resolution. As presented in [2.9], [2.10], a

linear combination of  $V_{BE}$  and  $\Delta V_{BE}$  can be used as the input to the ADC, leading to the following ratio:

$$\mu' = \frac{2\alpha \cdot \Delta V_{BE} - V_{BE}}{V_{BE} + \alpha \cdot \Delta V_{BE}}. \quad (2.10)$$

Now, as is evident from Figure 2.5, almost 90% of the ADC's dynamic range is used by the ratio  $\mu'$  over the military temperature range. In a recent work [2.11], a different combination of  $V_{BE}$  and  $\Delta V_{BE}$  is used to efficiently utilize the ADC's dynamic range over the clinical temperature range from 25°C to 45°C. Employing the generic readout over such a limited range will result in about 96% waste of the ADC's dynamic range. As presented,  $V_{BE}$  and  $\Delta V_{BE}$  are dynamically combined to obtain a linear ratio which varies by over 60% from 25°C to 45°C.

From equation (2.8), it can be seen that errors in the scaling factor  $\alpha$  will impact the accuracy of  $\mu$ , and thus the digital temperature reading  $D_{out}$ . As this scaling has to be done in the analog domain, precision techniques like dynamic element matching (DEM) are needed to realize precision temperature sensors [2.5], [2.12]. In addition,  $\alpha$  should be an integer or rational number to enable accurate on-chip implementation [2.13]. This requirement imposes a restriction on the values of  $I_C$  and the collector current density ratio  $p$ , through equation (2.7).

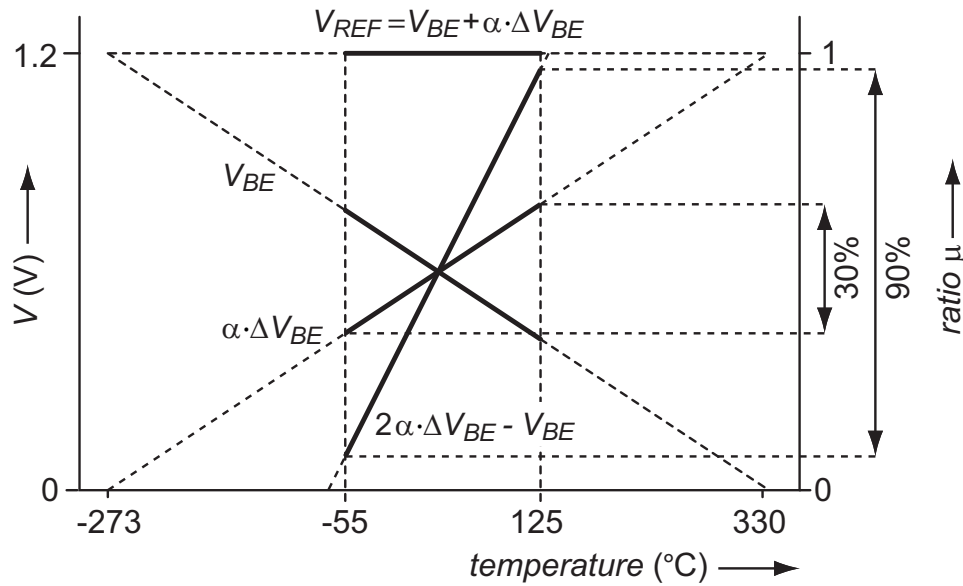


Figure 2.5: More efficient use of the ADC's dynamic range by using a different combination of  $V_{BE}$  and  $\Delta V_{BE}$  [2.5].



### 2.3.2 ADC Resolution

Essentially, the ADC resolution will depend on the target application. Given the diversity of temperature sensing applications, a wide range of resolutions are required, e.g. a couple of hundred milli-Kelvin in environmental monitoring, to sub milli-Kelvin for the temperature compensation of MEMS frequency references. Moreover, the need for accurate calibration (to meet a given accuracy specification) imposes further limits. As a rule of thumb, the sensor's resolution should be an order of magnitude less than the target accuracy to ensure negligible calibration error, e.g.  $\pm 0.01^\circ\text{C}$  resolution if an accuracy of  $\pm 0.1^\circ\text{C}$  is targeted.

In an ADC, the effective number of bits (ENOB), which expresses the total quantization noise of the ADC as a fraction of its full scale, can be expressed as:

$$ENOB = \log_2 \left( \frac{D_{FS}}{\max |D_{out} - D_{out,ideal}|} \right) - 1, \quad (2.11)$$

where  $D_{FS}$  and  $|D_{out} - D_{out,ideal}|$  represent the full-scale value and quantization error of the ADC, respectively [2.5].

For the generic bandgap temperature sensor,  $D_{FS}$  can be determined from Figure 2.2.b. As depicted,  $V_{REF} \approx 1.2\text{V}$ , which is the full scale value in the voltage domain, and corresponds to a temperature range of about  $600^\circ\text{C}$  i.e.  $D_{FS} \approx 600^\circ\text{C}$ . Assuming a target temperature resolution of  $\pm 0.01^\circ\text{C}$  in the calibration phase, an ENOB of 14.9 bits is required, which is not trivial to achieve together with extremely low-energy dissipation. By combining  $V_{BE}$  and  $\Delta V_{BE}$  at the ADC's input as in Figure 2.5,  $D_{FS}$  reduces to  $\approx 200^\circ\text{C}$ . This would mean that for a  $\pm 0.01^\circ\text{C}$  resolution, an ENOB of 13.3 bits will be required in the ADC, which is about 1.6 bits less than in the generic approach.

## 2.4 Energy-Efficiency of BJT-Based Sensors

A smart temperature sensor consists of two major blocks; the sensor front-end and the ADC. Since achieving high energy-efficiency is the main focus of this research, a careful energy-efficiency analysis of existing temperature sensor architectures is essential and will be discussed in the following sections.

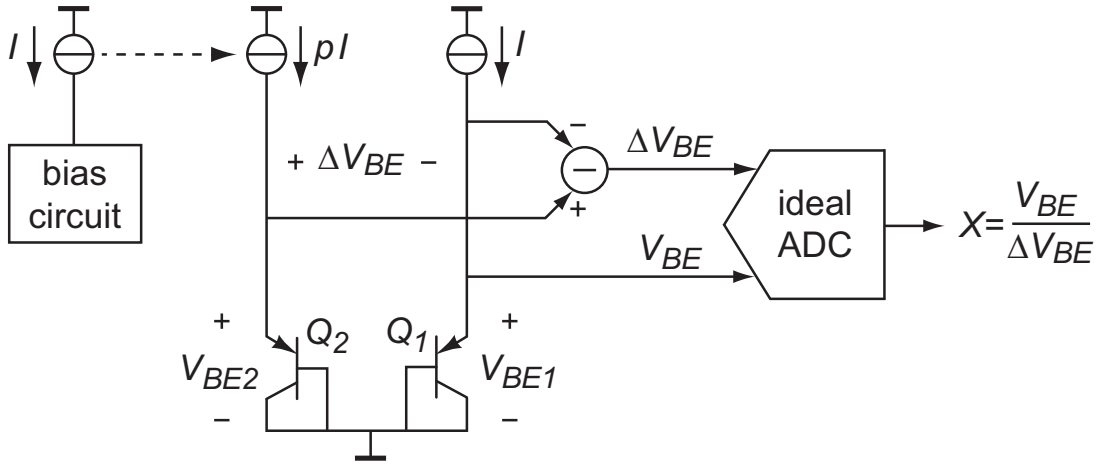


Figure 2.6: The circuit diagram used to determine the front-end's energy-efficiency using bias current scaling in the front-end.

### 2.4.1 Efficiency Limits of a BJT-Based Front-End

Regardless of the chosen ADC topology, the ultimate energy-efficiency of a smart temperature sensor is limited to that of its front-end. To calculate this efficiency limit, we assume the circuit diagram of Figure 2.6. Here, two identical diode-connected PNPs are biased with the bias currents  $I$  and  $p \cdot I$ , mirrored from a bias circuit (not shown). The sensor core generates  $V_{BE}$  and  $\Delta V_{BE}$ , whose ratio is then digitized by means of an ideal ADC, i.e. an ADC with infinite resolution and zero power dissipation.

Recall from Chapter 1, that the resolution figure-of-merit (FoM) is expressed as follows:

$$FoM = E_{conv} \cdot Resolution^2, \quad (2.12)$$

where,  $E_{conv}$  represents the dissipated energy per conversion. Neglecting the power dissipated in the bias circuit,  $E_{conv}$  can be calculated as follows:

$$E_{conv} = V_{DD} \cdot (1 + p) \cdot I \cdot T_{conv}, \quad (2.13)$$

where  $T_{conv}$  is the sensor's conversion time.

The sensor's resolution, on the other hand, is limited by the noise present in the sensor's output voltages i.e.  $V_{BE}$  and  $\Delta V_{BE}$ . The noise present in the base-emitter voltage of a diode-connected PNP is the combination of thermal noise and shot noise

associated with its base resistance and collector current, respectively:

$$\begin{aligned} v_{n,V_{BE}}^2 &= \frac{i_{n,c}^2}{g_m^2} + 4kTR_B B_n = 2qI B_n \left( \frac{kT}{qI} \right)^2 + 4kTR_B B_n \\ &= \frac{2kT}{g_m} B_n + 4kTR_B B_n, \end{aligned} \quad (2.14)$$

where  $g_m$  is the PNP's transconductance and  $B_n$  is the noise bandwidth. In this analysis the contribution of  $1/f$  noise is neglected, as it is relatively small in bipolar transistors. Also, the noise contribution of the base resistance  $R_B$  is usually negligible. This is due to the fact that, in practice, PNPs are often biased with current levels in the range of  $\mu\text{A}$  or even smaller. The associated  $1/g_m$  is, therefore, significantly larger than  $R_B$ , and thus dominates the noise in the base-emitter voltage  $V_{BE}$ . Assuming that the noise in the bias current  $I$  is equally dominated by the shot noise of the PNP transistors of the bias circuit, the overall noise in the base emitter voltage  $V_{BE}$  can be formulated as follows:

$$v_{n,V_{BE}}^2 = \frac{2kT}{g_m} B_n + \frac{2qI}{g_m^2} B_n = \frac{4kT}{g_m} B_n. \quad (2.15)$$

The noise present in  $\Delta V_{BE}$ , can then be obtained by adding the noise of two PNPs, as follows:

$$v_{n,\Delta V_{BE}}^2 = v_{n,V_{BE1}}^2 + v_{n,V_{BE2}}^2 = \frac{4kT}{g_m} B_n \cdot \left( 1 + \frac{1}{p} \right), \quad (2.16)$$

which is only slightly larger than  $v_{n,V_{BE}}^2$ .

Assuming that the succeeding (ideal) ADC integrates  $v_{n,V_{BE}}^2$  and  $v_{n,\Delta V_{BE}}^2$  during a conversion time  $T_{conv}$ , which is equivalent to filtering the noise with a sinc filter with a noise bandwidth  $B_n = 1/(2 \cdot T_{conv})$ , the integrated noise voltage can then be expressed as follows:

$$\begin{aligned} v_{n,V_{BE}} &= \sqrt{\frac{2kT}{g_m} \cdot \frac{1}{T_{conv}}} \\ v_{n,\Delta V_{BE}} &= \sqrt{\frac{2kT}{g_m} \cdot \frac{1}{T_{conv}} \cdot \left( 1 + \frac{1}{p} \right)}. \end{aligned} \quad (2.17)$$

The associated noise at the sensor output can now be readily found by calculating the

sensitivity of output  $D_{out}$  to  $V_{BE}$  and  $\Delta V_{BE}$ . This is achieved by differentiating equation (2.9) with respect to  $V_{BE}$  and  $\Delta V_{BE}$  [2.5]:

$$S_{V_{BE}}^{D_{out}}(T) = \frac{\partial D_{out}}{\partial V_{BE}} = A \cdot \frac{\partial \mu}{\partial V_{BE}} = -\frac{T}{V_{REF}}, \quad (2.18)$$

$$S_{\Delta V_{BE}}^{D_{out}}(T) = \frac{\partial D_{out}}{\partial \Delta V_{BE}} = A \cdot \frac{\partial \mu}{\partial (\Delta V_{BE})} = \frac{A - T}{V_{REF}} \cdot \alpha. \quad (2.19)$$

Clearly, the sensitivity to noise present in  $\Delta V_{BE}$  is considerably larger than that of  $V_{BE}$ . Moreover, given the fact that  $\Delta V_{BE}$  is always noisier than  $V_{BE}$  (see equation (2.17)), only the noise contribution of the former will be considered in the following calculations. The noise at the output  $D_{out}$  can be computed as:

$$\sigma_T^2 = S_{\Delta V_{BE}}^{D_{out}}(T)^2 \cdot v_{n,\Delta V_{BE}}^2 = \left(\frac{A - T}{V_{REF}}\right)^2 \cdot \alpha^2 \cdot \frac{2kT}{g_m} \cdot \frac{1}{T_{conv}} \cdot \left(1 + \frac{1}{p}\right). \quad (2.20)$$

The front-end FoM can now be calculated by substituting  $\sigma_T^2$  into equation (2.12):

$$FoM = V_{DD} \cdot (1 + p) \cdot I \cdot T_{conv} \left(\frac{A - T}{V_{REF}}\right)^2 \cdot \alpha^2 \cdot \frac{2kT}{g_m} \cdot \frac{1}{T_{conv}} \cdot \left(1 + \frac{1}{p}\right), \quad (2.21)$$

which can be simplified to:

$$FoM = 2 \cdot \frac{(p + 1)^2}{p} \cdot V_{DD} \cdot q \cdot V_T^2 \cdot \alpha^2 \cdot \left(\frac{A - T}{V_{REF}}\right)^2. \quad (2.22)$$

Surprisingly, the FoM is independent of both the bias current  $I$  or the conversion time  $T_{conv}$ . This is due to the fact that  $E_{conv} \propto (T_{conv} \cdot I)$ , while  $\sigma_T^2 \propto 1/(T_{conv} \cdot I)$ , and so these terms cancel out in the FoM expression. On the other hand, the parameter  $\alpha$  depends on  $p$  through equation (2.7). Therefore, the energy-efficiency of the sensor core exclusively depends on  $p$  and  $V_{DD}$  as design parameters. Figure 2.7 shows the FoM versus temperature and for different  $p$  values, assuming  $V_{DD} = 1.8V$ ; the nominal supply value in the target  $0.16\mu m$  CMOS process. As shown, FoM exhibits a parabolic profile with a peak around room temperature. By differentiating equation (2.22) with respect to  $T$ , it can be shown that FoM reaches to its maximum at  $T = A/2 \approx 300K$ .

Alternatively, the front-end circuit of Figure 2.8 can be used to generate  $V_{BE}$  and

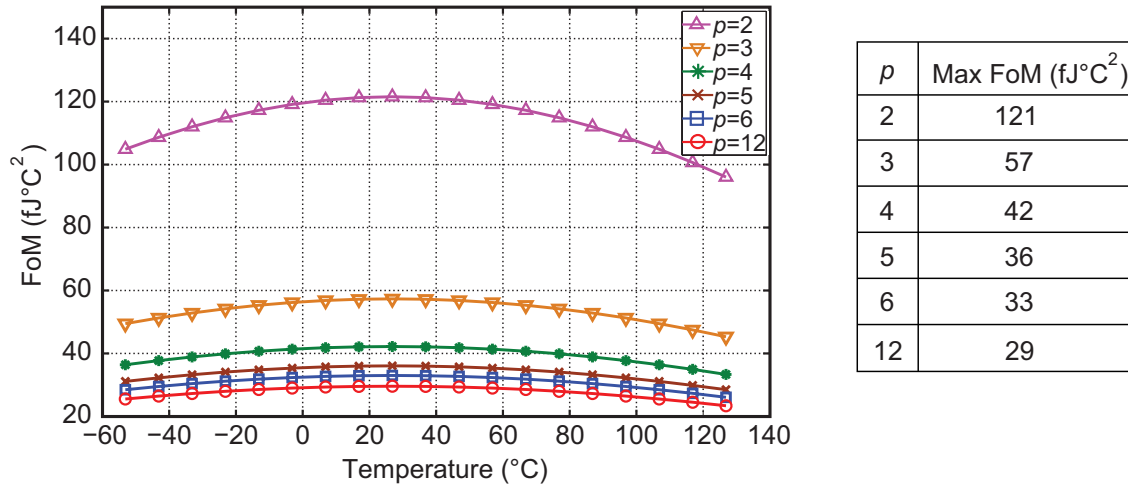


Figure 2.7: Figure-of-merit (FoM) as a function of temperature and for different  $p$  values in the bias current scaling topology.

$\Delta V_{BE}$ . In this circuit two PNPs with an emitter area ratio of  $1 : p$  are used, each of which is biased with equal bias current of  $I$ . The advantage of this biasing scheme is that both BJTs now have the same  $g_m$  and hence, generate the same noise. The collector current density ratio, essential to develop  $\Delta V_{BE}$ , is thus achieved by scaling the emitter area, as opposed to scaling the bias current in Fig 2.6. In this circuit the energy per conversion  $E_{conv}$  can be calculated as follows:

$$E_{conv} = 2 \cdot V_{DD} \cdot I \cdot T_{conv}. \quad (2.23)$$

To calculate the FoM of this circuit we first need to find  $v_{n,\Delta V_{BE}}^2$ . Neglecting the noise

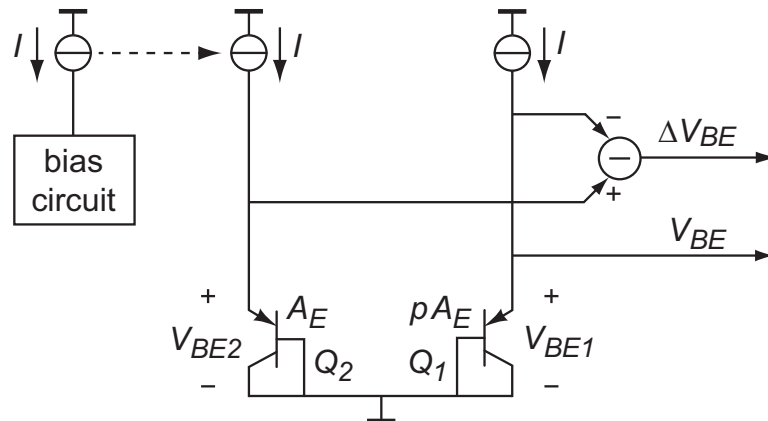


Figure 2.8: The alternative front-end circuit diagram using scaling of the emitter area.

associated with the base resistance  $R_B$ ,  $v_{n,\Delta V_{BE}}^2$  can be found as:

$$v_{n,\Delta V_{BE}}^2 = v_{n,V_{BE1}}^2 + v_{n,V_{BE2}}^2 \approx \frac{8kT}{g_m} B_n. \quad (2.24)$$

By multiplying by  $S_{\Delta V_{BE}}^{D_{out}}(T)^2$ , the noise at the sensor output  $\sigma_T^2$  and then the FoM can be calculated as:

$$FoM = 8 \cdot V_{DD} \cdot q \cdot V_T^2 \cdot \alpha^2 \cdot \left( \frac{A - T}{V_{REF}} \right)^2. \quad (2.25)$$

Again, the energy-efficiency of front-end depends on  $V_{DD}$  and  $p$  (via  $\alpha$ ) as design parameters. Figure 2.9 shows the FoM as a function of  $p$  and assuming  $V_{DD} = 1.8V$ . When compared to Figure 2.7, it can be seen that scaling the emitter area results in a more efficient sensor front-end. This becomes more evident by comparing equation (2.22) to (2.25). Clearly, scaling the bias current results in a  $(2 + p + 1/p)/4$  times larger FoM, and thus inferior efficiency. Figure 2.10 illustrates the efficiency comparison as a function of  $p$ : the larger the scaling factor  $p$ , the less efficient bias current scaling becomes compared to emitter area scaling, e.g. for  $p = 6$ , bias current scaling is  $\approx 2x$  less energy efficient.

Equations (2.22) and (2.25) set the lower bound on the energy-efficiency of BJT-based temperature sensors. In practice, the efficiency will drop significantly as a result of the extra noise and power consumption of other blocks such as the bias circuit, ADC, and the digital back-end. It is, therefore, worthwhile to compare the obtained theoretical limits

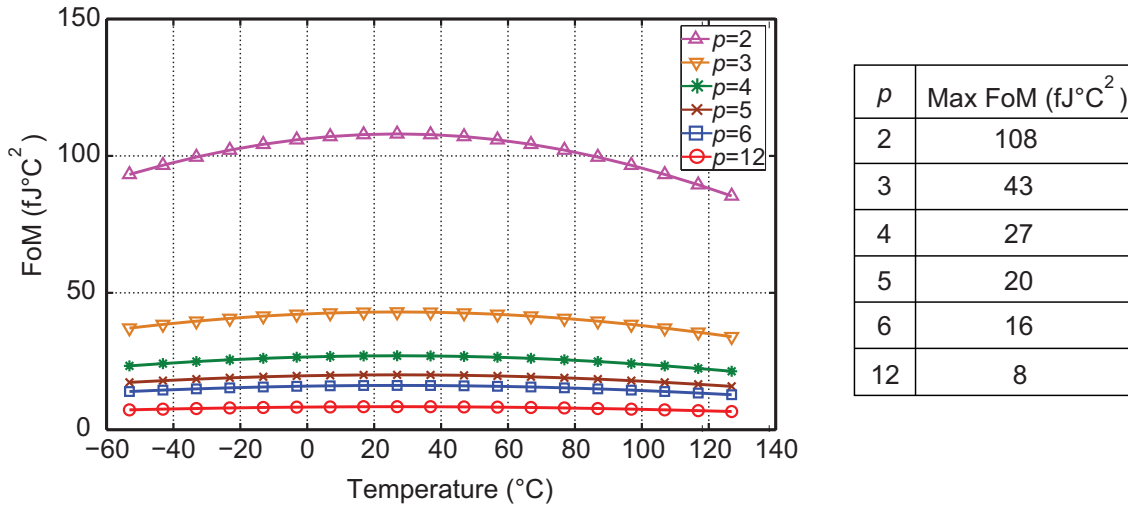


Figure 2.9: Figure-of-merit (FoM) as a function of temperature and for different  $p$  values in the emitter area scaling topology.

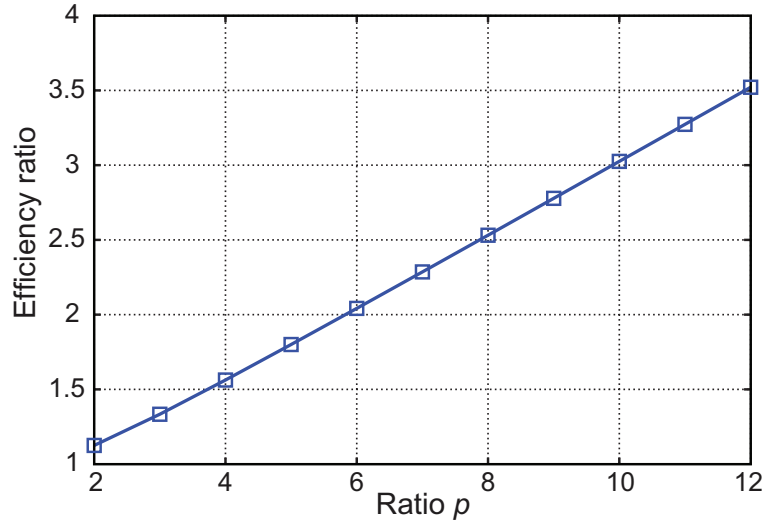


Figure 2.10: Efficiency comparison of bias current versus emitter area scaling as a function of  $p$  values.

with the efficiency of actual temperature sensors.

Figure 2.11 plots the dissipated energy per conversion versus sensor resolution of various smart temperature sensors [2.14]. As shown, the efficiency of smart sensors published prior to the start of this research (in 2009) was limited to about  $1\text{nJ}^\circ\text{C}^2$ , denoted by the solid FoM line. The worst-case theoretical limit of  $121\text{fJ}^\circ\text{C}^2$  (bias current scaling and  $p = 2$ ) has also been projected on the same plot, facilitating the comparison. Apparently, a huge efficiency gap of about four orders of magnitude existed, thus motivating this research. Since then, the work described in this thesis (among others) has improved the energy efficiency of BJT-based temperature sensors by a remarkable two orders of magnitude, represented by the dashed FoM line of  $10\text{pJ}^\circ\text{C}^2$ . One of the designs achieved a resolution FoM of  $11\text{pJ}^\circ\text{C}^2$ , which defined the state-of-the-art when it was published in 2013 [2.15].

Recently, some more energy-efficient designs have been reported: In 2015, a resistor-based sensor was reported with a FoM of  $0.65\text{pJ}^\circ\text{C}^2$  [2.16]. After a single-point calibration, the sensor achieves an inaccuracy of  $\pm 1^\circ\text{C}$  from  $-40^\circ\text{C}$  to  $125^\circ\text{C}$ . In 2014, a BJT-based sensor was reported which obtains a FoM of  $3.2\text{pJ}^\circ\text{C}^2$ , and an inaccuracy of  $\pm 0.15^\circ\text{C}$  using a single-point calibration [2.17]. The sensor employs the concept presented in Figure 2.5. However, its output consists of square-wave whose duty-cycle is a function of temperature. To accurately digitize this information, a counter running at several tens of MHz is required. In [2.17], this was done in an external FPGA, thus excluding this source of extra energy consumption from the FoM calculations.

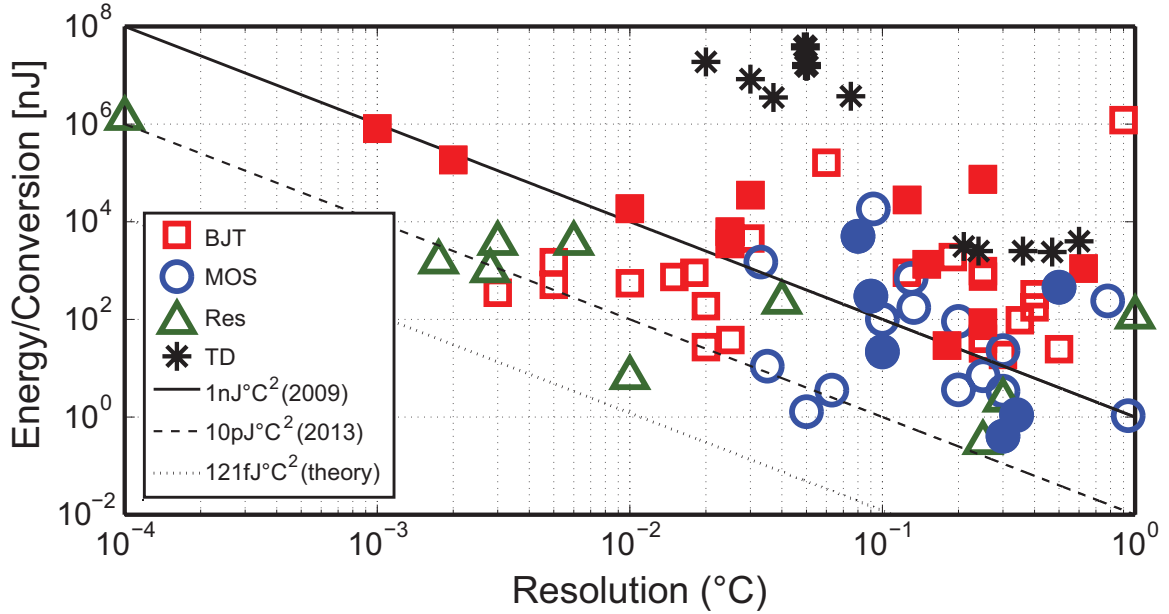


Figure 2.11: Energy per conversion versus resolution for various smart temperature sensors [2.14]. The solid and dashed FoM lines represent the state-of-the-art in 2009 and 2013, respectively. The solid symbols indicate the sensors published before 2009. The theoretical FoM line of a front-end using the (worst-case) bias current scaling ( $p = 2$ ) is illustrated by the dotted line.

### 2.4.2 Energy-Efficiency Gap

As discussed, the motivation for this research was to bridge the four orders of magnitude gap between the theoretical and practical FoMs. As has been shown in the preceding, the energy dissipated by the sensor's front-end is not the limiting factor. So the limiting factor is the sensor's readout circuitry, which dissipates energy without improving the sensor's resolution. In order to minimize the energy-efficiency gap, the dissipated energy as well as the excess noise introduced by readout circuitry should be minimized. Combined with the need to simultaneously achieve high resolution and accuracy, as required in most sensor applications, the result is a non-trivial task, which will be addressed in the following.

Table 2.1 presents an energy breakdown of several BJT-based sensors reported prior to the start of this research. The sensor presented in [2.18] uses  $36\mu\text{A}$  for biasing the BJTs which is less than 3% of the total of  $1.25\text{mA}$ . It achieves a resolution of  $0.25^\circ\text{C}$  in a conversion time of  $30\mu\text{s}$  and using a SAR-ADC, resulting in a FoM of  $5.3\text{nJ}^\circ\text{C}^2$ . In [2.2], which was the most power-efficient precision sensor back in 2009, a front-end based on bias current scaling ( $p = 5$ ) was presented. The sensor consumes  $25\mu\text{A}$ , while the current used for biasing the BJTs is only  $1.5\mu\text{A}$ : about 6% of the total. A  $2^{\text{nd}}$ -order  $\Delta\Sigma$ -ADC has been employed to achieve a resolution of  $0.025^\circ\text{C}$  in  $100\text{ms}$  conversion time, thus a FoM



	JSSC'98 [2.18]	JSSC'96 [2.19]	ISCAS'01 [2.20]	JSSC'05 [2.8]	JSSC'05 [2.1]	ISSCC'09 [2.2]	JSSC'09 [2.21]
<b>BJTs' bias current</b>	36 $\mu$ A	1.1 $\mu$ A	2 $\mu$ A	2 $\mu$ A	6 $\mu$ A	1.5 $\mu$ A	20 $\mu$ A - 100 $\mu$ A
<b>Total current</b>	1.25mA	25 $\mu$ A	100 $\mu$ A	130 $\mu$ A	75 $\mu$ A	25 $\mu$ A	1.52mA
<b>Conversion time</b>	30 $\mu$ sec	20msec	100msec	100msec	100msec	100msec	830 $\mu$ sec
<b>ADC topology</b>	SAR	$\Delta\Sigma$ -ADC	$\Delta\Sigma$ -ADC	$\Delta\Sigma$ -ADC	$\Delta\Sigma$ -ADC	$\Delta\Sigma$ -ADC	$\Delta\Sigma$ -ADC
<b>Resolution</b>	0.25 $^{\circ}$ C	0.625 $^{\circ}$ C	0.125 $^{\circ}$ C	0.03 $^{\circ}$ C	0.01 $^{\circ}$ C	0.025 $^{\circ}$ C	0.15 $^{\circ}$ C
<b>FoM</b>	5.3nJ $^{\circ}$ C <sup>2</sup>	430nJ $^{\circ}$ C <sup>2</sup>	440nJ $^{\circ}$ C <sup>2</sup>	32nJ $^{\circ}$ C <sup>2</sup>	1.9nJ $^{\circ}$ C <sup>2</sup>	3.9nJ $^{\circ}$ C <sup>2</sup>	30nJ $^{\circ}$ C <sup>2</sup>

Table 2.1: Performance summary of prior-art temperature sensors published before 2009.

of 3.9nJ $^{\circ}$ C<sup>2</sup>. A similar trend can be observed in other designs, i.e. the current used to bias BJTs is only a fraction of the total, ranging from 1.5% in [2.7] to 8% in [2.2]. This clearly indicates one major problem: the power consumed in the readout circuitry is by far dominant compared to that in the front-end.

### 2.4.3 ADC Topology

As can be seen from Table 2.1,  $\Delta\Sigma$ -ADCs have been extensively used to read out smart temperature sensors. This is because, unlike Nyquist-rate ADC topologies such as Pipeline, Flash, and Successive Approximation (SAR),  $\Delta\Sigma$ -ADCs have more relaxed matching requirements on the circuit elements and are capable of achieving high resolution by trading it for conversion speed. They are therefore well matched to the low speed (typically less than 10 samples/s) and high resolution (up to 15 bits) requirements of precision smart temperature sensors. In such applications,  $\Delta\Sigma$ -ADCs are usually operated in *incremental* mode; they are first reset at the beginning of a conversion and then operated for a fixed number of cycles [2.22].

Moreover, to mitigate the offset and component mismatch in the front-end e.g. current sources and PNP transistors, which would otherwise increase the sensor's inaccuracy, dynamic error correction techniques, such as chopping and dynamic-element-matching (DEM) are essential. This, in turn, results in dynamic error components in the output of the front-end i.e.  $V_{BE}$  and  $\Delta V_{BE}$ , which then need to be filtered out. The filtering characteristic of low-pass  $\Delta\Sigma$  converters can be used for this purpose, thus making them

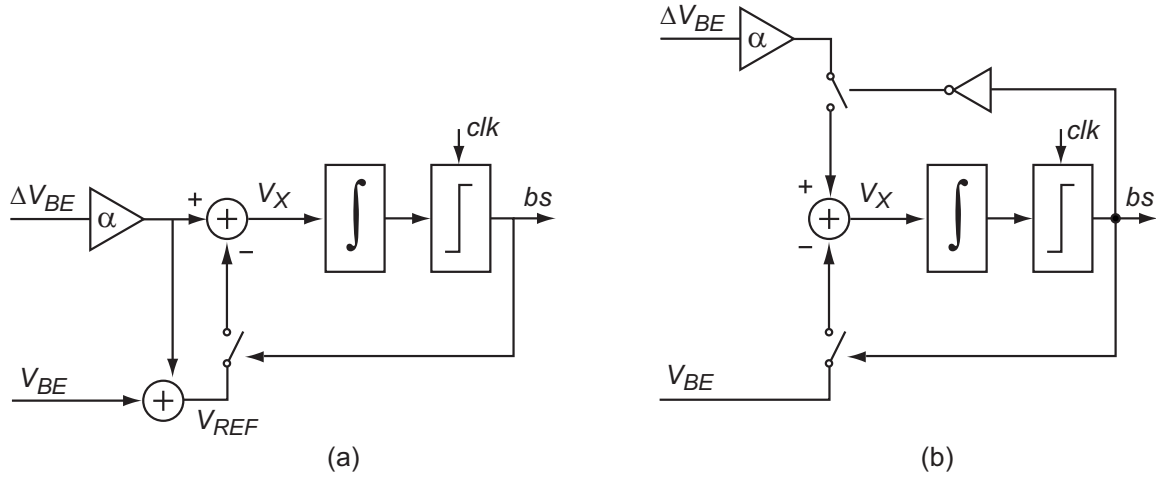


Figure 2.12: The charge balancing in  $\Delta\Sigma$  modulators is used to readout the output of the sensor front-end: (a): the straightforward approach, (b): the optimized approach.

a good choice for sensor applications with stringent accuracy requirements.

The use of  $\Delta\Sigma$  modulators to digitize the front-end's output can be explained with the help of Figure 2.12. As shown in Figure 2.12.a, a straightforward implementation involves generating  $\alpha \cdot \Delta V_{BE}$  and  $V_{REF}$  (see equation (2.5)) and then digitizing their ratio with a  $\Delta\Sigma$  modulator. In this example, a 1<sup>st</sup>-order modulator is formed by an integrator followed by a 1-bit quantizer in a feedback loop. Depending on the quantizer's output bitstream ( $bs$ ), the integrator's input voltage  $V_X$  will be either  $\alpha \cdot \Delta V_{BE}$ , when  $bs = 0$ , or  $-V_{BE}$ , when  $bs = 1$ . The negative feedback effectively tries to balance the charges in the integrator such that the average of the integrator's output voltage tends to zero. This implies that the average value of the quantized signal  $bs$  will track the input signal, and is equal to the desired  $\mu$  (see equation (2.8)). An equivalent charge balancing scheme can be realized with the optimized block diagram of Figure 2.12.b [2.5]. As shown,  $V_{REF}$  doesn't need to be explicitly generated, it can be dynamically created through the charge balancing in the loop, which simplifies the circuit implementation [2.1], [2.2].

For high energy-efficiency, the modulator's output should be limited by thermal rather than quantization noise. Once the modulator is in the thermal noise limited region, its energy-efficiency is theoretically no longer a function of conversion time. Compared to Nyquist-rate ADCs, however,  $\Delta\Sigma$ -ADCs inherently require much longer conversion times to become thermal noise limited. Although acceptable in typical temperature sensing applications, this comes at a cost of substantial energy dissipation per conversion without improving the front-end's energy-efficiency, as seen from equations (2.22) and (2.25). This would imply that the sensor's energy per conversion and thus FoM are substantially dominated by the readout. Using the generic approach with a 1<sup>st</sup>-order  $\Delta\Sigma$ -ADC, requires

at least 32000  $\Delta\Sigma$ -cycles to achieve a target resolution of  $\pm 0.01^\circ\text{C}$ , equivalent to an ENOB of  $\approx 15$  bits. Using higher-order  $\Delta\Sigma$ -ADC topologies significantly reduces the conversion time, e.g. 400 and 64  $\Delta\Sigma$ -cycles for the  $2^{\text{nd}}$ - and  $3^{\text{rd}}$ -order topologies, respectively. The drawback, however, is their greater complexity, reduced stable input range, larger size and higher power consumption. By using the charge-balancing scheme of Figure 2.12.b in a  $2^{\text{nd}}$ -order  $\Delta\Sigma$ -ADC, the sensor in [2.2] achieves a FoM of  $3.9\text{nJ}^\circ\text{C}^2$ . In recent work, this was significantly improved by using an optimized readout and the circuit flexibility offered by NPNs [2.23]. The sensor achieves a resolution of  $0.025^\circ\text{C}$  in a conversion time of 6ms, while drawing  $4.5\mu\text{A}$ , which translates to a FoM of  $24\text{pJ}^\circ\text{C}^2$ . The use of NPNs, however, requires an extra processing step, and therefore, is not well suited for use in low-cost applications.

An energy-efficient alternative to a  $\Delta\Sigma$ -ADC could be a SAR-ADC. This is a simple structure that employs a successive approximation (binary search) algorithm in a feedback loop including a 1-bit quantizer i.e. a comparator. The architecture operates in a “bit-at-a-time” manner, implying that an  $n$  bits conversion takes place within  $n$  steps. Compared to  $\Delta\Sigma$ -ADCs, this is a distinct advantage, resulting in much shorter conversions, e.g.  $30\mu\text{s}$  for 10 bits in [2.18]. The simplicity in the hardware implementation of a SAR-ADC and its capacity for low-power, fast conversions has made it the most attractive choice for applications in which energy-efficiency is a key requirement, e.g. biomedical and wireless sensors [2.24] - [2.27].

However, the resolution of SAR-ADCs depends on component matching, which means that extensive calibration is needed to achieve the minimum ENOB of 13 - 15 bits required by precision temperature sensors. This in turn, rules out their use in low-cost applications. Moreover, the low-pass characteristic of a  $\Delta\Sigma$ -ADC, which is essential to the use of dynamic error correction techniques, is not provided by the SAR-ADC structure. Clearly, a trade-off exists between the high resolution/accuracy provided by  $\Delta\Sigma$ -ADCs and the energy-efficiency of SAR-ADCs.

To summarize, the efficiency gap of sensors listed in Table 2.1 is predominantly due to their inefficient readout circuits. The employed ADCs suffer either from a long conversion time and poor power-efficiency, or are not capable of providing high resolution and/or accuracy without calibration. Design of a low-power, fast, and precision ADC is therefore critical to achieving better energy-efficiency in temperature sensing applications. In the next chapter, a *zoom*-ADC architecture is described which addresses this challenge.

## 2.5 Conclusions

In this chapter, a study of energy-efficiency limits in BJT-based temperature sensors has been presented. A significant gap was observed between the energy-efficiency of prior-art temperature sensors and the theoretical efficiency limits. The employed readout circuits either suffer from long conversion time and poor power-efficiency, or are not capable of providing the target resolution or accuracy. To bridge this efficiency gap, a new readout architecture is clearly required. In the next chapter, a new ADC architecture will be proposed to address this issue.

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## Chapter 3

# Energy-Efficient BJT Readout

### 3.1 Introduction

In analog circuit design, speed and precision are typically achieved at the expense of power dissipation. Therefore, the design of energy-efficient, high-resolution smart temperature sensors, as targeted in this work, is not trivial and involves fundamental trade-offs. Assuming that an ADC's power dissipation scales linearly with sampling frequency, then simply decreasing its conversion time will translate into a proportional increase in power dissipation; i.e. the ADC's energy consumption per conversion remains unchanged. Improving an ADC's energy-efficiency, thus, calls for architecture-level solutions.

The generic BJT readout of Figure 2.4 operates by digitizing  $\Delta V_{BE}$  w.r.t. a stable, temperature independent voltage reference  $V_{REF}$ . A key observation, however, is that all the necessary temperature information is present in the two voltages:  $V_{BE}$  and  $\Delta V_{BE}$ . So, one could alternatively use the ratio of  $V_{BE}$  and  $\Delta V_{BE}$  as a measure of temperature. In [3.1], a BJT-based analog front-end generates  $V_{BE}$  and  $\Delta V_{BE}$  voltages at the input of a  $\Delta\Sigma$ -ADC. The ADC then provides the digital back-end with the non-linear ratio  $\gamma = \Delta V_{BE}/V_{BE}$ . The additional signal processing, which is required to obtain the PTAT output  $\mu$  as in equation (2.8), is carried out by the digital back-end as follows:

$$\mu = \frac{\alpha \cdot \Delta V_{BE}}{V_{BE} + \alpha \cdot \Delta V_{BE}} = \frac{\alpha \cdot \frac{\Delta V_{BE}}{V_{BE}}}{1 + \alpha \cdot \frac{\Delta V_{BE}}{V_{BE}}} = \frac{\alpha \cdot \gamma}{1 + \alpha \cdot \gamma}, \quad (3.1)$$

implying that the scaling factor  $\alpha$  is implemented in the digital back-end. The ratio  $\gamma = \Delta V_{BE}/V_{BE}$  varies between 0.019 and 0.036 over the temperature range from 45°C to 175°C. Therefore, the use of an ADC that directly digitizes  $\gamma$ , as presented in [3.1], results



in a significant waste of dynamic range, much more than the 70% wasted in the generic architecture. This is demonstrated by the poor temperature resolution ( $0.45^\circ\text{C}$ ) and FoM ( $324\text{nJ}^\circ\text{C}^2$ ) achieved in [3.1].

The temperature-to-digital converter (TDC) presented in [3.2] switches the bias current of a diode-connected bipolar transistor from  $I_{bias}$  to  $m \times I_{bias}$  in order to obtain two different base-emitter voltages:  $V_{BE1}$  and  $V_{BE2}$ . These are then digitized by an ADC. The signal processing required to determine the ratio  $\mu$  is then performed in the digital back-end, e.g. the subtraction of  $V_{BE1}$  from  $V_{BE2}$  to obtain the PTAT voltage  $\Delta V_{BE}$  and implementation of the scaling factor  $\alpha$ , etc. Since  $\Delta V_{BE}$  is directly calculated in the digital domain, a high resolution, high accuracy ADC is required to maintain the accuracy of the digitized  $\Delta V_{BE}$  voltage. Such an ADC typically dissipates a high amount of power or requires a long conversion time, both of which would result in poor energy efficiency.

In [3.3], a sensor topology is presented in which a ratio  $\mu' = 2 \cdot \Delta V_{BE} / (V_{BE} + 2 \cdot \Delta V_{BE})$  is digitized by a  $\Delta\Sigma$ -ADC. Since  $\mu'$  is not a linear function of temperature, further processing in the digital back-end is performed in order to calculate the PTAT ratio as follows:  $\mu_{PTAT} = 9 \cdot \mu' / (1 + 8 \cdot \mu')$ . Again, the ADC's input  $2 \cdot \Delta V_{BE}$  is considerably smaller than its reference voltage ( $V_{BE} + 2 \cdot \Delta V_{BE}$ ), which translates to poor usage ( $\sim 10\%$ ) of the ADC's dynamic range. The sensor requires 455ms to achieve a resolution of 30mK and a FoM of  $4.1\text{nJ}^\circ\text{C}^2$ .

In all these examples, a non-linear ratio between  $V_{BE}$  and  $\Delta V_{BE}$  is first digitized by an ADC, while further processing is done in a digital back-end to obtain a linear function of temperature. In this approach, commonly referred to as digitally-assisted analog circuit design, the high-density and low-energy per computation of digital circuits is exploited to enable circuits based on minimal-precision, low-complexity analog building blocks [3.4]. Since digital logic is available in all CMOS processes, this approach is attractive because it does not increase production costs. Some other potential advantages of the digitally-assisted approach are scalability, lower chip area and thus lower cost, simpler architecture and faster post processing.

### 3.2 Proposed Sensor Topology

Let's consider equation (2.8) and reformulate it as follows:

$$\mu = \frac{\alpha \cdot \Delta V_{BE}}{V_{BE} + \alpha \cdot \Delta V_{BE}} = \frac{\alpha}{\frac{V_{BE}}{\Delta V_{BE}} + \alpha}. \quad (3.2)$$

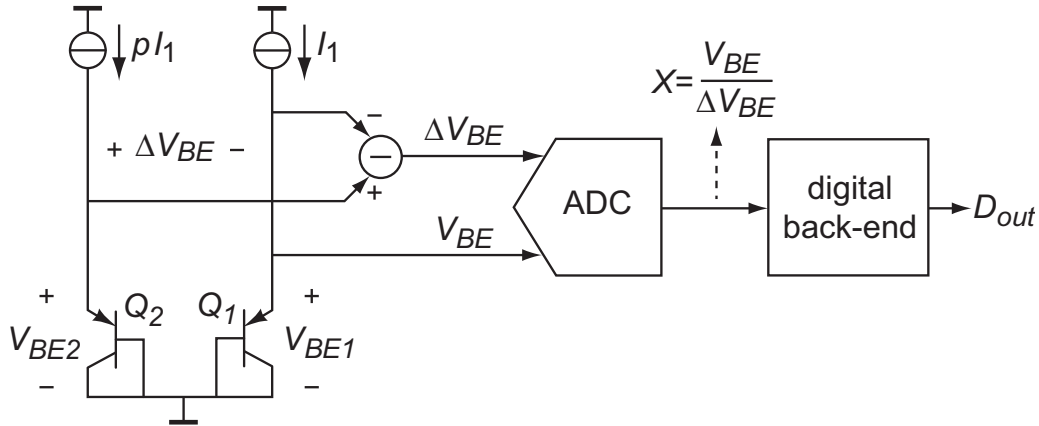


Figure 3.1: Circuit diagram of the proposed temperature sensor based on digitizing  $X = V_{BE}/\Delta V_{BE}$ .

By introducing a new parameter  $X = V_{BE}/\Delta V_{BE}$  the equation can be simplified to:

$$\mu = \frac{\alpha}{X + \alpha}. \quad (3.3)$$

This new interpretation suggests that to obtain the PTAT ratio  $\mu$ , only the ratio  $X$  needs to be digitized. A digitally implemented  $\alpha$  and a simple division in the digital domain can then be used to compute the ratio  $\mu$ .

Figure 3.1 shows the block diagram of the proposed digitally-assisted bandgap temperature sensor. Compared to the block diagram of the generic bandgap temperature sensor, depicted in Figure 2.4, the ADC only has to compute the ratio between  $V_{BE}$  and  $\Delta V_{BE}$  voltages. Therefore, the two voltages no longer need to be combined to generate  $V_{REF}$ . This is one step toward simplifying the system; the temperature-independent reference voltage  $V_{REF}$  is replaced by a temperature varying voltage, i.e.  $V_{BE}$ . The ratio  $X = V_{BE}/\Delta V_{BE}$ , will then depend on both temperature and the collector current density ratio  $p$  used to develop  $\Delta V_{BE}$ . Increasing temperature or the ratio  $p$  results in a larger PTAT voltage  $\Delta V_{BE}$ , which in turn, reduces the ratio  $X$  in a non-linear fashion.

Figure 3.2 illustrates the ratio  $X$  as a function of temperature and for different bias current density ratios  $p = 3, 5, 8,$  and  $12$ . As depicted,  $X$  is a non-linear function of temperature and decreases as  $p$  increases. A different point of view, depicted in Figure 3.3, compares the base-emitter voltage  $V_{BE}$  with integer multiples of  $\Delta V_{BE}$  for the case when  $p = 5$ . As shown, and as is also evident from Figure 3.2,  $V_{BE}$  ranges between  $6 \cdot \Delta V_{BE}$  and  $28 \cdot \Delta V_{BE}$  from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ , and is about  $14.4 \cdot \Delta V_{BE}$  at room temperature. Linearization of the ratio  $X$  can then be performed in the digital domain as described in equation (3.3) and graphically depicted in Figure 3.4.

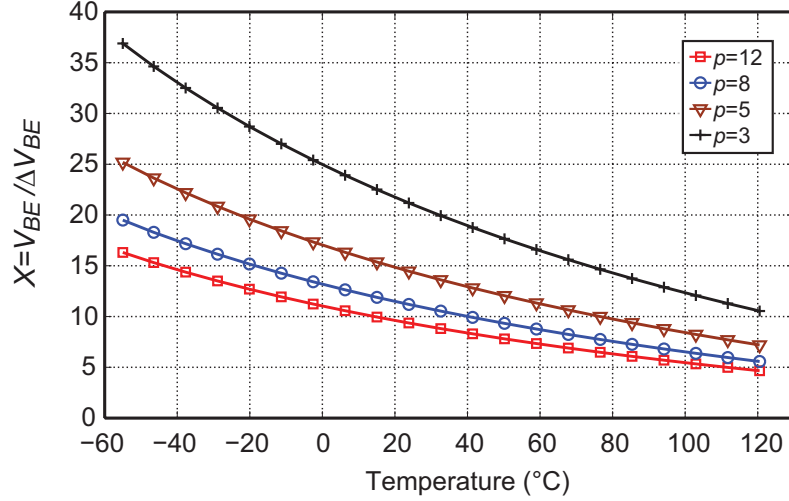


Figure 3.2: Temperature dependence of the ratio  $X$  for different bias current ratios  $p = 3, 5, 8,$  and  $12$ .

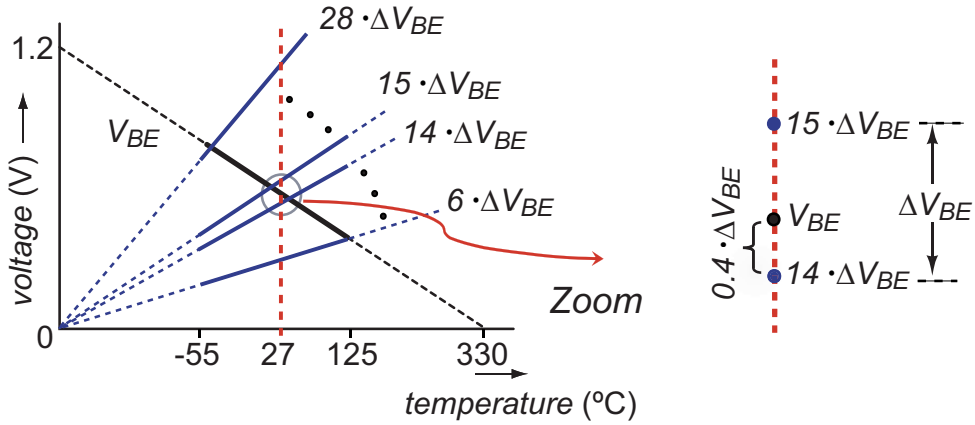


Figure 3.3: Temperature dependence of  $V_{BE}$  and integer multiples of  $\Delta V_{BE}$  from  $-55^{\circ}\text{C}$  to  $125^{\circ}\text{C}$ , e.g. at room temperature:  $14 \cdot \Delta V_{BE} < V_{BE} < 15 \cdot \Delta V_{BE}$ .

### 3.2.1 ADC's Resolution Requirement

For the proposed digitally-assisted topology, deriving the resolution requirement on the ADC requires an analysis of the non-linear ratio  $X = V_{BE}/\Delta V_{BE}$ . Considering equation (3.3), the derivative of ratio  $\mu$  with respect to  $X$  can be described as:

$$\frac{\partial \mu}{\partial X} = \frac{-\alpha}{(X + \alpha)^2}. \quad (3.4)$$

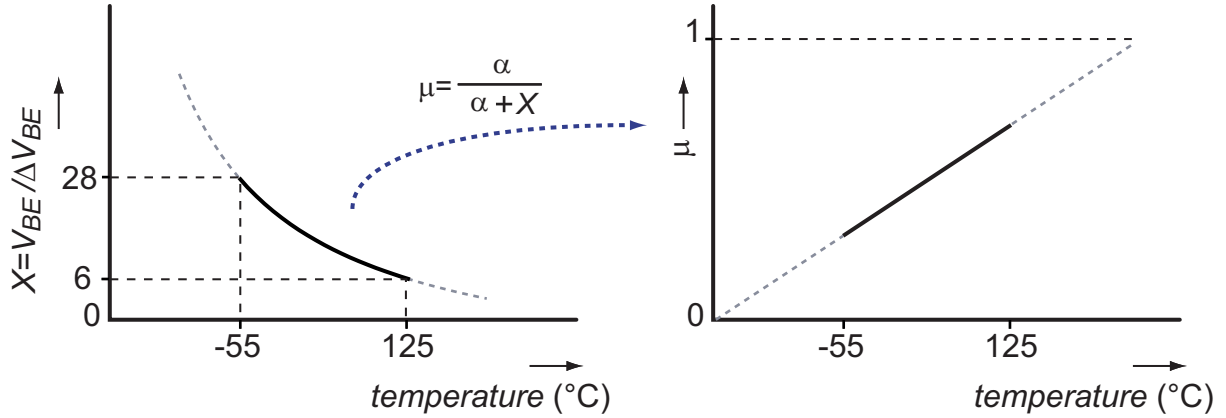


Figure 3.4: Non-linear  $X = V_{BE} / \Delta V_{BE}$  ( $p = 5$ ) and linearized  $\mu = \alpha / (\alpha + X)$  as a function of temperature.

The minimum detectable step  $\Delta X$ , can thus be found as:

$$\Delta\mu = \Delta X \cdot \frac{\partial\mu}{\partial X} \Rightarrow \Delta X = \frac{\Delta\mu}{\left| \frac{\partial\mu}{\partial X} \right|}, \quad (3.5)$$

implying a temperature-dependent resolution due to the temperature dependence of the derivative  $\partial\mu/\partial X$ . Moreover, for a given temperature resolution, the required resolution will depend on the coefficient  $\alpha$ , and thus on the current density ratio  $p$ , as is evident from equation (2.7).

The ADC's resolution can be calculated as follows:

$$ENOB \approx \log_2 \left( \frac{X_{FS}}{\Delta X} \right) - 1, \quad (3.6)$$

where  $X_{FS}$  is the full scale i.e.  $X_{FS} = X_{max} - X_{min} = 22$ , for  $p = 5$ . For the target sensor resolution of  $\pm 0.01^\circ\text{C}$ , the ADC's required resolution as a function of temperature is shown in Figure 3.5. As shown, the resolution requirement is slightly temperature dependent, with a maximum value of 14.3 bits at  $125^\circ\text{C}$ , which is about 0.5 bits less than that of the generic approach. The corresponding impact on conversion time and power consumption will depend on the chosen ADC architecture, which is the subject of the following discussion.

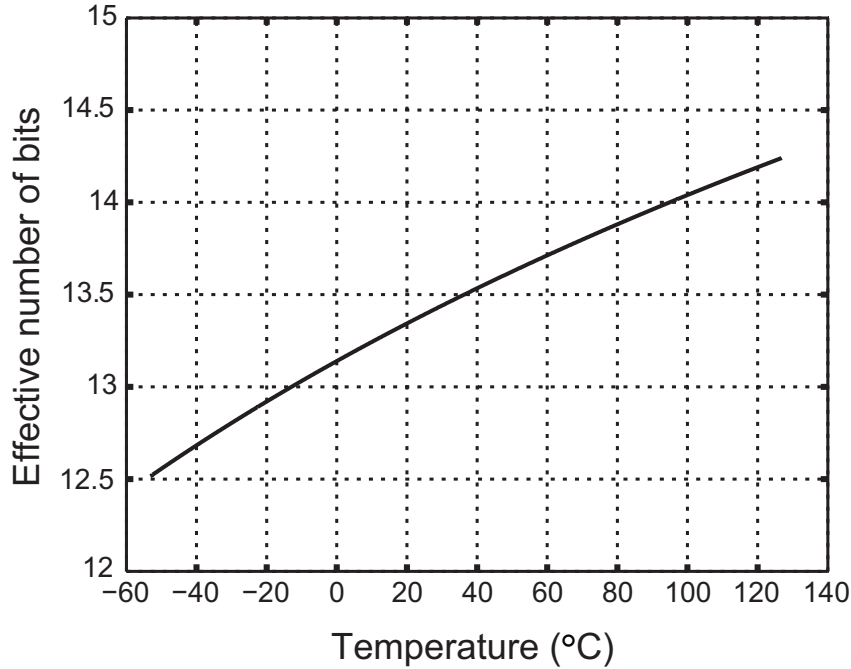


Figure 3.5: ADC’s resolution requirement versus temperature in the proposed readout approach (assuming  $p = 5$ ) and for a target temperature resolution of  $\pm 0.01^\circ\text{C}$ .

### 3.3 The Zoom-ADC: An Energy-Efficient ADC

#### 3.3.1 Introduction

In general, temperature changes are rather slow. Therefore, the ratio  $X$  can be accurately digitized by a two-step ADC, in which a full-range conversion first obtains a coarse estimate of the input level. This is followed by a low-range, but high resolution, fine conversion to obtain an accurate estimate of the input level. In the proposed zoom-ADC, the strengths of both SAR-ADC and  $\Delta\Sigma$  converters are combined into a two-step conversion scheme, as will be explained in the following.

#### 3.3.2 Topology

In the proposed sensor topology, the input to the ADC,  $X = V_{BE} / \Delta V_{BE}$ , is always greater than one, and thus can be expressed in terms of integer and fractional parts, as

follows:

$$X = \frac{V_{BE}}{\Delta V_{BE}} = n + \mu', \quad (3.7)$$

where  $n$  is the integer part of  $X$ , and  $\mu'$  is its fractional part i.e.  $0 < \mu' < 1$ . As we observed in Figure 3.2, the magnitude of  $X$ , and thus the integer value  $n$ , depends, apart from temperature, on the collector current density ratio  $p$  used to generate  $\Delta V_{BE}$ . In the rest of this chapter we assume  $p = 5$ . This implies that  $X$  ranges from 6 to 28 over the military temperature range, i.e. from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ .

To determine  $X$ , the two parameters  $n$  and  $\mu'$  are determined in a two-step manner, as shown in Figure 3.6; In the first step the integer  $n$  is found in a coarse conversion phase (Figure 3.6.a). Knowing  $n$ , the fraction  $\mu'$  can then be readily resolved in the succeeding fine conversion step by zooming into the range from  $n$  to  $n+1$ , as graphically depicted in Figure 3.6.b. Since the full scale range of the fine conversion is now quite small, the resolution requirements of the fine converter are greatly relaxed, thus leading to a simple

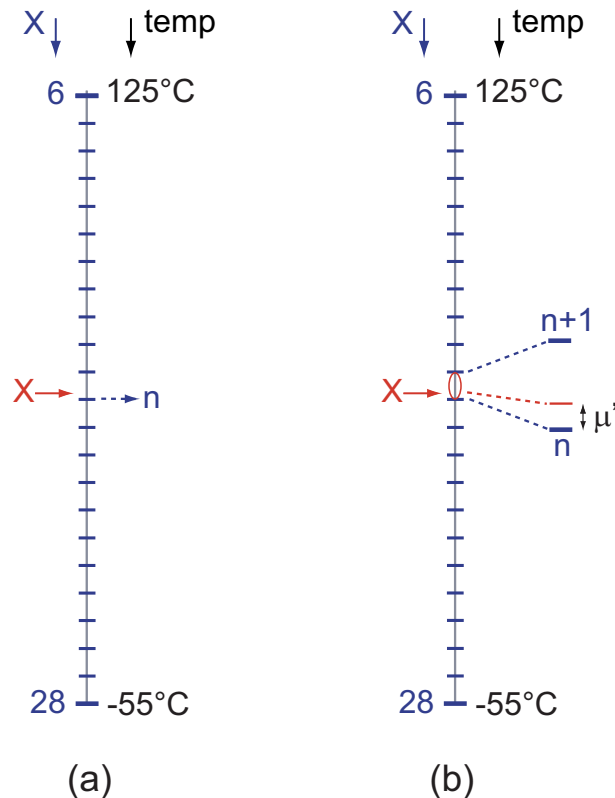


Figure 3.6: Temperature dependence of  $X = n + \mu'$  from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ . The integer  $n$  ranges between 6 and 28 (a), while the fraction  $\mu'$  ranges between 0 and 1 (b).

implementation and short conversion times, concurrently. Moreover, in contrast to the previously described readout architectures, the dynamic range of the zoom-ADC is now fully utilized, thus enabling energy-efficient conversions.

### 3.3.3 Coarse Converter

In order to determine the integer value  $n$ , a low-resolution converter is sufficient. As shown in Figure 3.2, for practical  $p$  ratios a resolution of about  $5 \sim 7$  bits in the coarse converter is sufficient to find the integer  $n$ . As explained before, this can be quickly done by means of a SAR-ADC. This combines a simple hardware implementation with low-power operation, both of which are key requirements for our target application. Assuming  $p = 5$  implies  $6 \leq n \leq 28$ , and thus a 5-bit SAR-ADC.

Figure 3.7 shows the proposed block diagram of the coarse converter. In a feedback loop, the base-emitter voltage  $V_{BE}$  is compared to integer multiples of  $\Delta V_{BE}$  i.e.  $k \cdot \Delta V_{BE}$ , where  $k = 1, 2, \dots, 28$ . A clocked comparator then performs the comparison by simply detecting the sign of:

$$V_X = V_{BE} - k \cdot \Delta V_{BE}. \quad (3.8)$$

The comparison result  $bs$  is then used by the SAR-logic, which, in turn, properly updates the multiplication factor  $k$  for the next comparison step. The procedure will be continued until the SAR-logic finds the region in which:

$$n \cdot \Delta V_{BE} < V_{BE} < (n + 1) \cdot \Delta V_{BE}. \quad (3.9)$$

At room temperature, this condition is satisfied for  $n = 14$  (see Figure 3.3). Since

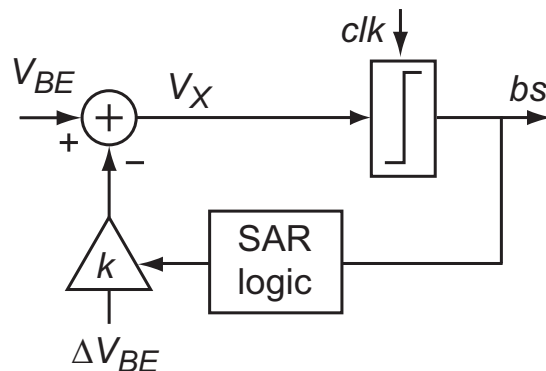


Figure 3.7: Block diagram of the zoom-ADC during the coarse conversion.

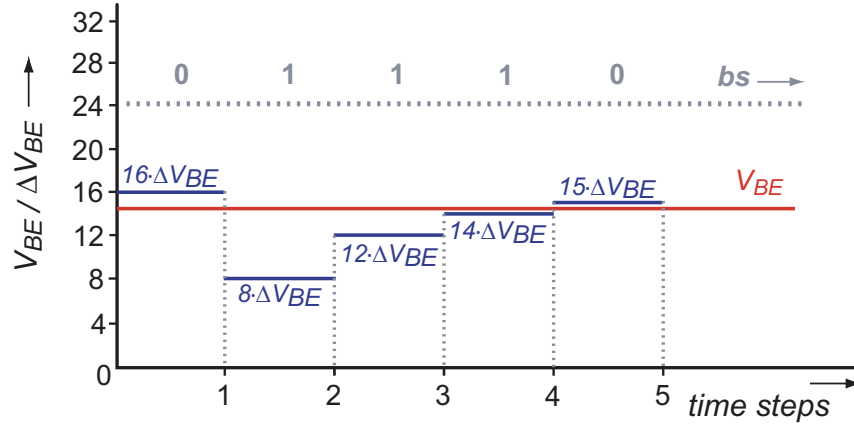


Figure 3.8: The successive approximation steps during the coarse conversion. The unknown  $n$  is determined within 5 approximation steps.

$n \leq 28$ , the SAR-logic only needs five steps to find  $n$ . Figure 3.8 shows the comparison steps performed by SAR-logic to find  $n$  at room temperature.

An alternative approach to further simplify the coarse conversion would be to replace the SAR-logic with a linear search algorithm as follows: the control logic sets  $k = 6$ , which corresponds to the high end of temperature range i.e.  $125^\circ\text{C}$ . Since  $V_{BE} < 6 \cdot \Delta V_{BE}$  over the temperature range from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ , the comparison result will be always 1. The control logic then starts ramping up the comparison level by 1 i.e.  $k = 7, 8, \dots$  until it detects a zero at the comparator's output implying that  $V_{BE} < (n+1) \cdot \Delta V_{BE}$ . At room temperature, 10 comparison steps are required, while at  $-55^\circ\text{C}$  the maximum number of 23 steps are required. Although this approach requires simpler control logic, the price to pay is a longer (and temperature dependent) conversion time in the coarse phase, e.g.  $1 \sim 23$  versus 5 comparison steps for the case of a SAR-logic.

### 3.3.4 Fine Converter

Having found the integer  $n$ , the next step is to determine the fraction  $\mu'$ . By rewriting equation (3.3), this can be expressed as follows:

$$\mu' = \frac{V_{BE} - n \cdot \Delta V_{BE}}{\Delta V_{BE}}. \quad (3.10)$$

In other words, the fraction  $\mu'$  can be found by resolving the *residue* ( $V_{BE} - n \cdot \Delta V_{BE}$ ) over the range  $\Delta V_{BE}$ . This can be done by using a  $\Delta\Sigma$ -ADC, whose input and reference



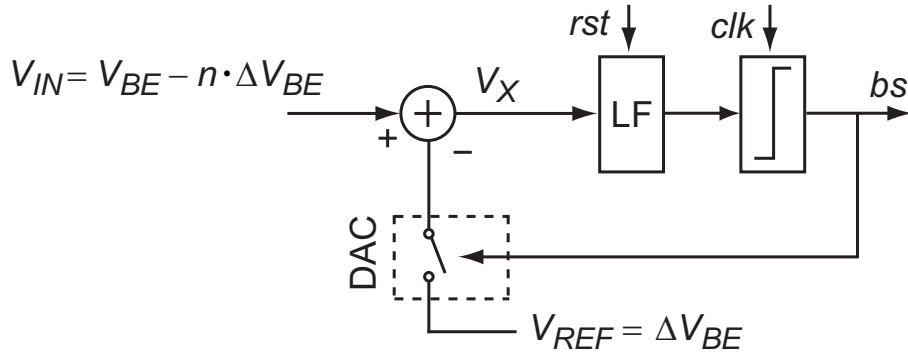


Figure 3.9: Block diagram of the zoom-ADC during the fine conversion.

voltages are set to  $(V_{BE} - n \cdot \Delta V_{BE})$  and  $\Delta V_{BE}$ , respectively, as shown in Figure 3.9. It consists of a summation node, an integrator, a 1-bit quantizer (a clocked comparator), and a 1-bit digital-to-analog converter (DAC). The sign of the integrator's output,  $V_{INT}$ , is determined on the rising edges of the clock. The comparison result  $bs$ , is then fed back via the DAC, which determines whether  $\Delta V_{BE}$  has to be subtracted from the input  $V_{IN} = V_{BE} - n \cdot \Delta V_{BE}$  or not. Therefore, the integrator's input  $V_X$  can be expressed as:

$$V_X = \begin{cases} V_{BE} - n \cdot \Delta V_{BE} & \text{if } bs = 0 \\ V_{BE} - (n + 1) \cdot \Delta V_{BE} & \text{if } bs = 1. \end{cases} \quad (3.11)$$

A loop filter (LF) then processes the difference between the input voltage  $V_{IN}$  and the feedback voltage. Since the net integrated charge is forced to be approximately zero by the feedback loop, the bitstream average is the desired fraction  $\mu' = (V_{BE} - n \cdot \Delta V_{BE}) / \Delta V_{BE}$ . In practice, a decimation filter processes  $bs$  to calculate  $\mu'$ . Depending on the order of loop filter, different types of decimation filters could be used to optimally decimate the produced bitstream  $bs$ . A discussion on the different types of decimation filters and their pros and cons is, however, out of the scope of this thesis, and can be found in [3.5].

A more elegant way to find the fraction  $\mu'$  is shown in the block diagram of Figure 3.10. This topology is equivalent to the one in Figure 3.9, but results in a much simpler circuit-level implementation. Depending on the bitstream value  $bs$ , either  $n \cdot \Delta V_{BE}$  or  $(n + 1) \cdot \Delta V_{BE}$  will be subtracted from the input  $V_{BE}$ , thus simply implementing equation (3.10). Likewise,  $(V_{BE} - n \cdot \Delta V_{BE})$  is integrated when  $bs = 0$ . When  $bs = 1$  the input to the integrator  $V_X$  will be reduced by  $\Delta V_{BE}$  i.e.  $V_X = V_{BE} - (n + 1) \cdot \Delta V_{BE}$ , and thereby reduces the accumulated value. Similarly the charge balancing forces the average charge accumulated in the integrator to be (approximately) zero, and hence equation (3.10) holds.

Comparing the block diagrams of the coarse and fine converters shown in Figure 3.10 and Figure 3.7 reveals the similarities between the two converters; they both sample  $(V_{BE}$

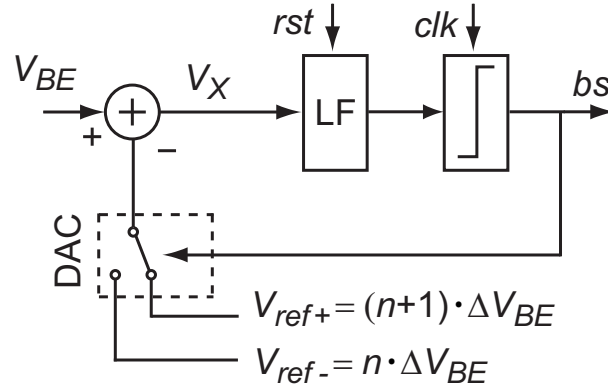


Figure 3.10: An equivalent block diagram to perform the fine conversion.

$-k \cdot \Delta V_{BE}$ ), and use a 1-bit quantizer i.e. a comparator. Sharing hardware between the two converters enables a compact implementation, as will be discussed in the following chapters.

A key observation is that the range of the  $\Delta\Sigma$ -ADC employed in the fine conversion is reduced to  $\Delta V_{BE}$ , which is  $\approx 40\text{mV}$  at room temperature ( $p = 5$ ). Compared to the bandgap voltage reference of  $\approx 1.2\text{V}$  used in conventional approaches, this results in significantly shorter and thus more energy-efficient conversions. As a result of the reduced full scale range in the fine step, the error signals processed by the loop filter have small amplitudes, which can be readily handled by low-power, low-swing integrators. Furthermore, since the fine conversion step incorporates a  $\Delta\Sigma$ -ADC, it leverages the high resolution capability of  $\Delta\Sigma$ -ADCs without relying on matching of circuit elements. Moreover, dynamic techniques such as chopping, auto-zeroing (AZ), and dynamic element matching (DEM) [3.5], [3.6] can still be used in the zoom-ADC, to achieve the precision required in temperature sensors. In the following, a more detailed system-level analysis of the proposed zoom-ADC is presented.

### 3.3.5 System-Level Considerations

In this section the system-level design of a zoom-ADC is addressed, together with some considerations for its practical implementation.

#### 3.3.5.1 Redundancy and Guard-Banding

In practice, the quantization errors of the coarse converter, which are caused by the comparator's input-referred noise and offset as well as by mismatch in the feedback DAC,

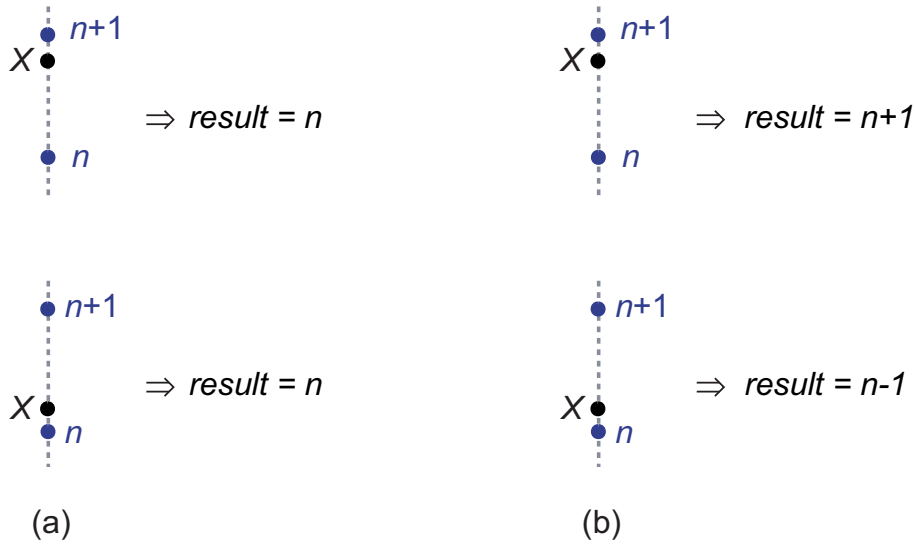


Figure 3.11: Accuracy in coarse step: Ideal (left) and practical (right) situations.

could result in incorrect  $n$  values, as shown in Figure 3.11. This is especially problematic when  $X \approx n$  or  $(n+1)$ . As shown, the coarse result could then be either  $(n-1)$  or  $(n+1)$ , when  $X \approx n$  or  $X \approx (n+1)$ , respectively. The fine conversion will then zoom into the wrong region, resulting in clipping or out-of-ranging, both of which would result in significant errors at the ADC's output, up to 1-LSB of the coarse conversion.

To avoid such out of range errors, redundancy, which is often employed in two-step ADCs [3.7], [3.8], is also employed in the zoom-ADC to accommodate the SAR-ADC's quantization errors and relax its required accuracy. This is done by extending the fine conversion range and making it equal to 2-LSBs of the coarse conversion, i.e.  $2 \cdot \Delta V_{BE}$ . This redundancy ensures that the input  $X$  will always lie safely within the input range of the fine  $\Delta\Sigma$ -ADC. Since extending the range comes at the price of reduced resolution in the  $\Delta\Sigma$  phase, the number of cycles should be appropriately increased to maintain the target ENOB.

The overlap between the two conversion steps is readily implemented by setting the feedback reference range of the  $\Delta\Sigma$  modulator to  $2 \cdot \Delta V_{BE}$ . However, ensuring that the input  $X$  is always within the extended range requires an extra comparison step. As shown in Figure 3.12, this involves comparing  $V_{BE}$  to  $(n+0.5) \cdot \Delta V_{BE}$  at the end of the coarse conversion, and is referred to as the “guard-banding” step throughout this thesis. The total number of coarse comparison steps, therefore, increases by one. Depending on the result, the references of the fine converter can then be set to cover the range from  $n$  to  $(n+2)$  or from  $(n-1)$  to  $(n+1)$ , for  $bs = 1$  and  $bs = 0$ , respectively. This can be implemented by modifying the block diagram of Figure 3.10 so as to set the feedback references to  $(n-1) \cdot \Delta V_{BE}$  and  $(n+1) \cdot \Delta V_{BE}$  or to  $n \cdot \Delta V_{BE}$  and  $(n+2) \cdot \Delta V_{BE}$ . Assuming a

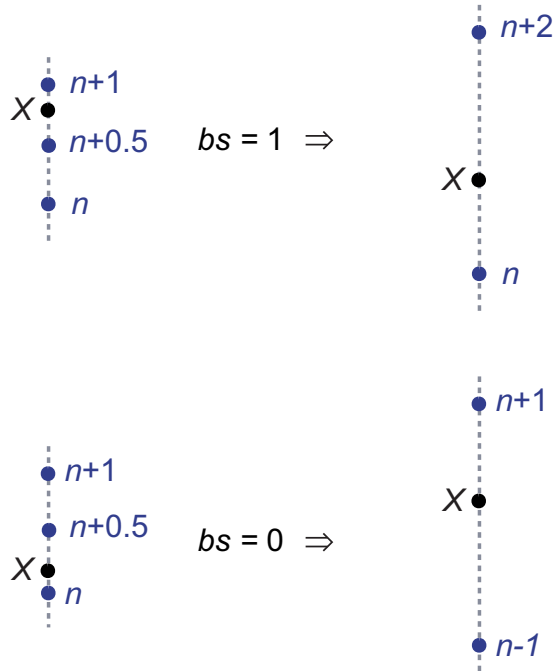


Figure 3.12: Guard-banding mechanism ensures that  $X$  is always roughly in the middle of the extended range, thus avoiding out-of-ranging.

5-bit coarse conversion, and a fine input range of  $2 \cdot \Delta V_{BE}$ , the error in the coarse ADC should be less than  $\pm 0.5 \cdot \Delta V_{BE}$ , which corresponds to a worst case error of  $\pm 15\text{mV}$  at  $-55^\circ\text{C}$ .

By combining the results from the two conversion steps, the digital output  $X_{out}$  of the zoom-ADC can then be expressed as follows:

$$X_{out} = n' + 2 \cdot \mu'', \quad (3.12)$$

where  $n'$  and  $\mu''$  are the results of the coarse and fine steps after guard-banding step, i.e.  $n'$  is equal to either  $(n-1)$  or  $n$ , while  $0 < \mu'' < 1$  is the resolved residue over the extended range of  $2 \cdot \Delta V_{BE}$ , and thus should be scaled before adding to  $n'$ .

### 3.3.5.2 Number of Cycles

When operated in incremental mode, the achievable ENOB of a  $\Delta\Sigma$  modulator depends on three parameters: the order of the loop filter, the quantizer resolution, and the number of clock cycles  $N_{\Delta\Sigma}$ . Due to its ideal linearity, a 1-bit quantizer is best suited for precision applications. It is also power and area efficient, since it only requires a single comparator.

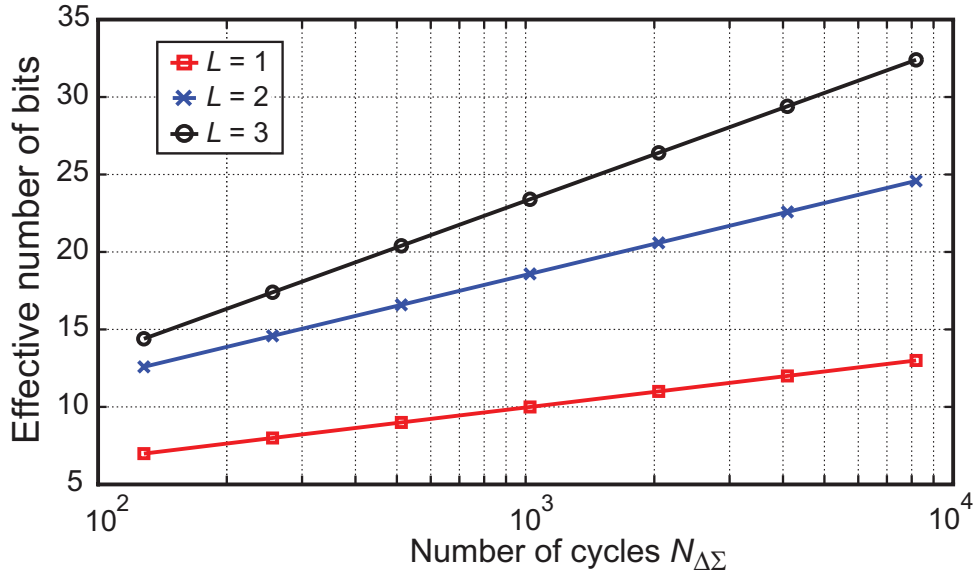


Figure 3.13: ENOB versus number of cycles  $N_{\Delta\Sigma}$  for conventional incremental  $\Delta\Sigma$ -ADCs.  $L$  is the loop filter's order.

To illustrate the modulator's design space, Figure 3.13 shows the simulated ENOB of a 1-bit incremental  $\Delta\Sigma$  modulator versus the number of cycles  $N_{\Delta\Sigma}$  for various modulator orders ( $L$ ). It can be seen that for the same ENOB,  $N_{\Delta\Sigma}$  can be reduced by using higher order loop filters. However, this is at the expense of more complexity and power consumption, as well as a reduction in the range of input signals for which the modulator is stable. For example, the usable input range of a 3<sup>rd</sup> order  $\Delta\Sigma$  modulator is only about 67% of its reference range [3.6].

The maximum ENOB obtainable from the 1<sup>st</sup>- and 2<sup>nd</sup>-order zoom-ADCs versus  $N_{\Delta\Sigma}$  at 25°C is illustrated in Figure 3.14, assuming a coarse ADC resolution of 5 bits. In order to achieve the target ENOB of 13.4 bits (to allow a temperature resolution of  $\pm 0.01^\circ\text{C}$  at 25°C, as in Figure 3.5), a 1<sup>st</sup>-order zoom-ADC would theoretically require  $\approx 680$  cycles. The use of a 2<sup>nd</sup>-order zoom-ADC reduces the required number of cycles to  $\approx 43$ . Considering the increased range due to guard-banding, these numbers will increase to about 1360 and 60 for the 1<sup>st</sup>- and 2<sup>nd</sup>-order zoom-ADCs, respectively. However, for high energy-efficiency, the modulator's output should be limited by thermal rather than quantization noise. This translates to a slight increase in the number of cycles shown in Figure 3.14. Moreover, the full scale range of the fine converter is temperature dependent ( $2 \cdot \Delta V_{BE}$ ), as is its quantization error. On the other hand, due to the non-linear ENOB requirements versus temperature, higher resolution is required at higher temperatures (see Figure 3.5). This again translates into an increase in the number of cycles in order to meet the target resolution over the whole temperature range.

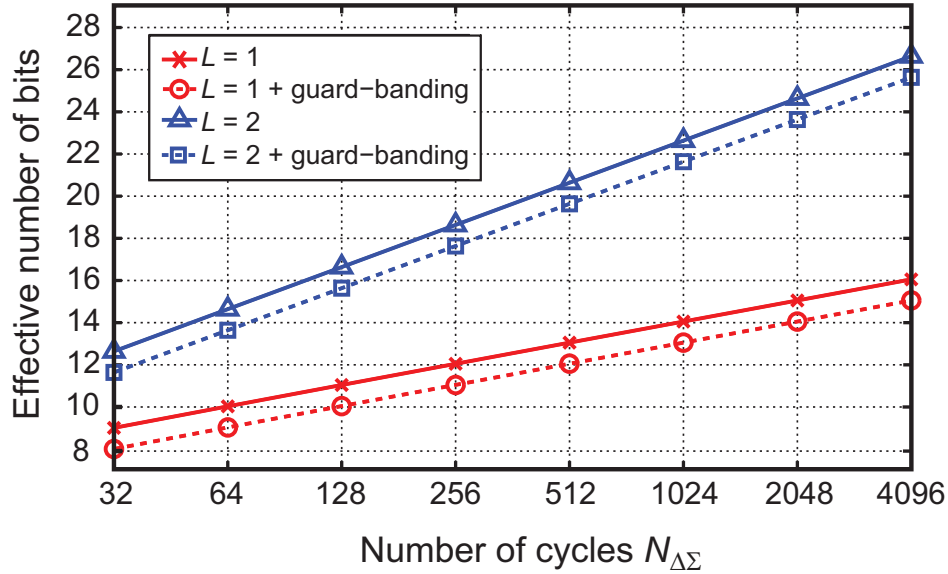


Figure 3.14: ENOB versus number of cycles  $N_{\Delta\Sigma}$  (at 25°C) for the 1<sup>st</sup> and 2<sup>nd</sup>-order zoom-ADC employing a 5-bit coarse converter.

### 3.3.5.3 Signal Swing

The loop filter of a  $\Delta\Sigma$  modulator consists of a number of integrator blocks, each of which will be implemented by an amplifier. The first integrator processes the error signal i.e. the difference between the input and reference signals, whose amplitudes define the integrator's output swing. In practice, the output swing defines the amplifier's topology, and eventually its power efficiency. In conventional single-bit  $\Delta\Sigma$  modulators, the output of the first integrator always includes a large input-signal related component, which results in poor power efficiency. In the proposed zoom-ADC, this issue is mitigated because the error signal processed by the loop filter of the fine  $\Delta\Sigma$ -ADC is proportional to the quantization error of the coarse conversion step, and is thus significantly smaller.

To verify this, a behavioral model for a 1<sup>st</sup>-order zoom-ADC was developed. Figure 3.15 shows a simplified block diagram of this model, in which various design parameters such as the current ratio  $p$ , the resolution of the coarse converter, the mismatch parameters of the feedback loop, the DC gain of the integrator, etc. can be adjusted. Using this model, Figure 3.16 shows the maximum swing of the integrator's output as a function of the coarse ADC's resolution at 25°C. As shown, the output swing is inversely proportional to the resolution of the coarse ADC. For a 5-bit coarse ADC and assuming a voltage gain of 0.5x, an output swing of only  $\pm 50\text{mV}$  is required. Clearly, guard-banding effectively reduces the resolution of coarse converter by 1 bits, resulting in larger error

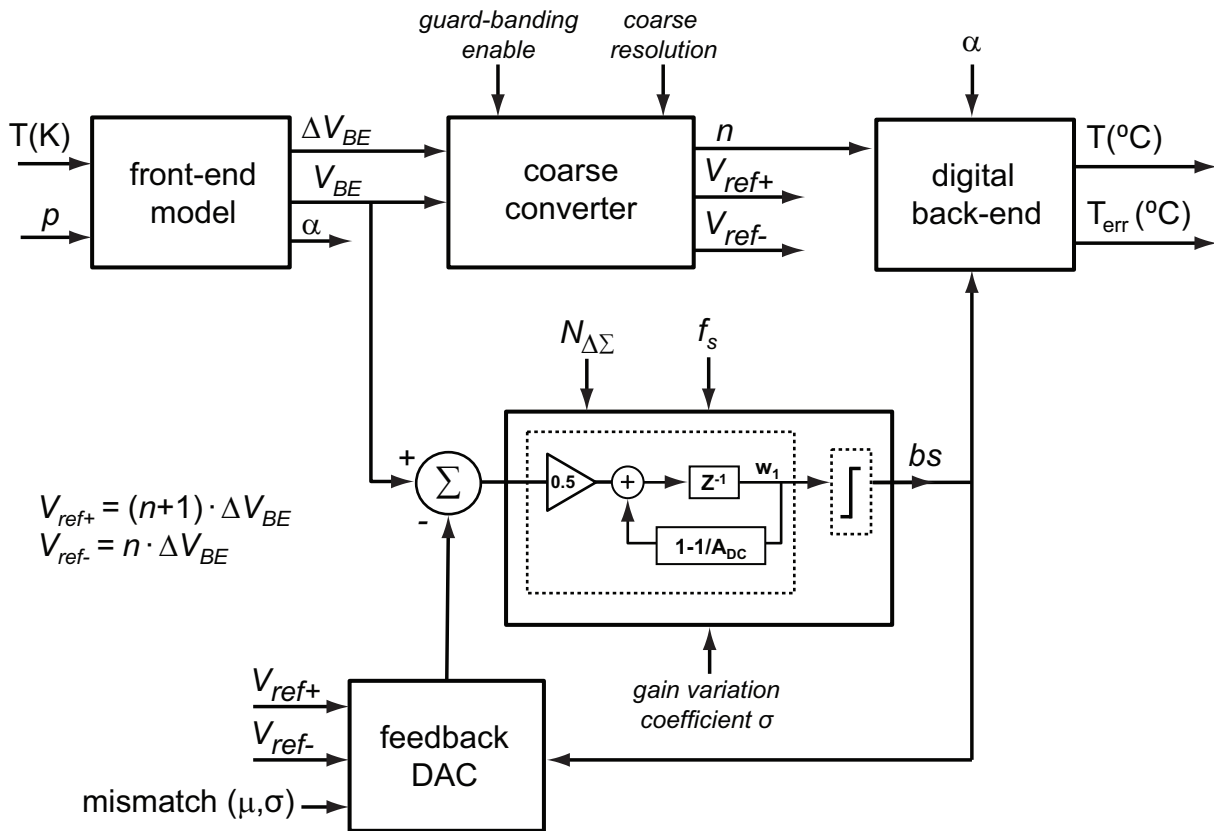


Figure 3.15: Simplified block diagram of the behavioral model of the zoom-ADC.

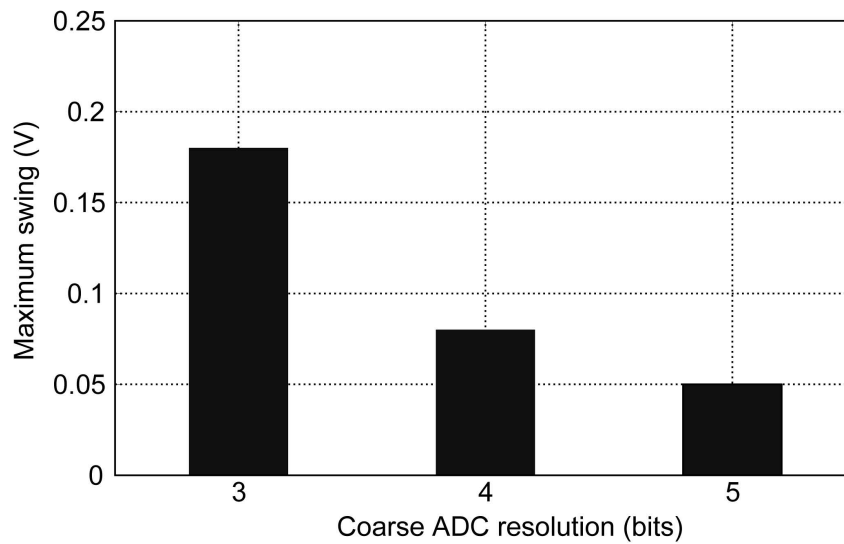


Figure 3.16: Maximum signal swing of the fine  $\Delta\Sigma$ -ADC's first integrator as a function of the coarse step resolution (a voltage gain of 0.5x was assumed).

signals and slightly increasing the output swing.

A similar model can be developed for the conventional charge balancing configuration (see Figure 2.12) [3.9], [3.10]. The result is that the first integrator has to handle an output swing of about  $\pm 300\text{mV}$  for a similar voltage gain of 0.5x. Apparently, zooming significantly relaxes the opamp's output swing and settling requirements, and thus enables a simple and power-efficient realization, as will be shown in the following chapters.

### 3.3.5.4 Integrator Gain

The finite DC gain and gain variation of the integrators in the loop filter are critical error sources. Due to the finite DC gain of the amplifier used in the integrator, the charge on the integration capacitors will leak away. The impact caused by such a leaky integrator is non-linearity, which reduces the achievable effective number of bits (ENOB). An amplifier's DC gain can usually be approximated by a third-order polynomial [3.11] as follows:

$$A_{DC}(V_{out}) \approx A_{DC} \cdot \left( 1 - \sigma \left| \frac{V_{out}}{V_{max}} \right|^3 \right), \quad (3.13)$$

where  $A_{DC}$  is the DC gain at the mid-level output,  $\sigma$  is the gain variation coefficient,  $V_{out}$  is the output swing, and  $V_{max}$  is the maximum output swing. The model in Figure 3.15 has been used to verify the impact of such variations on the zoom-ADC's performance. Figure 3.17 shows the simulated error in degrees Celsius due to the finite DC gain and gain variation of the integrator of a 1<sup>st</sup>-order zoom-ADC, assuming a 5-bit coarse conversion. In this plot, the impact of gain variations on the coarse converter is neglected, since in this step the integrator only acts as a sample-and-hold amplifier for the quantizer, and thus for practical DC gain values, its gain variations will not impact the coarse conversion result. As shown, a DC gain of 80dB is required to ensure a temperature error below  $\pm 10\text{mK}$ , assuming  $\sigma = 0.9$ , which represents a significant gain variation coefficient. In order to achieve similar performance with the conventional charge balancing scheme used in [3.9], [3.10], a minimum DC gain of 100dB would be required, which is considerably higher and more difficult to achieve, especially at low supply voltages in modern processes, e.g. 1.8V in the target 160nm CMOS and sub-1V in nano-scale processes.

Given the reduced output signal swing, the OTA of a zoom-ADC does not slew and so its settling follows a single pole response. In this case, incomplete settling only results in a fixed gain error in the integrator. As long as this error does not significantly alter the loop filter's transfer function, the modulator's performance will not be impaired.



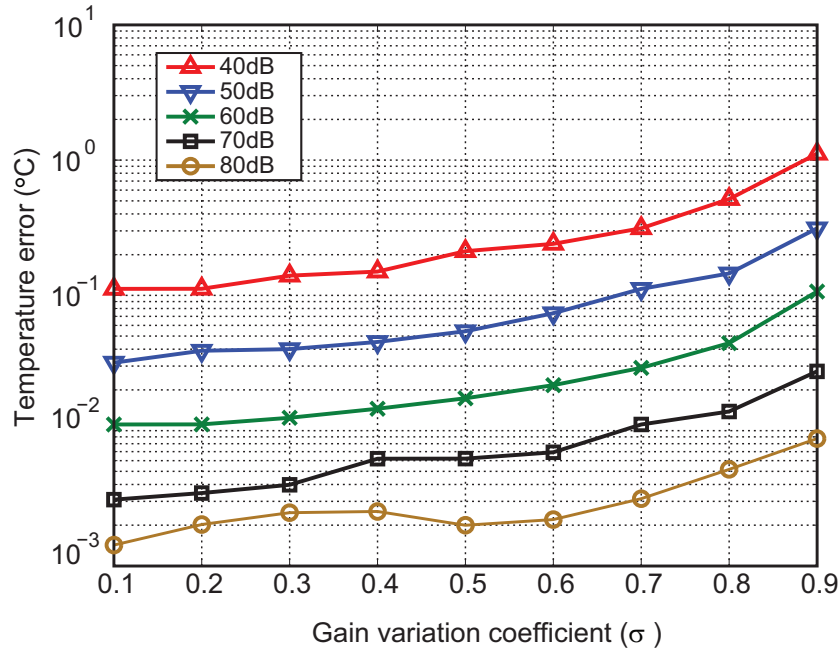


Figure 3.17: Simulated temperature error due to finite DC gain and gain variation in the integrator of a 1<sup>st</sup>-order zoom-ADC with 5-bit coarse conversion.

### 3.3.5.5 DAC Mismatch

As shown in Figure 3.7, the proposed zoom-ADC's  $\Delta\Sigma$  modulator uses a 1-bit quantizer, and so its linearity is not limited by quantizer offset and offset spread. However, the ADC's overall linearity is limited by the nonlinearity of its multi-bit DAC, which is caused by the mismatch of the various DAC elements, e.g. the mismatch between the unit capacitors of a unary-weighted capacitor DAC (CDAC). As in other two-step ADC structures, this mismatch could result in ADC non-linearity. The standard metric to evaluate this is to measure the ADC's integral non-linearity (INL) and differential non-linearity (DNL). While INL is considered as the error with respect to a best linear-fit, DNL reflects the discontinuities in the ADC's transfer function.

Although the resulting non-linearity can be improved by calibration and/or trimming of the DAC references, the associated calibration time significantly increases production costs, thus prohibiting the use of such techniques for large volume production. A suitable solution would be to apply DEM to the unit elements of the multi-bit DAC. During the fine conversion phase, a subset of DAC elements which form the  $\Delta\Sigma$ -ADC's references can be dynamically interchanged, thus averaging out their mismatch. The simulation model of Figure 3.15 can be used to investigate the effect of the mismatch between the unit elements of the feedback DAC. Figure 3.18, top, shows the simulated INL or temperature

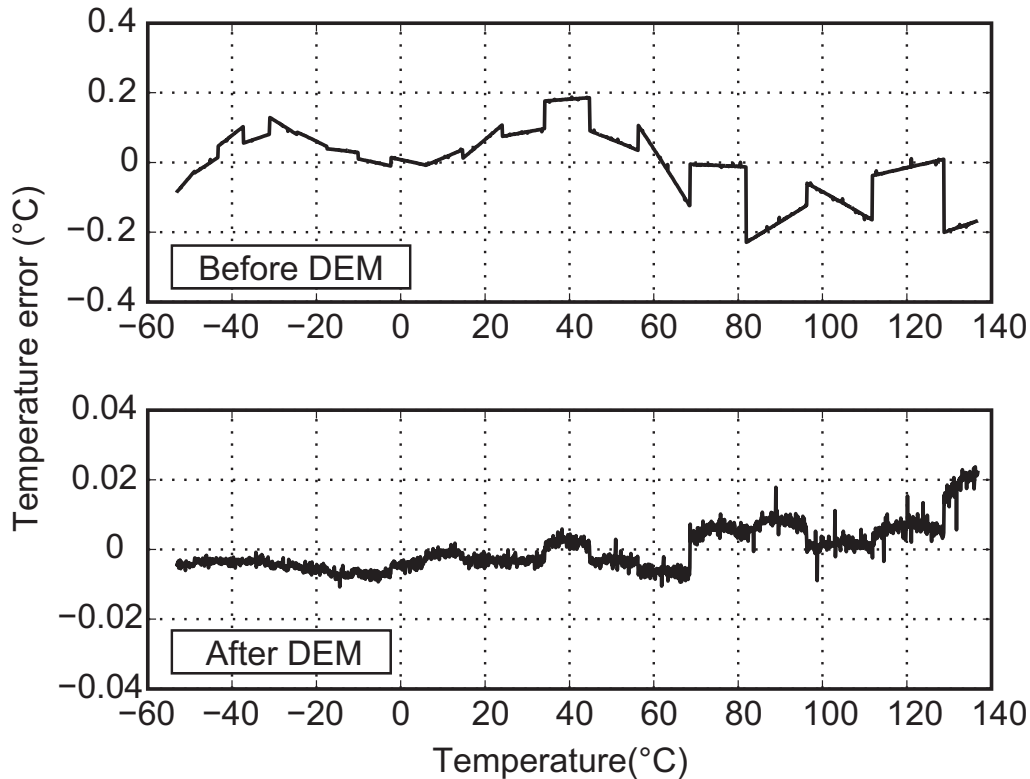


Figure 3.18: Simulated temperature error in the 1<sup>st</sup>-order zoom-ADC, assuming a mismatch with normal distribution (zero mean and  $\sigma = 1\%$ ) between the DAC elements: before (top) and after (bottom) applying DEM ( $N_{\Delta\Sigma} = 1024$ ).

error of a 1<sup>st</sup>-order zoom-ADC with the random mismatch of the feedback DAC, assuming a normal distribution with mean = 0 and  $\sigma = 1\%$ . As shown, the resulting temperature errors can be as high as  $\pm 0.2^\circ\text{C}$ , which is unacceptably large for precision applications. Applying DEM will improve the accuracy to  $\pm 0.02^\circ\text{C}$ , representing a 10x improvement as shown in Figure 3.18, bottom.

Measuring the DNL of a smart temperature sensor is not trivial, as it requires the generation of temperature steps that are smaller than the ADC's quantization error. A more practical alternative is to take the difference between consecutive ADC outputs, thereby defining a pseudo-DNL function. If the rate of change at the input of the ADC is less than its conversion speed, any discontinuities in the ADC's characteristics will be captured by this experiment, which can be readily performed by exposing the sensor to a thermal ramp. The simulated pseudo-DNL for a normal distribution (mean = 0 and  $\sigma = 1\%$ ) between the DAC's unit elements is shown in Figure 3.19, top. As shown, a significant discontinuities between the different fine segments is observed, e.g.  $\pm 0.2^\circ\text{C}$  at  $X = 7.5$  and  $X = 10.5$ . Applying DEM reduces these to well below  $\pm 0.015^\circ\text{C}$ . In

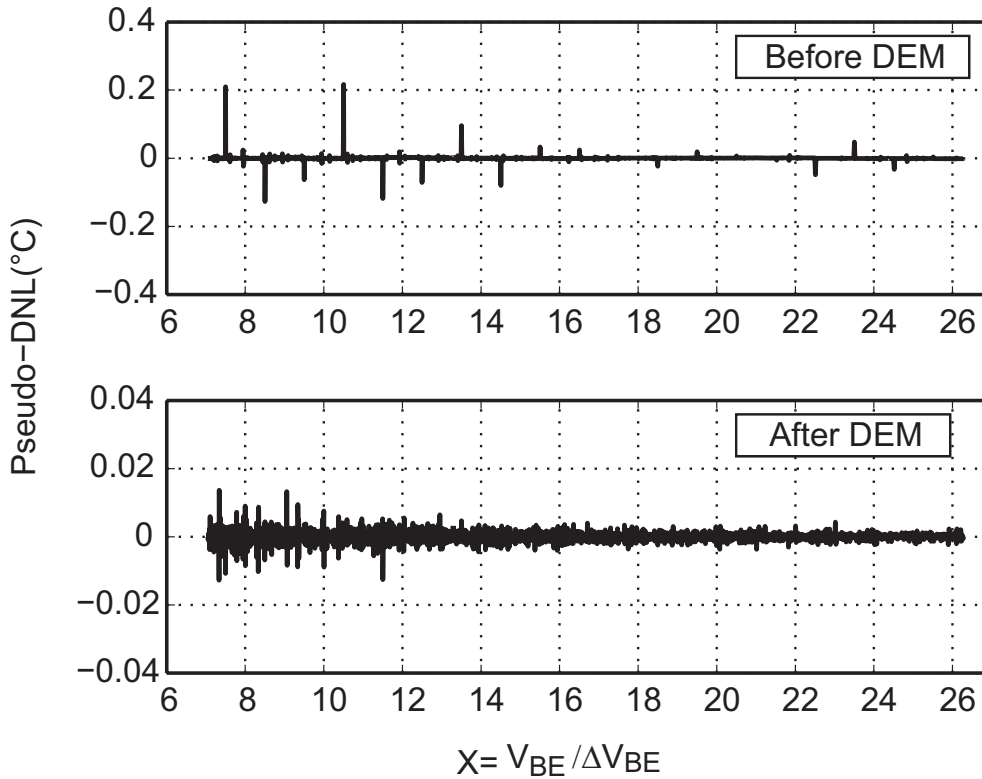


Figure 3.19: Simulated pseudo-DNL of the 1<sup>st</sup>-order zoom-ADC, assuming a mismatch with normal distribution (zero mean and  $\sigma=1\%$ ) between the DAC elements: before (top) and after (bottom) applying DEM ( $N_{\Delta\Sigma}=1024$ ).

practice, careful design and layout can reduce the mismatch of the DAC's unit elements to less than 1%.

### 3.4 Curve Fitting and Trimming

So far, the design of an ADC to digitize the ratio  $X$  has been discussed. However, further digital signal processing is required before an output  $D_{out}$  in degrees Celsius can be obtained. By combining equations (3.3) and (2.9),  $D_{out}$  can be expressed as follows:

$$D_{out} = A \cdot \frac{\alpha}{X + \alpha} + B, \quad (3.14)$$

indicating that the parameters  $\alpha$ ,  $A$ , and  $B$  should be determined before  $D_{out}$  can be obtained. The mapping parameter  $\alpha$  can be found by solving equation (2.7), which requires

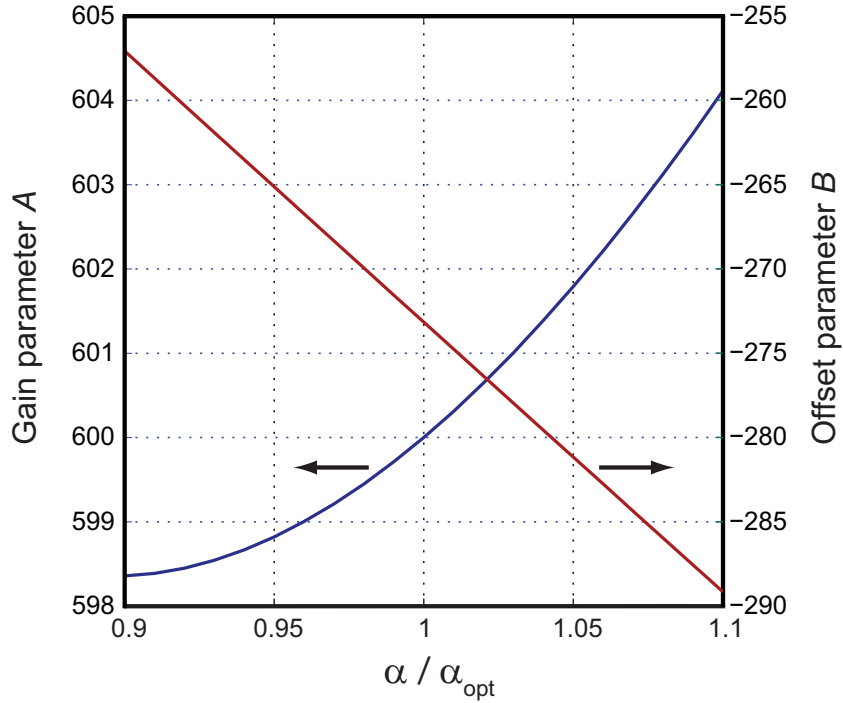


Figure 3.20: The required adjustments of the gain and offset parameters  $A$ ,  $B$ , as a result of  $\pm 10\%$  deviation of  $\alpha$  from  $\alpha_{opt}$ .

the knowledge of the design parameters  $p$  and  $S_{V_{BE}}^T$ . In practice, during a *calibration* phase the ADC's output  $X$  is carefully characterized over the temperature range of interest. During this phase, a *reference* sensor, in thermal equilibrium with the device under test (DUT), provides accurate temperature information  $T_{ref}$ , which is logged along with  $X$  [3.5]. The optimum mapping parameter  $\alpha_{opt}$  can then be found by applying a least-squares linear fit to the resulting  $\mu$  from equation (3.3). For a properly designed sensor front-end and readout circuit, the resulting  $\alpha_{opt}$  should match that obtained from equation (2.7). Assuming that  $p = 5$ , and  $S_{V_{BE}}^T = -2\text{mV}/^\circ\text{C}$ , a value of  $\alpha_{opt} = 15.218$  is obtained by both methods.

By means of linear fitting of  $\mu$  to the temperature information  $T_{ref}$ , obtained during the calibration phase, the scaling parameters  $A$  and  $B$  can also be found. It can be seen from equation (3.14) that their values will also depend on  $\alpha$ . This becomes more clear by applying a Taylor expansion of equation (3.14) around  $\alpha_{opt}$  as follows:

$$D_{out}(\alpha_{opt}) = A \cdot \frac{\alpha_{opt} \cdot \Delta V_{BE}}{V_{BE} + \alpha_{opt} \cdot \Delta V_{BE}} + B + A \cdot \frac{\Delta V_{BE} \cdot V_{BE}}{(V_{BE} + \alpha_{opt} \cdot \Delta V_{BE})^2} \cdot (\alpha - \alpha_{opt}) + \dots, \quad (3.15)$$

indicating that at a given temperature, any deviation of  $\alpha$  from  $\alpha_{opt}$  will require the adjustment of both  $A$  and  $B$  parameters, in order to achieve the same  $D_{out}$ . Figure 3.20 shows the resulting modifications in  $A$  and  $B$ , as a function of  $\alpha = \alpha_{opt} \pm 10\%$ . As can be seen,  $B$  linearly decreases with  $\alpha$ , while  $A$  increases parabolically. Given that for  $\alpha = \alpha_{opt}$ ,  $A = 600$  and  $B = -273.14$ , a 1% increase from  $\alpha_{opt}$  requires  $A$  and  $B$  to change by 0.052% and 0.59%, respectively, indicating a significantly larger sensitivity of  $B$  to variations in  $\alpha$ .

So far,  $V_{BE}$  is assumed to be linear with a sensitivity  $S_{V_{BE}}^T = -2\text{mV}/^\circ\text{C}$ , which is not an accurate assumption. In practice,  $V_{BE}$  exhibits non-linearity or *curvature* in the order of few mV, depending on the bias conditions, e.g. using a PTAT bias current for a substrate PNP results in a  $V_{BE}$  curvature up to  $\approx 3\text{mV}$  from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$  [3.5]. By applying the values of  $\alpha_{opt}$ ,  $A$ , and  $B$ , as obtained so far, a systematic non-linearity error up to  $\approx 1^\circ\text{C}$  at the sensor's output can be measured. Although this can be compensated by digital post processing, this approach increases the complexity of the digital back-end. As shown in [3.12], [3.13], a more straightforward compensation can be achieved by slightly increasing the value of  $\alpha_{opt}$ , in order to minimize the non-linearity of  $\mu$  due to the curvature in  $V_{BE}$ . The value of parameters  $A$  and  $B$  should then be slightly adjusted to obtain the required linear mapping between  $\mu$  and the output  $D_{out}$ . The resulting residual non-linearity is then reduced to less than  $\pm 0.1^\circ\text{C}$  from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$  [3.12], which is sufficiently small in most of applications.

In practice, the sensor's output deviates from the desired value due to the process variation of different on-chip components. During the calibration phase, the resulting error can be accurately measured and then nulled by adjusting one of the different possible *trimming* knobs. In a BJT-based sensor, the use of dynamic correction techniques ensures that the dominant source of inaccuracy is the process spread of  $V_{BE}$ , which is PTAT in nature, and therefore can be trimmed at a single temperature [3.5]. Assuming  $\pm 2\text{mV}$  spread of  $V_{BE}$  (at room temperature), Figure 3.21 shows that the resulting PTAT error over the temperature range from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$  is about  $\pm 0.9^\circ\text{C}$  at  $125^\circ\text{C}$ . The most effective way to trim such an error is to apply a PTAT compensation to  $V_{BE}$ . A straightforward way of doing this is to trim the BJT's bias current [3.5]. Alternatively, the adjustment can be made in the digital domain, by modifying equation (2.9) as follows:

$$D_{out} = A \cdot \frac{\mu}{1 - \gamma_D \mu} - B. \quad (3.16)$$

Here  $\gamma_D$  is a calibration constant that is determined as follows:

$$\gamma_D = \frac{1}{\mu} - \frac{1}{\mu_{ideal}}, \quad (3.17)$$

where  $\mu_{ideal}$  is the desired ratio at the calibration temperature [3.5]. The implementation of digital PTAT trim, therefore, requires extra digital signal processing, including a division and a multiplication, which translates to extra chip area.

However, equation (3.14) suggests that the parameters  $A$ ,  $B$ , and  $\alpha$  can also be used as trimming knobs. For example, an offset trim can be readily implemented by adjusting  $B$ . This choice, however, is not optimal as it does not completely cancel out the error due to the  $V_{BE}$ 's PTAT spread. For the  $\pm 2\text{mV}$   $V_{BE}$  spread, an offset trim at room temperature reduces the inaccuracy to around  $\pm 0.4^\circ\text{C}$  at  $125^\circ\text{C}$ , as shown in Figure 3.21. Depending on the target accuracy, offset trim can be considered as an attractive alternative, due to its simplicity of implementation. In the proposed readout architecture in this work,  $\alpha$  is implemented in the digital domain, and thus can also be used as a trim knob. As shown in Figure 3.21, a single-point  $\alpha$ -trim at room temperature results in a slightly greater inaccuracy compared to that obtained with an offset trim; around  $\pm 0.45^\circ\text{C}$  at  $125^\circ\text{C}$ . Since the PTAT error is a particular sort of gain error, a higher accuracy is expected by trimming  $A$ . As shown, this reduces the inaccuracy to about  $\pm 0.25^\circ\text{C}$  at  $125^\circ\text{C}$ .

The discussion above suggests that PTAT trim is the most suitable trimming method for BJT-based sensors, as it fully compensates for the PTAT spread of  $V_{BE}$ . In practice, however, any residual non-PTAT error will impact the effectiveness of trimming, and therefore a complete cancellation will not be possible. Various error sources can contribute

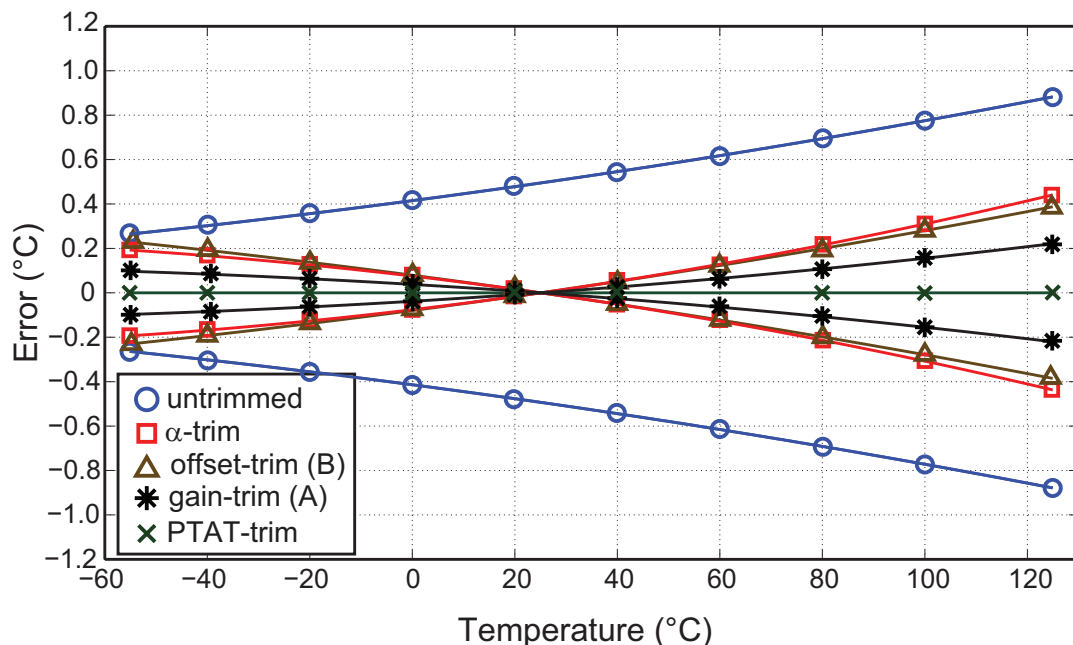


Figure 3.21: Error due to the PTAT spread in  $V_{BE}$  ( $\pm 2\text{mV}$  at  $25^\circ\text{C}$ ), before and after trimming: different single-point trim methods are employed.

to such non-PTAT residuals, e.g. the high-order non-linearity or curvature of  $V_{BE}$ , packaging stress, and other residual errors in the sensor front-end and readout circuit [3.5]. Moreover, any error during the calibration process will also directly impact the trimming, and therefore, the sensor's ultimate accuracy. The adopted trimming method, therefore, depends on the nature of any non-PTAT residual errors in the design, the target accuracy, the hardware constraints on the choice of trimming parameter, and the calibration accuracy.

### 3.5 Conclusions

In this chapter, a new readout architecture was proposed to bridge the energy-efficiency gap of the BJT-based temperature sensors. Given the fact that temperature is a slowly changing quantity, a two-step zoom-ADC architecture was proposed, which combines the speed of a coarse SAR-ADC, with the high resolution/accuracy of a fine  $\Delta\Sigma$ -ADC. In the zoom-ADC, the full-scale range of the fine converter is considerably reduced, thus notably relaxing various key requirements such as the number of  $\Delta\Sigma$ -cycles and the DC-gain and swing of the loop filter. Since both conversion time and power-efficiency can be improved in this architecture, a substantial energy-efficiency improvement can be achieved with the zoom-ADC. Lastly, the fact that dynamic correction techniques can be used in the fine conversion phase ensures that the accuracy of the zoom-ADC can be as good as that of conventional  $\Delta\Sigma$ -ADC architectures.

The required curve fitting to achieve the digital output in degrees Celsius was also discussed. This results in a set of parameters  $A$ ,  $B$ , and  $\alpha$  which determine the appropriate mapping. As shown, in order to minimize the non-linearity of the sensor's output, the value of  $\alpha$  should be optimally set. It was also observed that any deviation from this optimal value will call for small adjustments of the gain and offset parameters  $A$  and  $B$ . Finally, different trimming methods were explored to compensate for the sensor's dominant source of inaccuracy: the PTAT spread of  $V_{BE}$ . As discussed, a PTAT trim is the most suitable trimming method to cancel out such error. The other fitting parameters such as  $A$ ,  $B$ , and  $\alpha$ , can also be used as trimming knobs, but result in significantly greater inaccuracy when compared to PTAT trim.

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## Chapter 4

# BJT-Based, Energy-Efficient Temperature Sensors

As was shown in the previous chapter, the zoom-ADC is well suited for use in energy-efficient temperature sensors. It combines the strengths of SAR- and  $\Delta\Sigma$ -ADCs to realize an accurate, and energy-efficient temperature to digital conversion. In this chapter, a sensor prototype that employs a 1<sup>st</sup>-order zoom-ADC is described. It is compact and power efficient, requiring only a few  $\mu\text{W}$  to operate. Its energy-efficiency, however, is limited, due to the use of an inherently slow 1<sup>st</sup>-order  $\Delta\Sigma$  modulator. To improve energy-efficiency, a second prototype is presented, which requires less power to operate, while employing a 2<sup>nd</sup>-order zoom-ADC and a faster sampling scheme. Finally, a third prototype for sensing very high temperatures ( $>150^\circ\text{C}$ ) is presented, which uses robust techniques to overcome the different sources of temperature sensing errors at such temperatures.

### 4.1 A Micropower Temperature Sensor<sup>1</sup>

A zoom-ADC can be used to digitize the ratio  $X = V_{BE}/\Delta V_{BE}$ , which is a measure of temperature. As shown in Figure 4.1, for  $p = 5$ , the ratio  $X$  is a non-linear function of temperature, which ranges between 7 and 24 from  $-40^\circ\text{C}$  to  $125^\circ\text{C}$ . Once  $X$  is known, the PTAT function  $\mu$  (see Figure 4.1) can be easily determined as follows:

$$\mu = \frac{\alpha \cdot \Delta V_{BE}}{V_{BE} + \alpha \cdot \Delta V_{BE}} = \frac{\alpha}{X + \alpha}. \quad (4.1)$$

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<sup>1</sup>K. Souri and K. A. A. Makinwa, "A 0.12mm<sup>2</sup> 7.4 $\mu\text{W}$  Micropower Temperature Sensor with an Inaccuracy of  $\pm 0.2^\circ\text{C}$  ( $3\sigma$ ) from  $-30^\circ\text{C}$  to  $125^\circ\text{C}$ ," *IEEE Journal of Solid-State Circuits*, vol. 46, is. 7, pp. 1693 - 1700 July. 2011.

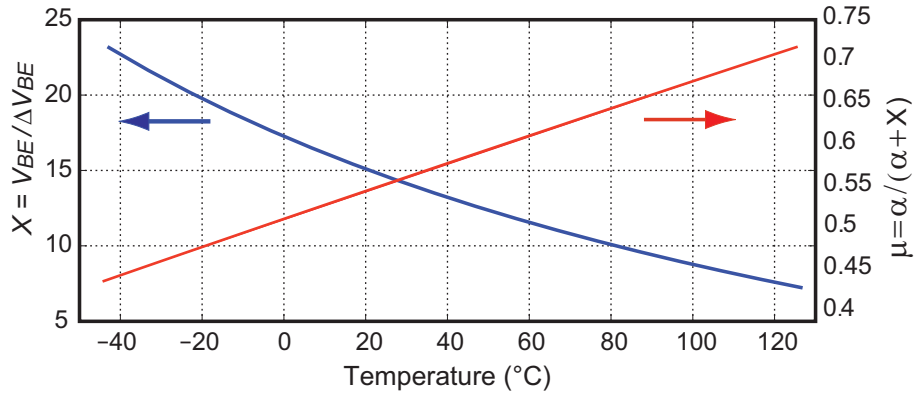


Figure 4.1: Non-linear  $X = V_{BE}/\Delta V_{BE}$  ( $p = 5$ ) and linearized  $\mu = \alpha/(\alpha+X)$  as a function of temperature.

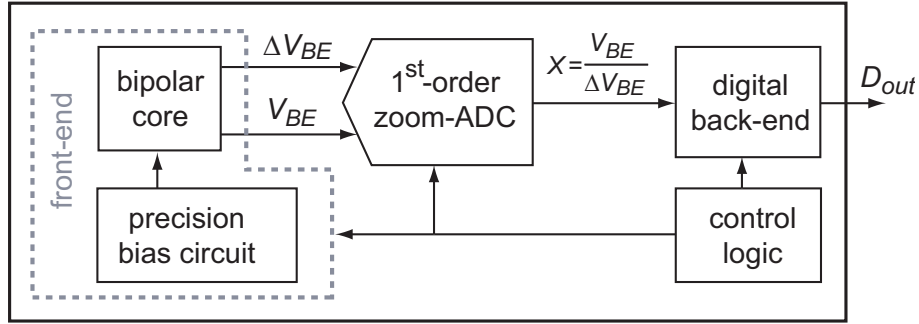


Figure 4.2: Block diagram of the smart temperature sensor.

The block diagram of the proposed temperature sensor is shown in Figure 4.2. It consists of a 1<sup>st</sup>-order zoom-ADC and an analog front-end: a precision bias circuit whose PTAT output current biases the substrate PNP transistors of a bipolar core. The  $V_{BE}$  and  $\Delta V_{BE}$  voltages extracted from the bipolar core are digitized by the zoom-ADC, which outputs  $X$  to a digital back-end, which, in turn, determines the PTAT function  $\mu$  and  $D_{out}$ , the temperature in degrees Celsius, as given by equation (2.9).

## 4.1.1 Analog Front-End

### 4.1.1.1 Topology

Figure 4.3 shows the circuit diagram of the analog front-end. Although the temperature dependence of the bias current  $I_b$  doesn't impact the accuracy of  $\Delta V_{BE}$  (see equation (2.4)), it does impact the systematic non-linearity or curvature of  $V_{BE}$ , and hence the sensor's systematic error. The curvature in  $V_{BE}$  can be reduced by using a PTAT bias

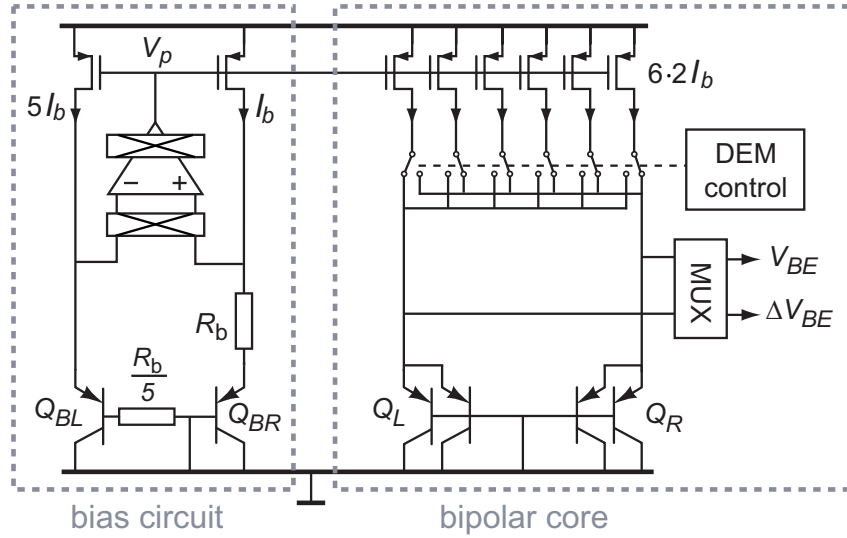


Figure 4.3: Circuit diagram of the analog front-end.

current [4.1]. As shown in Figure 4.3, the bias circuit uses two PNP transistors, biased at a 5 : 1 current ratio. The opamp in the feedback loop then forces  $\Delta V_{BE}$  across the resistor  $R_b$ , resulting in an accurate PTAT current  $I_b$ , which is then used to bias the PNPs of the bipolar core. Each of the PNPs in the front-end has an emitter area of  $A_E = 5\mu\text{m} \times 5\mu\text{m}$ .

#### 4.1.1.2 Effect of Forward Current Gain $\beta_F$

Since a substrate PNP transistor must be biased via its emitter, its collector current and, thus, the resulting  $V_{BE}$  will depend on the transistor's current gain  $\beta_F$ . The spread (up to 50%) and temperature dependence of  $\beta_F$  will then impact the accuracy of  $V_{BE}$ . This effect becomes more significant as the current gain  $\beta_F$  decreases, as is the case in modern CMOS processes (in the  $0.16\mu\text{m}$  CMOS process used  $\beta_F \approx 4.5$ ). The technique known as  $\beta_F$ -compensation mitigates this problem by modifying the PTAT bias circuit to generate a  $\beta_F$ -dependent current [4.2]. This is done by adding a resistor of  $R_b/5$  in series with the base of  $Q_{BL}$ . The opamp in the feedback loop then ensures that:

$$I_b = \ln(5) \cdot \left(\frac{kT}{q}\right) \cdot \left(\frac{1}{R_b}\right) \cdot \left(\frac{1 + \beta_f}{\beta_f}\right). \quad (4.2)$$

When biased with this emitter current, the collector current of the PNPs in the bipolar core ( $Q_L$ ,  $Q_R$ ) will be equal to  $\ln(5) \cdot (kT/q)/R_b$ , and hence their base-emitter voltages will be insensitive to variations in  $\beta_F$ . However, the accuracy of this technique is limited by current-mirror and  $\beta_F$  mismatch. Therefore, careful layout of the PNPs and the current

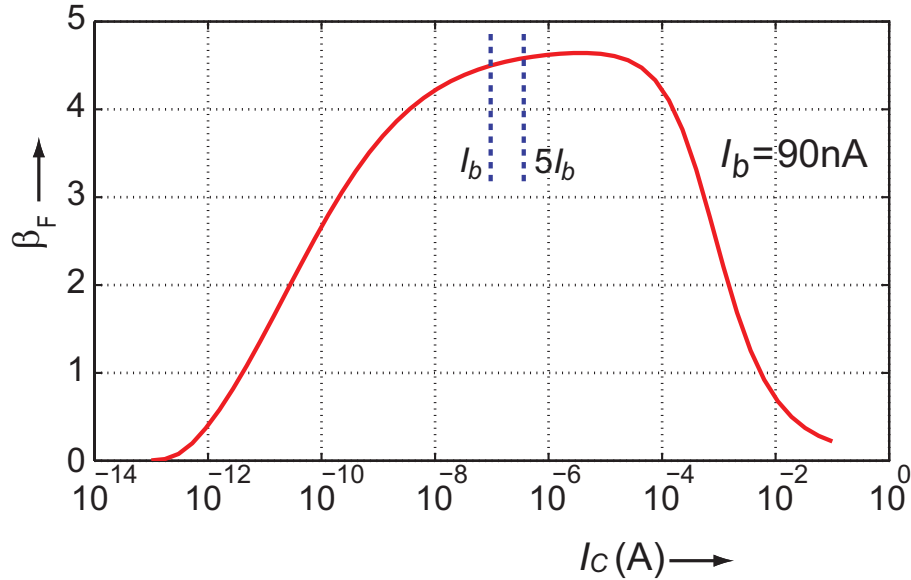


Figure 4.4: The forward current gain  $\beta_F$  (at 25°C) of a substrate PNP ( $A_E = 5\mu\text{m} \times 5\mu\text{m}$ ) as a function of collector current  $I_C$  in the 0.16 $\mu\text{m}$  CMOS process used.

sources is essential. A further source of  $V_{BE}$  error is the current dependency of  $\beta_F$ . In this design, the value of  $I_b$  ( $= 90\text{nA}$  at 25°C) was optimized to ensure that both  $I_b$  and  $5I_b$  are in a relatively flat part of the PNP's  $\beta_F$  versus collector current characteristic (see Figure 4.4) [4.1].

#### 4.1.1.3 Offset Cancellation

Besides the spread in  $R_b$ , the offset  $V_{OS}$  of the opamp in the bias circuit is a major source of bias current inaccuracy. For the temperature sensor to achieve an inaccuracy of less than  $\pm 0.2^\circ\text{C}$ , this offset needs to be less than  $100\mu\text{V}$  [4.2]. However, in CMOS, this cannot be practically achieved by transistor sizing and careful layout. Therefore, the opamp is chopped, so that the resulting bias current will be switched between  $I_b + I_{off}$  and  $I_b - I_{off}$  where  $I_{off} = V_{OS}/R_b$ . Thus, the average of the resulting  $V_{BE}$  in  $Q_L$ ,  $Q_R$  will be, to first-order, independent of  $V_{OS}$ . However, the use of chopping means that there will be square-wave ripple at the opamp's output  $V_P$ . While its amplitude is not important, the complete settling of  $V_P$  at the end of each chopping phase is critical, since this is the moment when the zoom-ADC samples the resulting  $V_{BE}$ . Due to the large input capacitance of the current-mirror MOSFETs, a typical single-stage high-output-impedance opamp will require a relatively large bias current. In [4.3] for example, the opamp was chopped at the ADC's sampling rate and drew a large portion of the sensor's supply current ( $1.7\mu\text{A}$  out of a total of  $6\mu\text{A}$ , or 28%).

## 4.1.1.4 Opamp Topology

As shown in Figure 4.5, an adaptive self-biasing opamp [4.4] - [4.6] is used in this work. It consists of a PMOS input pair ( $M_{5,6}$ ) with diode-connected NMOS loads ( $M_{1,2}$ ). Since the opamp's input voltage is chopped, switch  $S_1$  is used to maintain the correct feedback polarity. The voltage output of the input stage ( $V_{gs,M1}$  or  $V_{gs,M2}$ ) is converted into a current via  $M_3$  and fed back to the differential pair through the  $M_{10}:M_{11}$  current mirror. As a result, the tail current of the input stage is derived from its output voltage. The aspect ratio of the transistors has been chosen such that:  $(W/L)_3 = 4 \cdot (W/L)_{1,2}$  and  $(W/L)_{10} = 2 \cdot (W/L)_{11}$ . This ensures that the current gain of the loop formed by  $M_{1,2} : M_3, M_{10} : M_{11}$  is equal to 1 for zero input. When operated in an open-loop configuration, the positive feedback in this current loop would result in ever-increasing/decreasing output currents for negative/positive differential input voltages, corresponding to a very high DC gain. In the bias circuit, however, the amplifier is operated in a negative feedback loop which stabilizes the circuit and enforces a PTAT current  $I_b$ . Furthermore, since  $M_{10}$  is diode-connected,  $V_P$  is a low impedance node, which reduces the time constant associated with the settling of  $V_P$ . The opamp's bias current, therefore, can be reduced to meet the relaxed load requirements, while maintaining the gain required.

As any current mirror mismatch will result in input-referred offset, high over-drive voltages (260mV and 130mV for the PMOS and NMOS devices respectively, at 25°C) and careful layout are essential. To minimize the effect of channel length modulation on the current gain of the positive feedback loop, a replica circuit drives  $M_8$  and ensures that the  $V_{ds}$  of  $M_3$  is equal to that of  $M_{1,2}$ . At 25°C, the opamp draws only 630nA; significantly

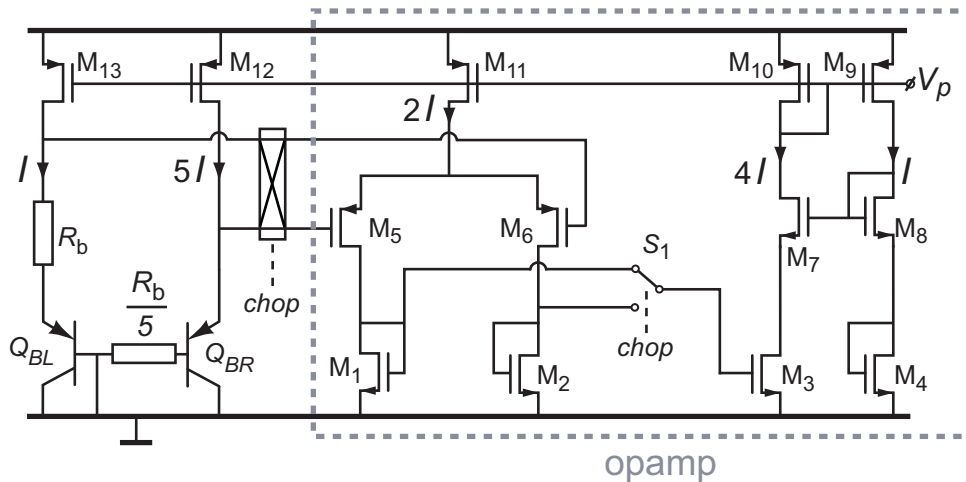


Figure 4.5: Simplified circuit diagram of the bias current circuit (left hand side) and the positive feedback opamp.

less (63%) than in our previous work [4.3]. The entire front-end draws only  $2.1\mu\text{A}$  from a 1.8V supply.

#### 4.1.1.5 Precision Issues

Mismatch in current sources and bipolar devices impacts the accuracy of  $\Delta V_{BE}$ . Even with careful layout, a relative current ratio mismatch  $\Delta p/p$  in the order of 0.1% can be expected, at best. For  $p = 5$ , this will lead to an error of about  $0.14^\circ\text{C}$  at  $T = 25^\circ\text{C}$  [4.1]. Therefore, dynamic element matching (DEM) of the six current sources and two bipolar transistors in the bipolar core is essential to generate the accurate 1 : 5 current ratio required for an accurate  $\Delta V_{BE}$ . As shown in Figure 4.3, DEM switches are placed in series with each current source, in order to swap their current direction according to the DEM control logic. Since  $V_{BE}$  and  $\Delta V_{BE}$  are sampled at the emitter terminal of the BJTs, the voltage drop across DEM switches doesn't impact the accuracy of the sampled voltages.

### 4.1.2 Zoom ADC

#### 4.1.2.1 Topology

In this work, a 1<sup>st</sup>-order zoom-ADC is used to digitize the ratio  $X = V_{BE}/\Delta V_{BE}$ . As previously discussed, it combines a coarse SAR-ADC and a fine 1<sup>st</sup>-order  $\Delta\Sigma$  converter in a two-step conversion scheme [4.3], [4.4]. In this topology, the digital ratio  $X$  is accurately resolved, with a high resolution and within a fairly short conversion time.

As shown in Figure 4.1, the ratio  $X = V_{BE}/\Delta V_{BE}$  ranges from 7 to 24 from  $-40^\circ\text{C}$  to  $125^\circ\text{C}$  ( $p = 5$ ).  $X$  can thus be expressed as  $X = n + \mu'$ , where  $n$  and  $\mu'$  are its integer and fractional parts, respectively, and can be determined separately, as shown in Figure 4.6. In a coarse conversion,  $n$  is determined by a SAR algorithm, which compares  $V_{BE}$  to integer multiples of  $\Delta V_{BE}$ . In a succeeding fine conversion, the fraction  $\mu'$  is then determined by a 1<sup>st</sup>-order  $\Delta\Sigma$ -ADC, whose references are chosen so as to zoom into the region determined by the SAR algorithm, i.e. from  $n \cdot \Delta V_{BE}$  to  $(n+1) \cdot \Delta V_{BE}$ . The region of interest is now quite small (less than  $18^\circ\text{C}$ ), thus relaxing the resolution requirement on the  $\Delta\Sigma$ -ADC, which in turn, leads to energy-efficient conversions [4.3], [4.4].

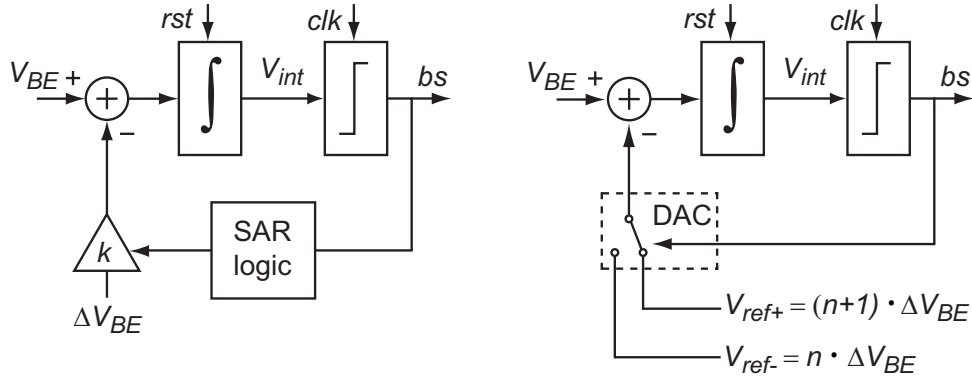


Figure 4.6: Block diagram of the 1<sup>st</sup>-order zoom-ADC during the coarse (left) and fine (right) conversions.

### 4.1.2.2 Implementation

As shown in Figure 4.7, the zoom-ADC is basically a modified 1<sup>st</sup>-order switched-capacitor (SC)  $\Delta\Sigma$ -ADC with 24 unit sampling capacitors. At the start of each comparison step of the coarse conversion, the integrator is reset, and it therefore functions as a sample-and-hold circuit. As shown in Figure 4.8.b,  $V_{BE}$  is then sampled on a single unit capacitor and integrated during one full clock cycle. In the next clock cycle,  $-\Delta V_{BE}$  is sampled on  $k$  unit capacitors and also integrated, thus a total charge proportional to  $(V_{BE} - k \cdot \Delta V_{BE})$  is integrated. The comparator's output  $bs$  then indicates the result of the comparison ( $V_{BE} > k \cdot \Delta V_{BE}$ ). The control logic implements the SAR algorithm, with which  $n$  can

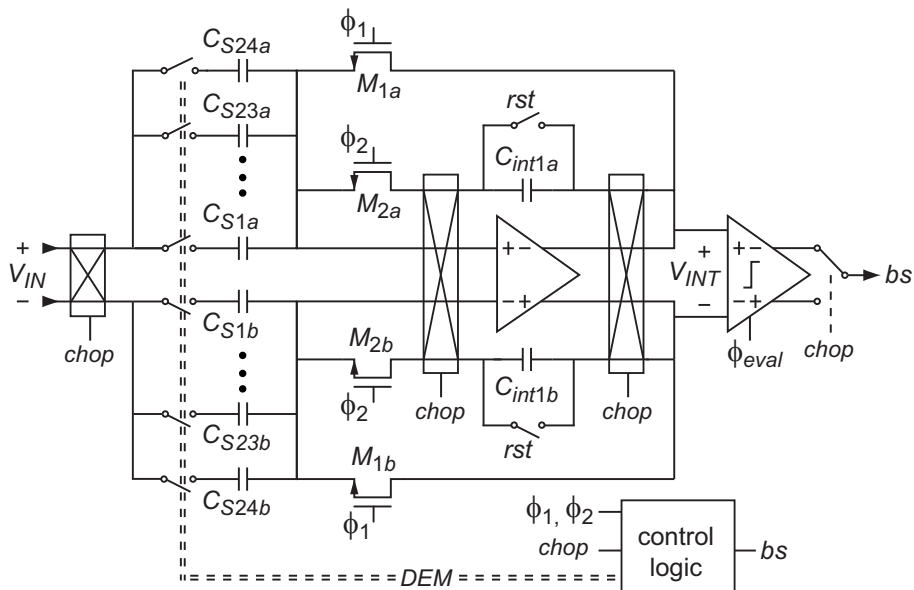


Figure 4.7: Simplified circuit diagram of the 1<sup>st</sup>-order zoom-ADC.



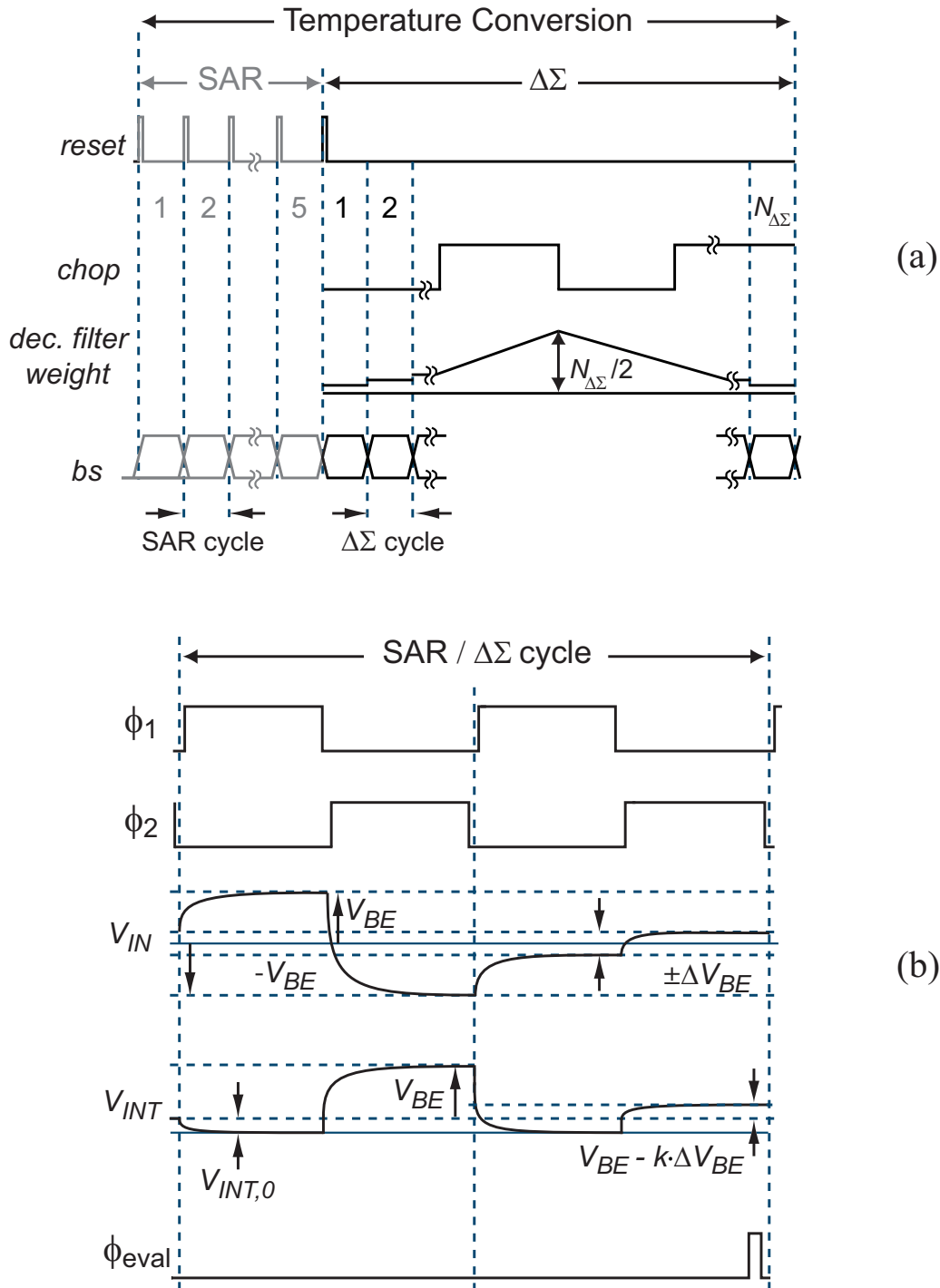


Figure 4.8: (a) Timing diagram of a temperature conversion: (b) Waveforms of a full SAR  $\Delta\Sigma$  cycle.  $V_{IN}$ : Zoom ADC's input voltage,  $V_{INT}$ : Integrator's output voltage.  $V_{INT,0}$ : integrator's initial voltage. The gain factor  $k$  is set by the SAR-logic in the coarse conversion, while  $k = n$  or  $n+1$  when  $bs = 0$  or  $1$  in the fine conversion step.

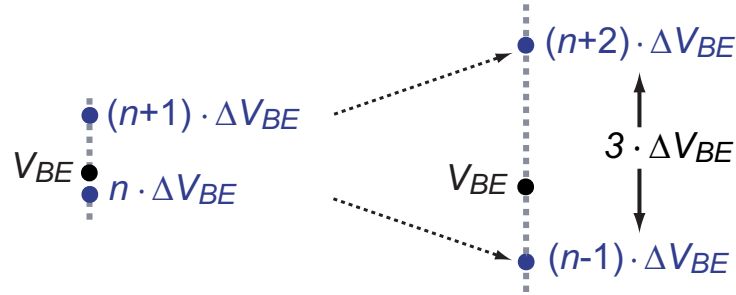


Figure 4.9: The range is extended to  $3 \cdot \Delta V_{BE}$  to avoid out-of-range errors during the fine conversion step.

be determined within five comparison steps, since  $n \leq 24$ . Once  $n$  is known, the fine conversion step is determined with a  $\Delta\Sigma$  charge-balancing scheme (Figure 4.6). After an initial reset, the modulator operates as follows: when  $bs = 0$ ,  $(V_{BE} - n \cdot \Delta V_{BE})$  is integrated, and when  $bs = 1$ ,  $(V_{BE} - (n+1) \cdot \Delta V_{BE})$  is integrated. As shown in Figure 4.8, such integrations require two clock cycles: one to integrate  $V_{BE}$  and one to integrate  $-k \cdot \Delta V_{BE}$ . Since the net integrated charge is approximately zero, the bitstream average is the desired  $\mu = (V_{BE} - n \cdot \Delta V_{BE}) / \Delta V_{BE}$ .

When  $V_{BE} \sim k \cdot \Delta V_{BE}$ , non-idealities such as comparator offset, noise and mismatch during the coarse conversion could lead to incorrect  $n$  values, and therefore clipping in the fine conversion. As discussed in Chapter 2, this can be avoided by performing an extra guard-band cycle. In this work, the fine conversion is appropriately extended to  $3 \cdot \Delta V_{BE}$ , thus relaxing the requirements on the coarse conversion [4.4]. As shown in Figure 4.9, the range is set-up to cover from  $(n-1) \cdot \Delta V_{BE}$  to  $(n+2) \cdot \Delta V_{BE}$ , in such a way that  $V_{BE}$  is always roughly in the middle of this range, and hence avoiding any out-of-ranging in the fine conversion step. In order to accommodate the required accuracy of the gain factor  $k$ , DEM is applied to the elements of the sampling cap-DAC during the  $\Delta\Sigma$  conversion step.

As illustrated in Figure 4.7, the main element of the zoom-ADC is a SC integrator built around a fully-differential folded-cascode opamp with a gain of 86dB, which is required to maintain the linearity of overall ADC across the fine conversion segments. However, due to the relaxed requirements on the ADC's resolution, no gain boosting is required, unlike [4.2], [4.7], thus reducing the area and power of the opamp. A SC common-mode feedback scheme is applied to maintain the opamp's output common mode at a well defined value ( $\approx V_{DD}/2$ ). The sampling capacitors are also quite small:  $C_S = 120\text{fF}$ , while the integration capacitors are  $2 \cdot C_S$ , in order to maintain the integrator swing within the folded-cascode topology limits. The size of sampling caps are defined based on the target resolution at sensor's output, which is determined by  $kT/C$  noise of the

sampling capacitors, and the conversion time.

The opamp's offset and  $1/f$  noise are reduced by correlated double-sampling (CDS) during both the coarse and fine conversions. Another source of inaccuracy in a SC integrator is the charge injection of the switches. A fully-differential structure combined with the bottom-plate sampling technique is known to be the most effective way to mitigate this, as the injected charge then appears as a common-mode effect [4.8]. The effectiveness of this technique, nevertheless, is mainly limited by the mismatch between switches, and thus their charge injection. To overcome this, the entire ADC is chopped twice per fine conversion. This is done by swapping the polarity of input voltage and the quantizer output, i.e. bitstream  $bs$ . State preserving switches around the opamp (Figure 4.7) are essential to maintain the integrator state when chopping is applied [4.2]. The timing diagram of the chopping signal is shown in Figure 4.8.a.

### 4.1.3 Measurement Results

The temperature sensor was realized in a standard  $0.16\mu\text{m}$  CMOS process with five metal layers (Figure 4.10). The chip has an active area of  $0.12\text{mm}^2$ , and consumes  $8.2\mu\text{W}$  from a  $1.8\text{V}$  supply at  $25^\circ\text{C}$ . The digital back-end, the control logic and the fine conversion's  $\text{sinc}^2$  decimation filter were implemented off-chip for flexibility. For characterization, the prototypes were packaged in ceramic DIL packages and placed in a climate chamber, in good thermal contact with an aluminum block containing a platinum Pt-100 resistor calibrated to  $20\text{mK}$ . With this setup, 19 samples from one batch were then characterized over

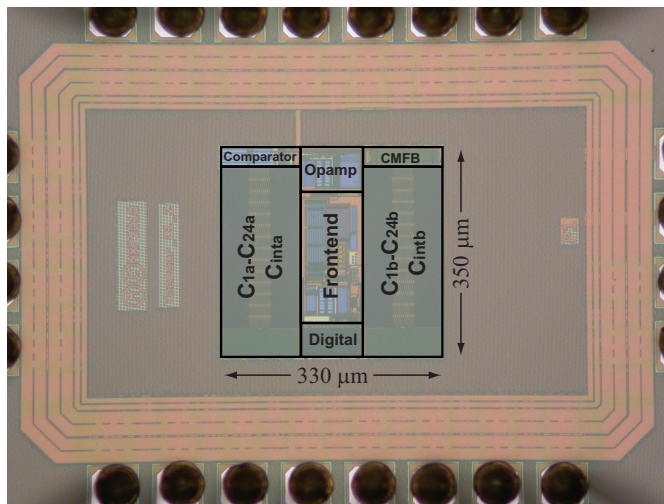


Figure 4.10: Chip micrograph of the first prototype sensor employing a 1<sup>st</sup>-order zoom-ADC.

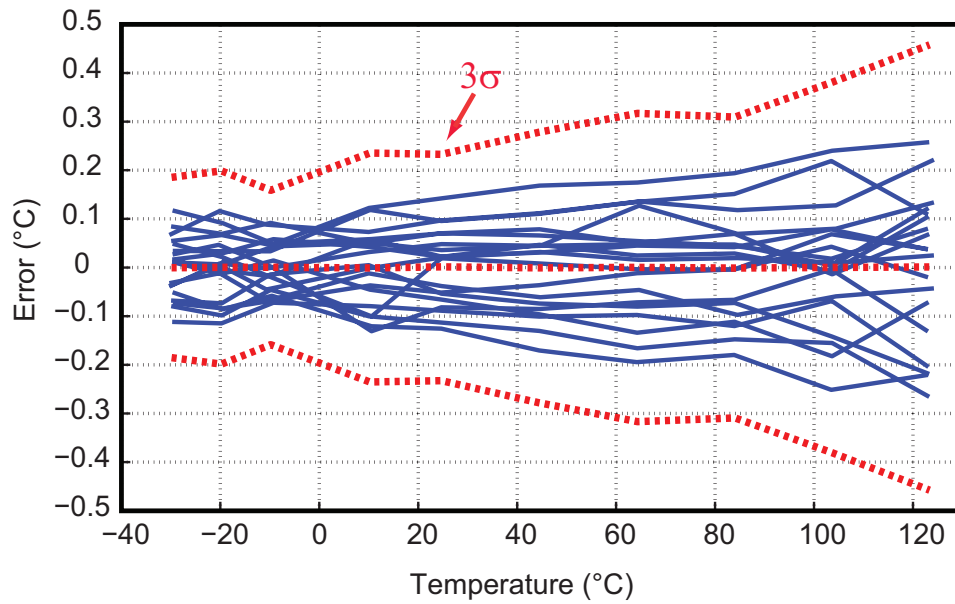


Figure 4.11: Measured temperature error of 19 sensors before trimming; dashed lines refer to the average and  $\pm 3\sigma$  limits.

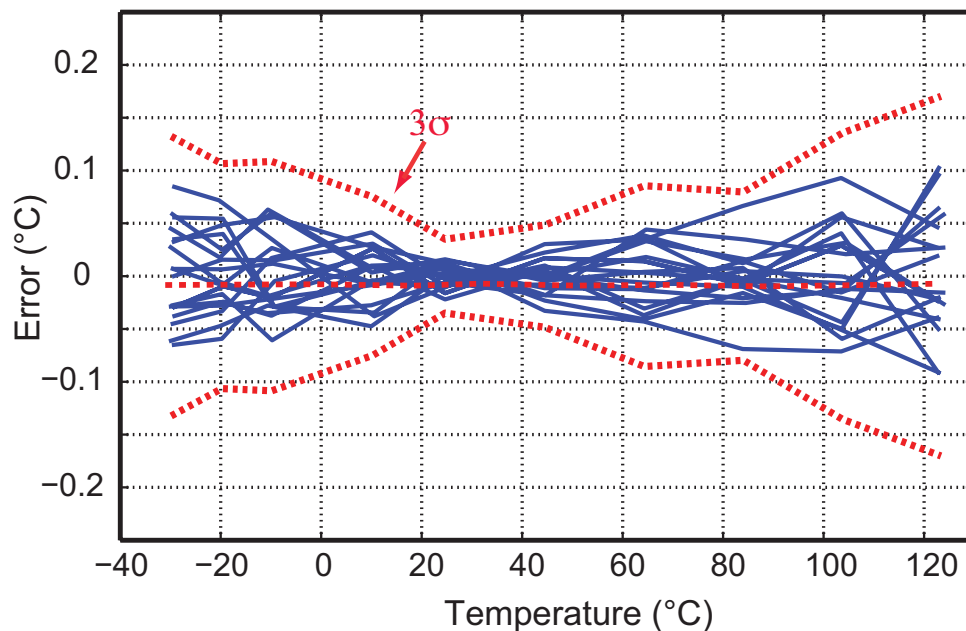


Figure 4.12: Measured temperature error of 19 sensors after a single  $\alpha$ -trim at  $25^\circ\text{C}$ ; dashed lines refer to the average and  $\pm 3\sigma$  limits.

the temperature range from  $-30^\circ\text{C}$  to  $125^\circ\text{C}$ . As shown in Figure 4.11, the resulting batch-calibrated inaccuracy was  $\pm 0.5^\circ\text{C}$  ( $3\sigma$ ), after digital compensation for residual curvature ( $\pm 0.25^\circ\text{C}$ ). The standard deviation  $\sigma$  is an estimated value, which is obtained based on a

limited number of samples: 19 samples in this case. A single digital trim at 25°C was used to compensate for  $V_{BE}$ 's PTAT spread [4.3], [4.4]. This was done by individually tuning and embedding the  $\alpha$  value for each sensor in the digital back-end, thereby reducing the inaccuracy to  $\pm 0.2^\circ\text{C}$  ( $3\sigma$ ) as shown in Figure 4.12. At 10 conversions/s (1024  $\Delta\Sigma$  cycles), the sensor achieves a resolution of 15mK (rms). The sensor operates from a 1.6V to 2V supply with a sensitivity of  $0.1^\circ\text{C}/\text{V}$ . The sensor's performance is summarized in Table 4.1, and compared to other accurate temperature sensors [4.2], [4.7].

To assess the sensor's energy-efficiency, its resolution figure-of-merit (FoM) can be calculated. The sensor operates from a minimum supply voltage of 1.6V and draws  $4.6\mu\text{A}$ . Given the resolution of 15mK in a conversion time of 100ms, this translates into a resolution FoM of  $0.17\text{nJ}^\circ\text{C}^2$ . Compared to the other designs in Table 4.1, the energy-efficiency has been improved by more than 11x, while achieving comparable accuracy.

Parameter	1 <sup>st</sup> Prototype JSSC'11 [4.10]	JSSC'05 [4.2]	ISSCC'09 [4.7]
CMOS technology	0.16 $\mu\text{m}$	0.7 $\mu\text{m}$	0.7 $\mu\text{m}$
Chip area	0.12 $\text{mm}^2$	4.5 $\text{mm}^2$	4.5 $\text{mm}^2$
Supply current (RT) <sup>†</sup>	4.6 $\mu\text{A}$	75 $\mu\text{A}$	25 $\mu\text{A}$
Supply voltage	1.6V - 2V	2.5V - 5.5V	2.5V - 5.5V
Inaccuracy (trim points)	$\pm 0.2^\circ\text{C}$ ( $3\sigma$ ) (1)	$\pm 0.1^\circ\text{C}$ ( $3\sigma$ ) (1)	$\pm 0.1^\circ\text{C}$ ( $3\sigma$ ) (1)
Temperature range	-30°C - 125°C	-55°C - 125°C	-55°C - 130°C
Resolution ( $T_{\text{conv}}$ )	0.015°C (100msec)	0.01°C (100msec)	0.025°C (100msec)
Resolution FoM	0.17 $\text{nJ}^\circ\text{C}^2$	1.9 $\text{nJ}^\circ\text{C}^2$	3.9 $\text{nJ}^\circ\text{C}^2$

<sup>†</sup> Excluding the off-chip digital

Table 4.1: Performance summary of the first prototype employing a 1<sup>st</sup>-order zoom-ADC and compared to previous work.

## 4.2 An Energy-Efficient Temperature Sensor<sup>2</sup>

As demonstrated in the previous section, using a 1<sup>st</sup>-order zoom-ADC results in a low-power sensor. However, its energy-efficiency is limited by the low conversion rate of its 1<sup>st</sup>-order modulator. Furthermore, to maintain linearity across the zoom segments, an opamp with a DC gain in excess of 80dB was necessary. Combined with the large swing requirement ( $>\pm 350\text{mV}$ ) in the  $V_{BE}$  integration phase (see Figure 4.8), this led to a topology with limited power efficiency in the target  $0.16\mu\text{m}$  CMOS process. In this section a temperature sensor based on a 2<sup>nd</sup>-order zoom-ADC will be presented that is significantly more energy-efficient, and achieves the required loop-gain while dissipating less power.

### 4.2.1 Improving Energy-Efficiency

As in the first prototype, a 5-bit SAR-ADC performs the coarse conversion step. The fine step, however, uses the block diagram of Figure 4.13. It is based on a single-bit feed-forward 2<sup>nd</sup>-order  $\Delta\Sigma$ -ADC. This architecture requires about 8x less  $\Delta\Sigma$  cycles to achieve almost the same resolution as in the first prototype. Moreover, the overall loop-gain is now achieved by using two low-gain, low-power integrators, thus significantly improving power- and energy-efficiency.

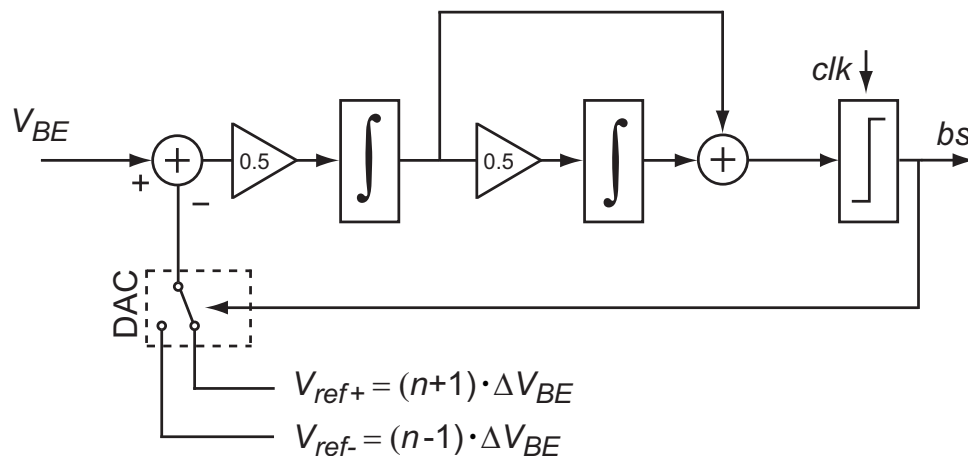


Figure 4.13: Block diagram of the proposed 2<sup>nd</sup>-order zoom-ADC during the fine conversion.

<sup>2</sup>K. Souri, Y. Chae and K. A. A. Makinwa, "A CMOS Temperature Sensor with a Voltage-Calibrated Inaccuracy of  $\pm 0.15^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ ," *IEEE Journal of Solid-State Circuits*, vol. 48, is. 1, pp. 292 - 301 Jan. 2013.

The modulator's stability is achieved by a single feed-forward path around the second integrator. This leverages a distinct advantage of the zooming algorithm: the fact that since  $V_{BE} \sim n \cdot \Delta V_{BE}$  during the fine conversion step, the error signal processed by the loop filter is quite small, thus reducing the output swing of the two opamps. As a result, no extra direct feed-forward path between the input terminal and the quantizer's input is required, as is the case in the well-known low-swing feed-forward architecture [4.9]. This simplifies the modulator's implementation, reduces the loading on the analog front-end that generates  $V_{BE}$ , and eliminates a potential source of parasitic coupling into the summing node at the quantizer's input.

As previously discussed, the range in the fine conversion step of the first prototype was extended by 3x to accommodate small errors during the coarse conversion phase. In this work, the range is extended by only 2x, which further improves the conversion speed, since the modulator's full range has been reduced. The necessary information is obtained during a guard-band step, in which  $V_{BE}$  is compared to  $(n+0.5) \cdot \Delta V_{BE}$ . Depending on the result, the references of the  $\Delta\Sigma$ -ADC are then set to either  $(n-1) \cdot \Delta V_{BE}$  and  $(n+1) \cdot \Delta V_{BE}$ , or  $n \cdot \Delta V_{BE}$  and  $(n+2) \cdot \Delta V_{BE}$ . In the rest of this section, for simplicity, we shall assume that the former is the case, as shown in Figure 4.13.

### 4.2.2 An Energy-Efficient Integration Scheme

During the fine conversion, as shown in Figure 4.13, every  $\Delta\Sigma$  cycle requires the integration of either  $(V_{BE} - (n-1) \cdot \Delta V_{BE})$  or  $(V_{BE} - (n+1) \cdot \Delta V_{BE})$ , when the comparator's output  $bs$  is either 0 or 1, respectively. For simplicity, let's assume that a charge proportional to  $(V_{BE} - k \cdot \Delta V_{BE})$  is integrated during one  $\Delta\Sigma$  cycle, where  $k$  is either  $(n-1)$  or  $(n+1)$  depending on the polarity of  $bs$ . In the first prototype, this was performed using a SC integrator and in two clock cycles: in a first clock cycle a charge proportional to  $V_{BE}$  was integrated, while in a second clock cycle a charge proportional to  $-k \cdot \Delta V_{BE}$  was integrated. A folded-cascode opamp was therefore required to accommodate the large swing during the  $V_{BE}$  integration phase, and thus the low-swing advantage of zooming was not fully exploited.

In this design, the two clock cycles are combined i.e. both  $V_{BE}$  and  $\Delta V_{BE}$  are simultaneously sampled and then integrated in only *one* clock cycle. As shown in Figure 4.14 during the sampling phase  $\phi_1$ ,  $V_{BE}$  is sampled on  $C_S$  while  $-\Delta V_{BE}$  is simultaneously sampled on  $k \cdot C_S$ , thus a charge proportional to  $(V_{BE} - k \cdot \Delta V_{BE})$  is stored on the sampling capacitors. The polarity of both input voltages is swapped during  $\phi_2$ , and therefore a charge proportional to  $2 \cdot (V_{BE} - k \cdot \Delta V_{BE})$  is integrated during each clock cycle. Due to the charge cancellation between  $V_{BE}$  and  $-k \cdot \Delta V_{BE}$ , the integrated charge difference

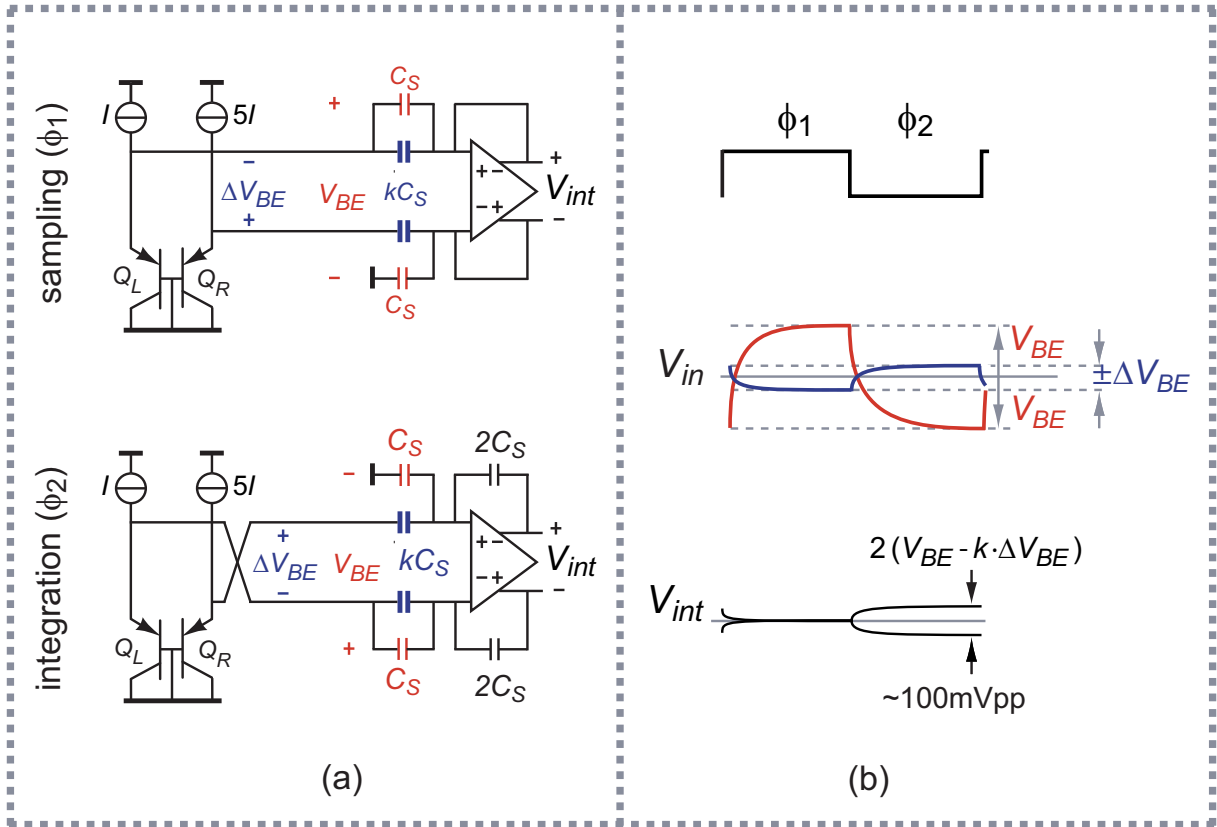


Figure 4.14: (a) The proposed integration scheme during the sampling ( $\phi_1$ ) and integration ( $\phi_2$ ) phases. (b) Waveforms of a full SAR/ $\Delta\Sigma$  cycle.  $V_{in}$ : zoom-ADC's input voltage,  $V_{int}$ : integrator's output voltage. The gain factor  $k$  is set by the SAR-logic in the coarse conversion, while  $k = n-1$  or  $n+1$  when  $bs = 0$  or  $1$  in the fine conversion step.

is quite small, and can be accommodated by a low-swing, and power-efficient telescopic opamp. Moreover, this approach also halves the conversion time, thus improving the energy-efficiency by another factor of two.

### 4.2.3 Implementation

#### 4.2.3.1 Circuit Diagrams

Figure 4.15 shows the simplified circuit level diagram of the proposed  $2^{nd}$ -order zoom-ADC. A capacitor DAC with 28, 120fF unit capacitances realizes the gain factor  $k$  required for  $\Delta V_{BE}$  sampling, while an extra capacitor  $C_G = 0.5 \cdot C_S$  is used during the guard-band step. To simultaneously sample  $V_{BE}$  and in order to cover the military temperature range, the number of unit elements in the capacitor DAC is increased to 29. During the



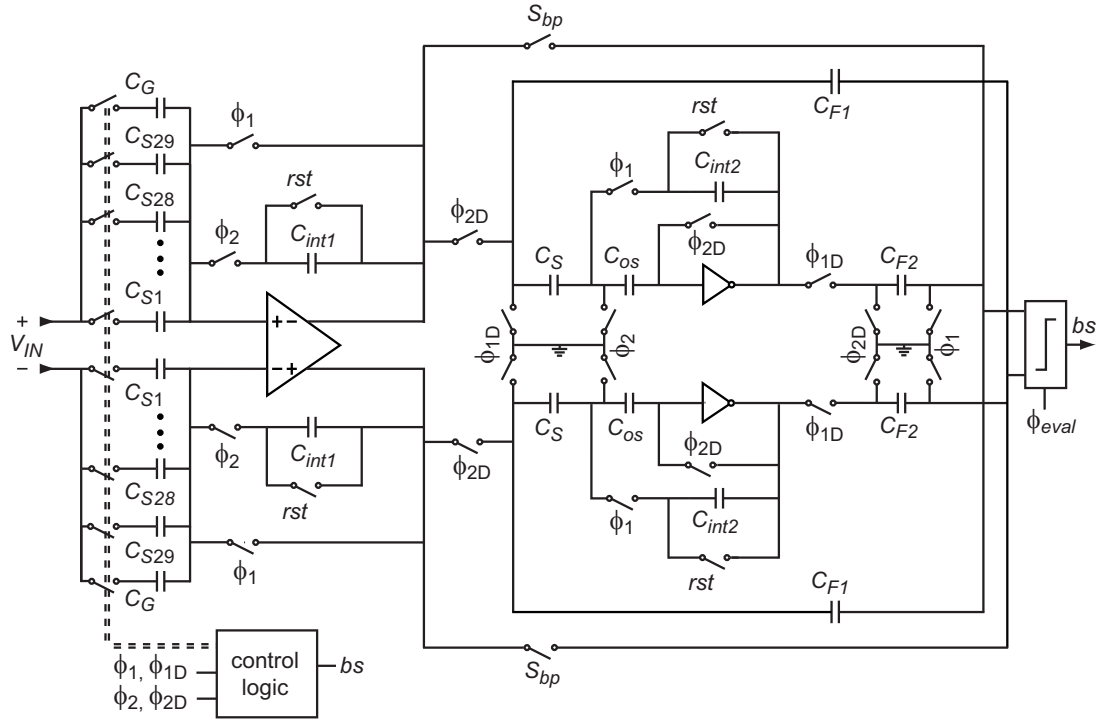


Figure 4.15: Simplified circuit diagram of the proposed 2<sup>nd</sup>-order zoom-ADC.

coarse conversion step, a switch  $S_{bp}$  bypasses the 2<sup>nd</sup> integrator, thus directly connecting the output of the first integrator to the comparator. Moreover, at the start of each comparison step, the first integrator is reset, and therefore acts as a sample-and-hold.

The first integrator is built around a power-efficient, fully differential telescopic opamp, which only draws 600nA, has a gain of 76dB, and a maximum swing of about  $\pm 200$ mV. As shown in Figure 4.15, a pseudo-differential inverter-based OTA forms the second integrator [4.11]. At 25°C, it draws 140nA, occupies only 0.002mm<sup>2</sup>, and has a gain of  $\sim 44$ dB. During  $\phi_2$ , when the output of 1<sup>st</sup> integrator is sampled on capacitors  $C_S$ , the two inverter-based OTAs are in unity-gain configuration and auto-zeroed via offset storing capacitors  $C_{OS}$ . Due to the feedback path through the integration capacitors in  $\phi_1$ , a virtual ground is formed, thus pushing the sampled charge into the integration capacitors  $C_{int2} = 2 \cdot C_S$ . Figure 4.16 shows the implementation details of the inverter-based OTA. To decrease the inverter's sensitivity to power supply and process spread, a dynamic current-biasing technique is proposed. During the auto-zeroing phase  $\phi_2$ ,  $M_{N1}$  and  $M_{P1}$  are diode-connected and biased with two current sources (45nA each), while their operating bias voltages are stored on offset storing capacitors  $C_{OS}$ . The bias voltages  $V_{b1}$  and  $V_{b2}$  are chosen such that  $M_{N2}$ ,  $M_{P2}$  are essentially off during  $\phi_2$ . The two currents are mirrored from the front-end's precision bias circuit to ensure robustness to supply and process variations. After disconnecting the two current sources in  $\phi_1$ ,  $M_{N1}$  and  $M_{P1}$  are

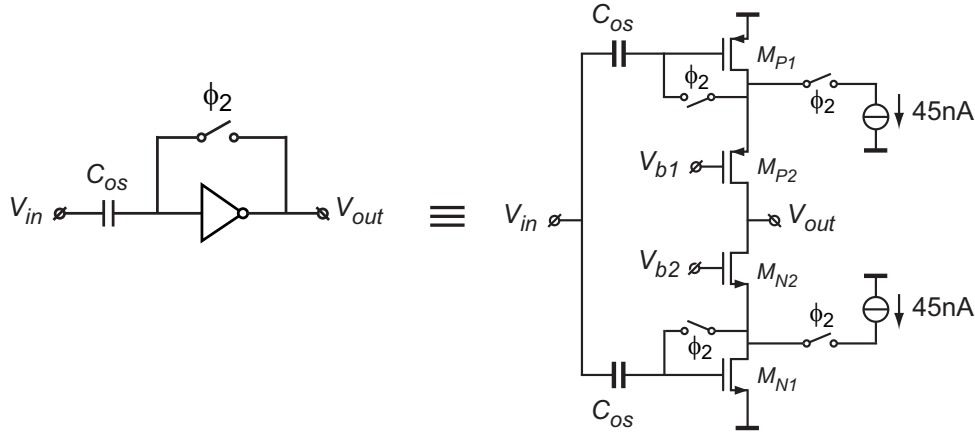


Figure 4.16: Circuit diagram of the proposed inverter-based OTA.

configured as common-source and form a class-AB amplifier, with a virtual ground at  $V_{in}$ . Since the output voltage swing requirement is reduced to about  $\pm 100\text{mV}$  by the prior coarse conversion step, cascoding of  $M_{N1}$  and  $M_{P1}$  is readily possible, thus enhancing the inverter's output resistance, and hence its DC gain. A passive summation network at the input of quantizer combines the output of  $2^{\text{nd}}$  integrator with that of the  $1^{\text{st}}$  integrator via the feed-forward capacitor  $C_{F1}$ , as shown in Figure 4.15. To set-up the various biasing voltages, the ADC requires a startup time of  $120\mu\text{s}$  (3 clock cycles) before each conversion.

### 4.2.3.2 Precision Techniques

During the fine conversion, the accuracy of the ratio  $k$  is determined by the matching between the unit capacitor that samples  $V_{BE}$  and the  $k$  capacitors which sample  $\Delta V_{BE}$ . Any mismatch will lead to a non-linear ADC transfer function. The matching of the references should, therefore, be commensurate with the ADC's target resolution, i.e. 13 bits. Since this cannot be achieved by layout alone, a dynamic element matching (DEM) scheme was used.

Figure 4.17 shows the block diagram of the sensor and the timing of a full temperature conversion. The analog front-end consists of a bias circuit and a bipolar core. As in the first prototype, the bias circuit generates a PTAT current  $I = 90\text{nA}$  (at  $25^\circ\text{C}$ ) with the help of a low power, self-biased chopped opamp and two auxiliary PNPs [4.10]. As in the first prototype, a  $\beta_F$ -compensation technique is employed to ensure that  $Q_R$  and  $Q_L$  are biased with  $\beta_F$ -independent collector currents, thus improving the robustness of the resulting  $V_{BE}$  to process spread. Furthermore, the six current sources and the two bipolar transistors in the bipolar core are dynamically matched to achieve (on average) the accurate 1 : 5 current ratio required to generate an accurate  $\Delta V_{BE}$ . As in [4.10], each

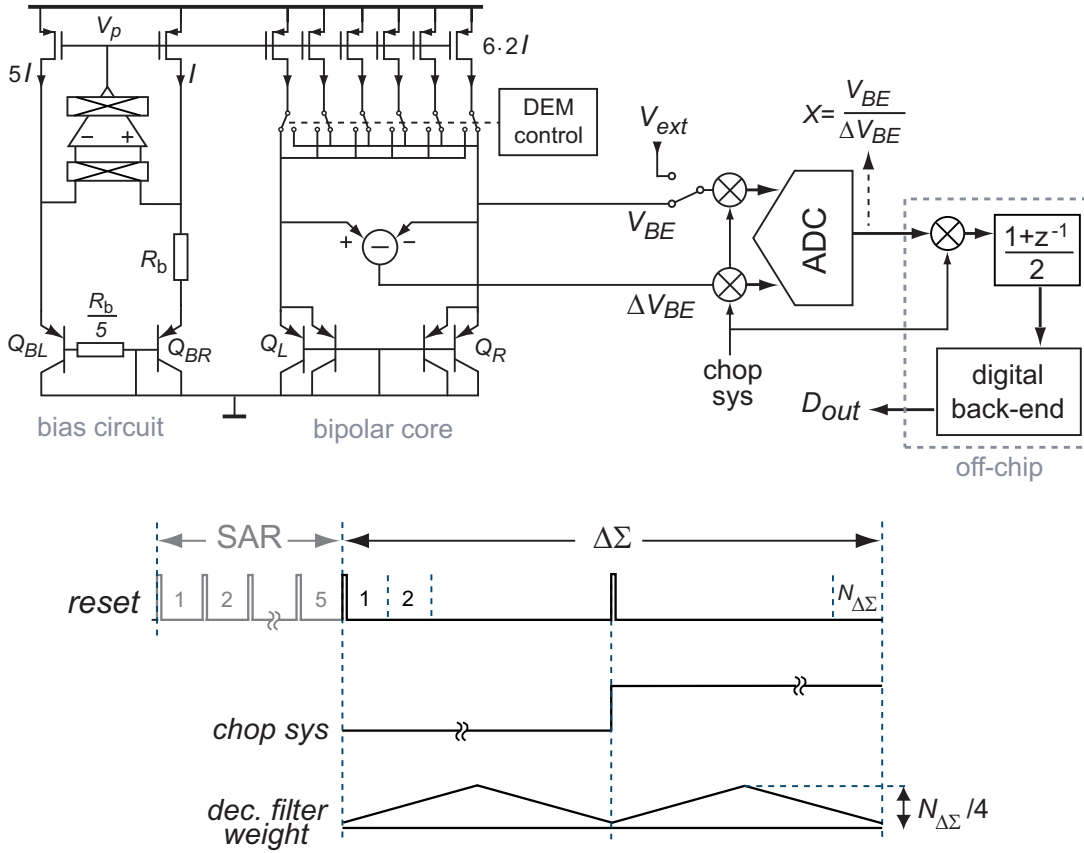


Figure 4.17: Top: Block diagram of the temperature sensor. Bottom: timing diagram of a full temperature conversion.

of the PNPs in the front-end has an emitter area of  $A_E = 5\mu\text{m} \times 5\mu\text{m}$ .

As previously discussed in section 3.1.2.2, a major source of inaccuracy, the offset and  $1/f$  noise of integrators, can be reduced by employing correlated-double sampling (CDS) during the coarse and fine conversions. In order to minimize the effect of charge injection, both integrators use differential topologies with minimum-size switches around the integration capacitors. In contrast to the first prototype, however, a digital rather than an analog implementation of system-level chopping is employed. As shown in Figure 4.17, after an initial coarse conversion, the  $\Delta\Sigma$  conversion is performed twice with swapped input voltage polarities, and the two digital results are then averaged [4.11]. This eliminates the need for state-preserving choppers around the integration capacitors, which simplifies the layout and eliminates a potential source of charge injection, which could otherwise cause ADC non-linearity. Compared to the conventional analog approach, however, this results in a small loss in resolution: up to 0.5 bits if the ADC is quantization-noise limited.

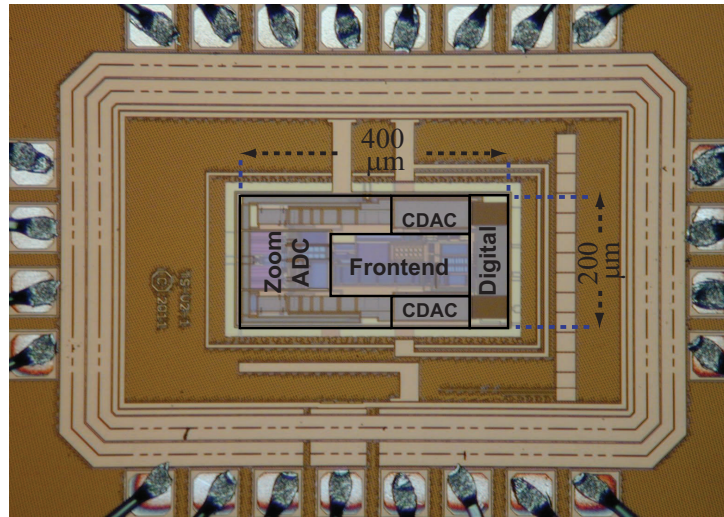


Figure 4.18: Chip micrograph of the second prototype sensor employing a  $2^{nd}$ -order zoom-ADC.

#### 4.2.4 Realization and Measurements

The sensor was realized in a standard  $0.16\mu\text{m}$  CMOS process with five metal layers, and has an active area of  $0.08\text{mm}^2$ , as shown in Figure 4.18. For flexibility, the digital back-end, the control logic and the fine conversion's  $\text{sinc}^2$  decimation filter [4.12] were implemented off-chip. At  $25^\circ\text{C}$ , the sensor draws  $3.4\mu\text{A}$  and operates from a  $1.5\text{V}$  to  $2\text{V}$  supply with a supply sensitivity of  $0.5^\circ\text{C}/\text{V}$ . Running at a clock frequency of  $25\text{kHz}$ , it requires a conversion time of  $5.3\text{ms}$  ( $128 \Delta\Sigma$  cycles) to achieve a  $kT/C$  limited resolution of  $20\text{mK}$  (rms), which improves to about  $5\text{mK}$  (rms) if the conversion time is extended to  $100\text{ms}$ . For characterization, 18 devices from one batch were packaged in ceramic DIL packages and measured over the military temperature range from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ . As shown in Figure 4.19, the resulting inaccuracy after batch calibration was  $\pm 0.6^\circ\text{C}$  ( $3\sigma$ ), with a residual curvature of only  $\pm 0.03^\circ\text{C}$ , which is significantly less than that of the first prototype ( $\pm 0.25^\circ\text{C}$ ). This was achieved by optimizing the design and layout of the Cap-DAC, in order to minimize the cross-coupling between different unit elements, thereby enhancing the overall linearity of the converter. The Cap-DAC layout is also very compact, as can be seen by comparing Figure 4.18 to Figure 4.10. To further improve the sensor's accuracy, individual calibration and trimming is essential. In the following, two different approaches based on thermal and electrical measurements are presented.

### 4.2.5 Thermal Calibration

Individual calibration of an integrated temperature sensor requires accurate information about its die temperature. Conventionally, this is obtained by bringing the device under test (DUT) and a reference temperature sensor to exactly the same temperature, whereupon the outputs of both devices are logged. As for the first prototype, the reference sensor is a platinum Pt-100 resistor calibrated to an inaccuracy of 20mK. Both sensors are embedded in a large metal block, which acts as a thermal low-pass filter and facilitates measurements with milli-Kelvin stability [4.13].

Three different single-parameter trimming methods were investigated. First, for each sensor, the offset parameter  $B$  in equation (2.9) was adjusted so as to cancel the error at the calibration temperature ( $30^\circ\text{C}$ ). After this offset trim, the sensor's inaccuracy is less than  $\pm 0.25^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ . Alternatively, the parameter  $\alpha$  in equation (3.1) can be adjusted, as was done in the first prototype [4.10]. The resulting inaccuracy, however, is almost exactly the same as that obtained with offset trim. Since the dominant source of sensor inaccuracy, i.e. the spread in  $V_{BE}$ , is PTAT in nature (see Figure 4.19), a digital PTAT trim [4.1], as discussed in section 3.4 is also employed. The resulting inaccuracy is then less than  $\pm 0.15^\circ\text{C}$  ( $3\sigma$ ), as shown in Figure 4.20.

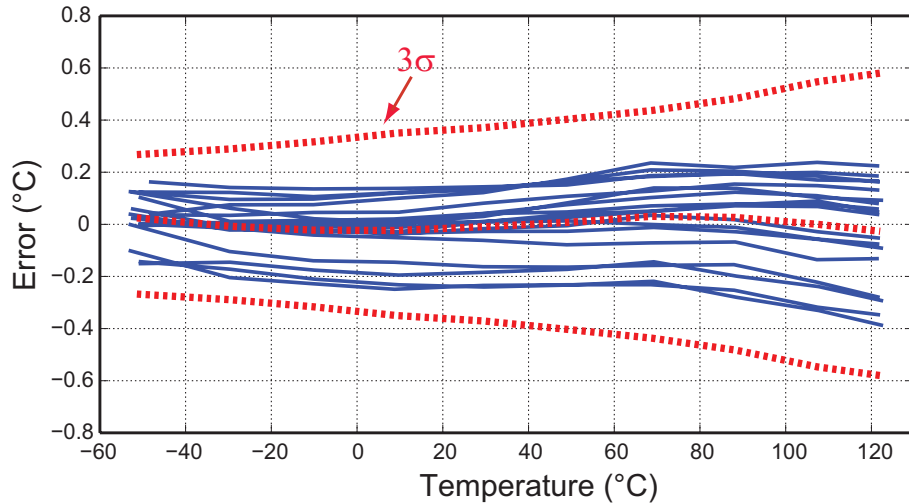


Figure 4.19: Measured temperature error of 18 sensors before trimming; dashed lines refer to the average and  $\pm 3\sigma$  limits.

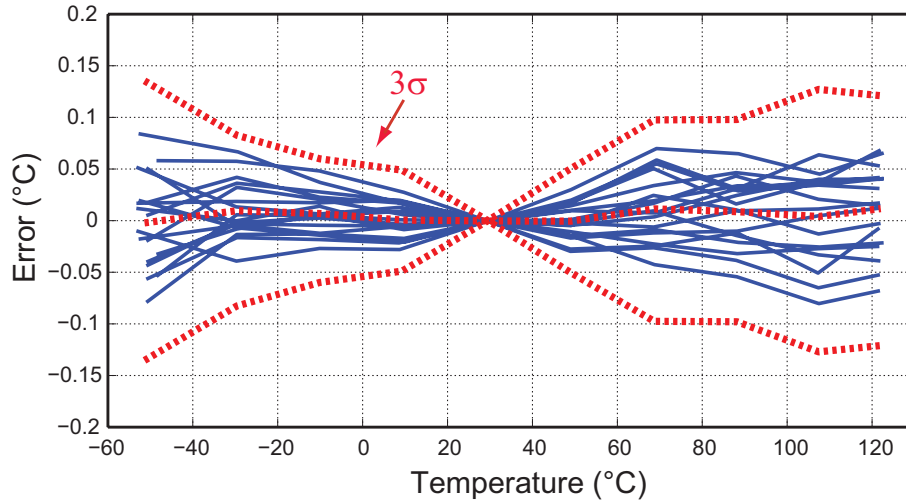


Figure 4.20: Measured temperature error of 18 sensors after thermal calibration and PTAT trimming at 30°C; dashed lines refer to the average and  $\pm 3\sigma$  limits.

### 4.2.6 Voltage Calibration

Although thermal calibration can be performed very accurately, the long stabilization time required for the DUT and the reference sensor to reach thermal equilibrium prohibits its use as a low-cost calibration method. In [4.14], a voltage calibration method was proposed, in which die temperature is established by measuring an on-chip  $\Delta V_{BE}$ . By applying DEM to the six current sources and the two PNPs (see Figure 4.17), the collector current ratio  $p$ , and therefore  $\Delta V_{BE}$  can be made robust to process spread. The process-dependent non-ideality factor  $\eta$  ( $= 1.0042$ ) can also be extracted by batch calibration. As shown in Figure 4.21, the die temperature can then be determined by the following procedure. First,  $V_{BE}$  is replaced by an accurate external voltage  $V_{ext}$  (see Figure 4.21.a). The on-chip ADC then digitizes the ratio  $X_{ext} = V_{ext}/\Delta V_{BE}$  accurately and with high resolution, whereupon the actual die temperature  $T_D$  can be calculated:

$$\Delta V_{BE} = \eta \cdot \frac{kT_D}{q} \cdot \ln(p), X_{ext} = \frac{V_{ext}}{\Delta V_{BE}} \Rightarrow T_D = \frac{V_{ext}}{C_m \cdot X_{ext}}, C_m = \eta \cdot \frac{k}{q} \cdot \ln(p). \quad (4.3)$$

In a second step,  $V_{ext}$  is replaced by the on-chip  $V_{BE}$  and a normal conversion is performed to determine  $X = V_{BE}/\Delta V_{BE}$ , and hence the sensor's untrimmed output (see Figure 4.21.b). In contrast to thermal calibration, this approach can be performed at room temperature, and is much faster, requiring only two ADC conversions. Since the sensor achieves a resolution of 5mK in a conversion time of 100ms, which is commensurate with the expected  $\pm 0.15^\circ\text{C}$  inaccuracy, this means that the total calibration time is only

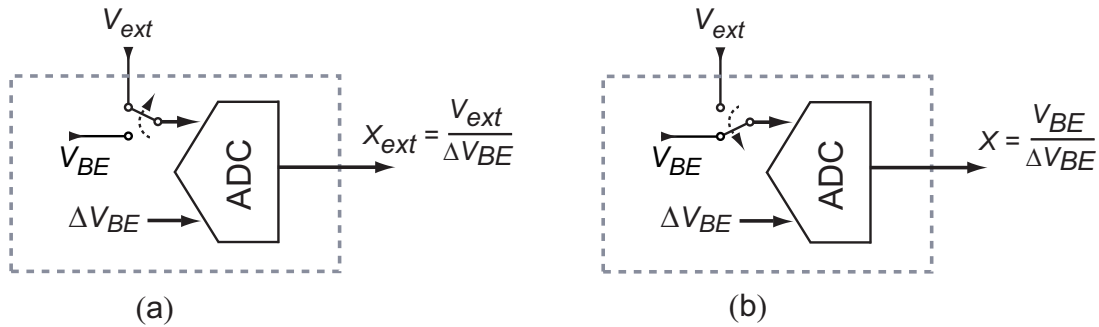


Figure 4.21: Voltage calibration requires two ADC conversions: one to obtain the actual die temperature (a) and another to obtain the untrimmed output (b).

200ms.

Compared to the results of thermal calibration, the results of voltage calibration followed by an offset or digital PTAT trim are only slightly worse around room temperature. The worst-case inaccuracy from  $-55^{\circ}\text{C}$  to  $125^{\circ}\text{C}$ , however, is almost exactly the same as shown in Figure 4.22. This confirms the fact that the inaccuracy of  $\Delta V_{BE}$  is negligible, and so voltage calibration is a robust alternative to thermal calibration.

#### 4.2.7 Batch-to-Batch Spread and Plastic Packaging

To verify the effect of batch-to-batch spread on sensor inaccuracy, devices from a different process batches were characterized. As before, 18 devices from one batch were packaged

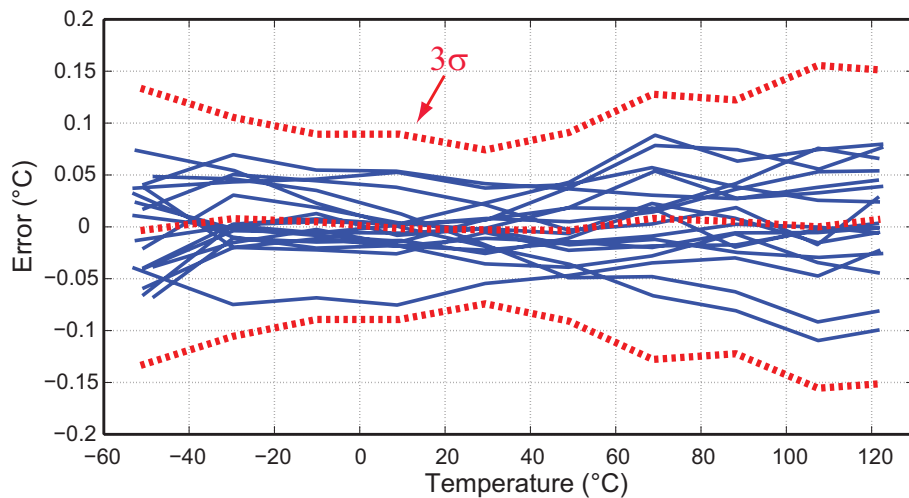


Figure 4.22: Measured temperature error of 18 sensors after voltage calibration and PTAT trimming at room temperature; dashed lines refer to the average and  $\pm 3\sigma$  limits.

in ceramic DIL packages and measured from  $-55^{\circ}\text{C}$  to  $125^{\circ}\text{C}$ . Table 4.2 compares the resulting inaccuracy and calibration parameters (i.e.  $A$ ,  $B$ , and  $\alpha$ ) of the two batches. As shown, the resulting inaccuracy after batch calibration has increased to  $\pm 0.25^{\circ}\text{C}$  ( $3\sigma$ ). Moreover, the obtained gain and offset parameters  $A$ ,  $B$  after batch calibration show a batch-to-batch spread of about 0.4% and 0.3%, respectively. Since at room temperature the PTAT ratio  $\mu \approx 0.5$ , this translates to a temperature shift of about  $-0.5^{\circ}\text{C}$ .

However, the optimal mapping coefficient  $\alpha$  changes by less than 0.1% from batch-to-batch. According to the discussion in section 3.4, such a small variation results in a minimal impact on the parameters  $A$  and  $B$ , and so  $\alpha$  can be regarded as a digital constant. Finally, the non-ideality factor  $\eta$  only changes by about 0.02% from batch-to-batch, which corresponds to a maximum calibration error of about 50mK.

In production, low-cost plastic packages are preferred to ceramic packages. The associated mechanical stress, however, impacts the sensor's accuracy, an effect which is referred to as packaging shift [4.15], and results in a fairly systematic modification to the base-emitter voltage  $V_{BE}$  [4.16], [4.17]. To evaluate this, 22 samples from the same batch of the second batch were packaged in plastic DIP packages and then characterized. As shown in Table 4.3, the untrimmed inaccuracy after batch calibration increased to about  $\pm 0.8^{\circ}\text{C}$  ( $3\sigma$ ). However, a PTAT trim reduced the inaccuracy to about  $\pm 0.25^{\circ}\text{C}$  ( $3\sigma$ ), which is equivalent to that obtained with ceramic packaging. The optimal mapping coefficient  $\alpha$  changed by about 0.2%, while the fitting parameters  $A$ ,  $B$  changed by about 0.15% and 0.35% respectively, which corresponds to a packaging shift of about  $-0.36^{\circ}\text{C}$  at room temperature.

From these measurements, it can be concluded that in order to achieve high accuracy over different batches and different packages, batch calibration is essential. Once the fitting parameters  $A$  and  $B$  are known, individual devices can be trimmed on the basis of

	<b>Batch-1</b>	<b>Batch-2</b>
<b>Untrimmed inaccuracy (<math>3\sigma</math>)</b>	$\pm 0.6^{\circ}\text{C}$	$\pm 0.6^{\circ}\text{C}$
<b>PTAT-trimmed inaccuracy (<math>3\sigma</math>)</b>	$\pm 0.15^{\circ}\text{C}$	$\pm 0.25^{\circ}\text{C}$
<b><math>\alpha</math> (mapping coefficient)</b>	15.44	15.45
<b>A (gain parameter)</b>	613.31	610.74
<b>B (offset parameter)</b>	283.70	282.93
<b><math>\eta</math> (non-ideality factor)</b>	1.0042	1.0044

Table 4.2: Impact of batch-to-batch spread on sensor accuracy and calibration parameters.



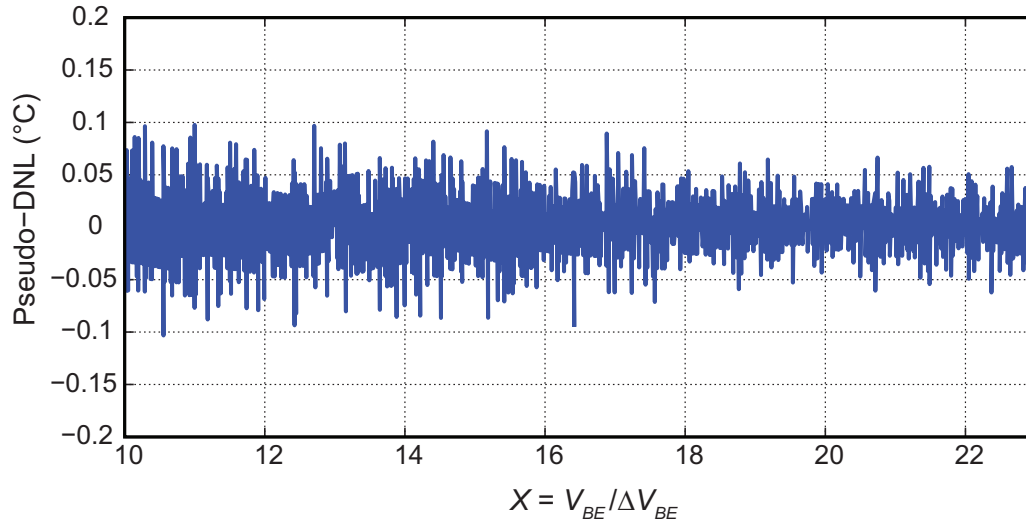


Figure 4.23: Measured Pseudo-DNL versus  $X$ . The sensor's conversion time is 5.3ms.

a fast voltage calibration, since the non-ideality factor appears to be essentially constant over different batches and packages.

#### 4.2.8 Noise and ADC Characteristics

As in other two-step ADC structures, mismatch between the references used in the various fine conversion steps could result in discontinuities in the ADC's characteristic. To examine this, the ADC's input range was swept by slowly sweeping the oven temperature from  $-40^{\circ}\text{C}$  to  $100^{\circ}\text{C}$  over a three hour period, while continuously logging the sensor's output. This corresponds to a temperature slope of  $\approx 13\text{mK/s}$ , which implies that the

	<b>Ceramic Package</b>	<b>Plastic Package</b>
<b>Untrimmed inaccuracy (<math>3\sigma</math>)</b>	$\pm 0.6^{\circ}\text{C}$	$\pm 0.8^{\circ}\text{C}$
<b>PTAT-trimmed inaccuracy (<math>3\sigma</math>)</b>	$\pm 0.25^{\circ}\text{C}$	$\pm 0.25^{\circ}\text{C}$
<b><math>\alpha</math> (mapping coefficient)</b>	15.45	15.47
<b>A (gain parameter)</b>	610.74	611.59
<b>B (offset parameter)</b>	282.93	283.94
<b><math>\eta</math> (non-ideality factor)</b>	1.0044	1.0044

Table 4.3: Effect of mechanical stress on sensor accuracy and calibration parameters.

Parameter	2 <sup>nd</sup> Prototype JSSC'13 [4.18]	1 <sup>st</sup> Prototype JSSC'11 [4.10]	ISSCC'14 [4.19]	JSSC'05 [4.20]	ISSCC'09 [4.21]
Sensor type	PNP	PNP	NPN	Resistor	MOSFET
CMOS technology	0.16 $\mu$ m	0.16 $\mu$ m	0.18 $\mu$ m	0.18 $\mu$ m	0.18 $\mu$ m
Chip area	0.08mm <sup>2</sup>	0.12mm <sup>2</sup>	0.085mm <sup>2</sup>	0.18mm <sup>2</sup>	0.032mm <sup>2</sup>
Supply current (RT) <sup>†</sup>	3.4 $\mu$ A	4.6 $\mu$ A	4.5 $\mu$ A	20 $\mu$ A	0.4 $\mu$ A
Supply voltage	1.5V - 2V	1.6V - 2V	1.4V - 2V	1.2V - 2V	0.9V - 1.1V
Supply sensitivity	0.5°C/V	0.1°C/V	-	0.625°C/V	8°C/V
Inaccuracy	$\pm 0.15^\circ\text{C}$ ( $3\sigma$ )	$\pm 0.2^\circ\text{C}$ ( $3\sigma$ )	-	$\pm 0.5^\circ\text{C}$ (max)	-0.8 / +1°C (max)
Temperature range	-55°C - 125°C	-30°C - 125°C	-45°C - 85°C	0°C - 100°C	0°C - 100°C
Calibration (points)	voltage (1)	thermal (1)	-	thermal (1)	thermal (2)
Resolution ( $T_{\text{conv}}$ )	0.02°C (5.3msec)	0.015°C (100msec)	0.025°C (6msec)	0.25°C (12.5 $\mu$ sec)	0.3°C (1msec)
Resolution FoM	11pJ°C <sup>2</sup>	170pJ°C <sup>2</sup>	24pJ°C <sup>2</sup>	19pJ°C <sup>2</sup>	32pJ°C <sup>2</sup>

<sup>†</sup> Excluding the off-chip digital

Table 4.4: Performance summary and comparison to previous work.

temperature change between successive measurements is less than 1mK, i.e. much smaller than the sensor's own resolution. Taking the difference between successive sensor outputs then results in a pseudo-DNL function, which reflects the ADC's resolution and possible discontinuities between the various fine conversion segments. As shown in Figure 4.23, the sensor achieves a resolution of 20mK (rms) into 5.3ms around room temperature ( $X \approx 14.5$ ), which is enough to calibrate it rapidly to 0.2°C inaccuracy. Moreover, there are no discontinuities between the different fine segments. Lastly, it can be seen that the sensor's resolution is slightly temperature dependent. This is due to the fact that the full-scale range of each fine conversion is not constant, but is equal to  $2 \cdot \Delta V_{BE}$ .

### 4.2.9 Comparison to Previous Work

The sensor's performance is summarized in Table 4.4 and compared to the first prototype as well as to other energy-efficient, low power state-of-the-art temperature sensors, with

digital output. It employs a low-cost, room-temperature voltage calibration technique, and it also achieves the highest accuracy over a wide temperature range:  $\pm 0.15^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ . Compared to the first prototype, this work achieves comparable resolution in about 18x less conversion time, while consuming 25% less supply current, thus improving the energy-efficiency by over 20x. This is also in line with the performance of state-of-the-art thermistor- and MOSFET-based sensors, and is evidenced by a resolution FoM of  $11\text{pJ}^\circ\text{C}^2$ , which is the best among others.

### 4.3 Sensing High Temperatures <sup>3</sup>

While most applications only require temperature sensors that operate up to  $125^\circ\text{C}$ , automotive and industrial applications require operation at much higher temperatures ( $>150^\circ\text{C}$ ). Achieving this in CMOS is not a trivial task and so off-chip sensors such as thermistors and thermocouples are often used. The main challenge lies in the fact that the leakage and saturation currents of CMOS components increase rapidly at high temperatures, which, in turn, leads to significant temperature-sensing errors. While dissipating  $65\mu\text{W}$ , the sensor in [4.22] achieves an inaccuracy of  $\pm 0.1^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$  after a one-point trim. However, its inaccuracy increases to about  $\pm 1^\circ\text{C}$  at  $150^\circ\text{C}$ . While drawing  $180\mu\text{A}$  and after a two-point trim, the sensor in [4.23] achieves an inaccuracy of  $\pm 0.5^\circ\text{C}$  and  $\pm 2.5^\circ\text{C}$  from  $-55^\circ\text{C}$  to  $175^\circ\text{C}$  in the current and voltage modes of operation, respectively. In a commercially available product, an inaccuracy of  $\pm 1^\circ\text{C}$  was achieved up to  $175^\circ\text{C}$  (after trimming) [4.24]. In another product, an inaccuracy of  $\pm 3^\circ\text{C}$  was achieved at  $200^\circ\text{C}$  [4.25].

As explained in Chapter 1, CMOS temperature sensors based on the thermal diffusivity (TD) of silicon generate a temperature-dependent phase-shift, and are thus insensitive to leakage currents. This allows them to operate up to very high temperatures without compromising their accuracy. In [4.26], an inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) from  $-70^\circ\text{C}$  to  $200^\circ\text{C}$  was achieved with a one-point trim. However, the sensor's operation requires the generation of heat pulses, resulting in a power dissipation of  $2.6\text{mW}$ .

In this section, a modified version of the second prototype is presented, which achieves an inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $200^\circ\text{C}$  after a one-point trim, while drawing only  $22\mu\text{A}$ . It achieves a resolution FoM of  $59\text{pJ}^\circ\text{C}^2$ , which is much higher than that of the second prototype, mainly due to the modifications necessary for high temperature operation.

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<sup>3</sup>K. Souri, K. Souri, and K. A. A. Makinwa, "A  $40\mu\text{W}$  CMOS Temperature Sensor with an Inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $200^\circ\text{C}$ ," in *Proc. ESSCIRC*, pp. 221 - 224, Sept. 2013.

### 4.3.1 Analog Front-End

Recall from Chapter 1, the base-emitter voltage  $V_{BE}$  of a PNP can be expressed as:

$$V_{BE} = \frac{kT}{q} \ln \left( \frac{I_C}{I_S} + 1 \right), \quad (4.4)$$

where  $k$  is the Boltzmann constant,  $q$  is the electron charge,  $T$  is the temperature in Kelvin,  $I_C$  is the collector current and  $I_S$  is the saturation current. The base-emitter voltage difference of two identical PNPs biased at a 1 :  $p$  collector current ratio can then be expressed as:

$$\Delta V_{BE} = V_{BE2} - V_{BE1} = \frac{kT}{q} \ln \left( \frac{p \cdot I_C + I_S}{I_C + I_S} \right). \quad (4.5)$$

When operation at very high temperatures is targeted, the temperature dependence of the saturation current  $I_S$  must be taken into account. The saturation current  $I_S$  is given by:

$$I_S(T) = CT^\eta \exp \left( \frac{-qV_{g0}}{kT} \right), \quad (4.6)$$

where  $C$  and  $\eta$  are constants, while  $V_{g0}$  is the extrapolated bandgap voltage at 0K. The constant  $C \propto A_E$ , where  $A_E$  is the BJT's emitter area, while  $\eta \approx 4$ , implying that  $I_S$  is strongly temperature dependent. The sensitivities of  $V_{BE}$  and  $\Delta V_{BE}$  to temperature variations can be determined from the derivatives:

$$S_{V_{BE}}^T = \frac{\partial V_{BE}}{\partial T}, S_{\Delta V_{BE}}^T = \frac{\partial \Delta V_{BE}}{\partial T}. \quad (4.7)$$

Figure 4.24 shows the simulated temperature dependency of  $S_{V_{BE}}^T$  and  $S_{\Delta V_{BE}}^T$  for various emitter areas and for a PTAT bias current  $I_b = 90\text{nA}$  (at  $25^\circ\text{C}$ ) as in [4.18]. For  $I_S \ll I_C$ ,  $S_{V_{BE}}^T$  and  $S_{\Delta V_{BE}}^T$  are near-constant, monotonic functions of temperature. At high temperatures, however, this condition no longer holds, leading to an inflection point in  $S_{V_{BE}}^T$  and  $S_{\Delta V_{BE}}^T$ . This means that the spread of  $I_S$ , will introduce extra temperature-sensing errors that cannot be simply corrected by a one-point trim. Since  $C \propto A_E$ , the inflection point shifts to higher temperatures as  $A_E$  decreases. Alternatively,  $I_C$  can be boosted in order to reduce the impact of increases in  $I_S$ .

The sensor's analog front-end is shown in Figure 4.25. Two PNPs, biased at a 5 : 1 current ratio, generate an accurate bias current  $I_b$ , which is then used to bias the

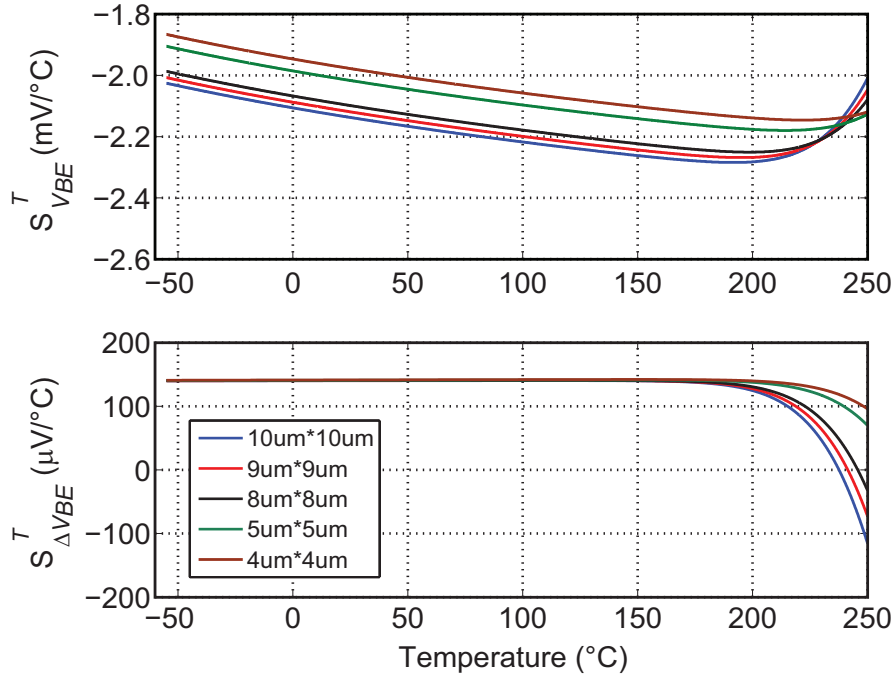


Figure 4.24:  $S_{V_{BE}}^T$  and  $S_{\Delta V_{BE}}^T$  versus temperature for different emitter area  $A_E$  and for the bias current  $I_b = 90\text{nA}$ .

PNPs ( $A_E = 4\mu\text{m} \times 4\mu\text{m}$ ) of the bipolar core. In order to investigate the effect of bias current on sensor accuracy [4.22], three programmable bias currents can be generated by appropriately connecting the bias resistors  $R_b$ . A set of three resistors (each equal to  $R_b/5$ ) in series with the base of  $Q_{BL}$  are used to implement the  $\beta_F$ -compensation, as in [4.2].

As in [4.18], the impact of the opamp's offset  $V_{OS}$  is reduced by chopping it twice per conversion. In [4.18], a positive feedback opamp topology was used. At high temperatures, however, leakage currents severely impact its gain. In this work, a robust folded-cascode topology is used. Moreover, DEM is applied to mitigate the mismatch of the current sources and PNPs. At high temperatures, the leakage current and the  $R_{off}$  of the DEM switches will also modify  $p$  (see equation (4.5)), leading to extra error in  $\Delta V_{BE}$  [4.27]. To mitigate this, the DEM switches are realized with thick oxide NMOS devices and sized appropriately.

### 4.3.2 ADC Design

As in [4.18], an energy-efficient  $2^{nd}$ -order zoom-ADC was used to digitize the ratio  $X = V_{BE}/\Delta V_{BE}$ . As previously discussed, since the fine  $\Delta\Sigma$ -ADC only processes the quanti-

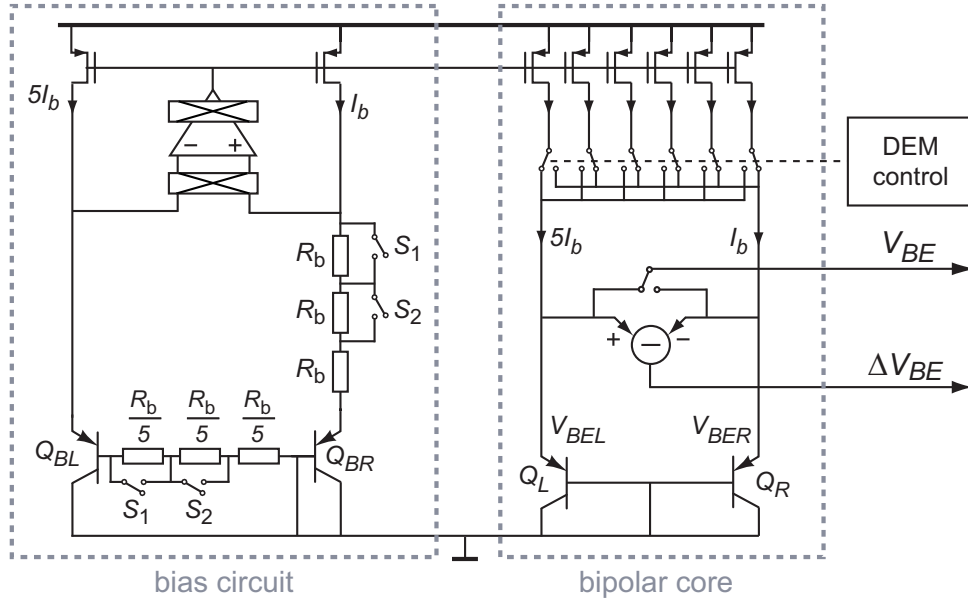


Figure 4.25: Circuit diagram of the sensor front-end.

zation error of the coarse phase, the output swing of the two integrators is quite small: less than  $\pm 3 \cdot \Delta V_{BE}$  ( $\pm 125\text{mV}$  at  $25^\circ\text{C}$ ). In [4.18], therefore, a telescopic opamp was used in the first integrator while the second integrator used a pair of inverter-based OTAs. At  $200^\circ\text{C}$ , however, the swing increases to about  $\pm 200\text{mV}$ . Moreover, leakage currents then reduce the gain of the inverter-based OTAs significantly. In this work, therefore, both integrators were implemented as folded-cascode OTAs, drawing  $1\mu\text{A}$  each.

Figure 4.26.a shows a pair of unit elements of the capacitor DAC that are used in the sampling network of the ADC, described in section 4.2.3.1 (see Figure 4.15). Each of the unit capacitors is connected to the bipolar core via two switches  $S_1$  and  $S_2$ . Each pair of capacitors samples either  $V_{BE}$  ( $S_1$  on) or  $\Delta V_{BE}$  ( $S_2$  on) or none (both  $S_1$  and  $S_2$  off), depending on the conversion state. At temperatures above  $\approx 120^\circ\text{C}$ ,  $X$  is less than 6, implying that most of the DAC switches are in the off-state. At these temperatures, however, the leakage current of these switches impacts the current ratio  $p$ , while their lower off-resistances causes charge integration errors. This is mitigated by using thick oxide transistors and modifying the switching scheme as shown in Figure 4.26.b. When  $S_1$  and  $S_2$  are in the off-state, an extra switch  $S_3$  ties the nodes  $A$  and  $B$  to an auxiliary base-emitter voltage  $V_{BE,aux}$ . This configuration will decrease the voltage across the off switches to either  $(V_{BER} - V_{BE,aux})$  or  $(V_{BEL} - V_{BE,aux})$ , which is small, and thereby significantly reducing the leakage currents through the off switches. It also prevent the nodes  $A$  and  $B$  from floating and absorbing any leakage currents.

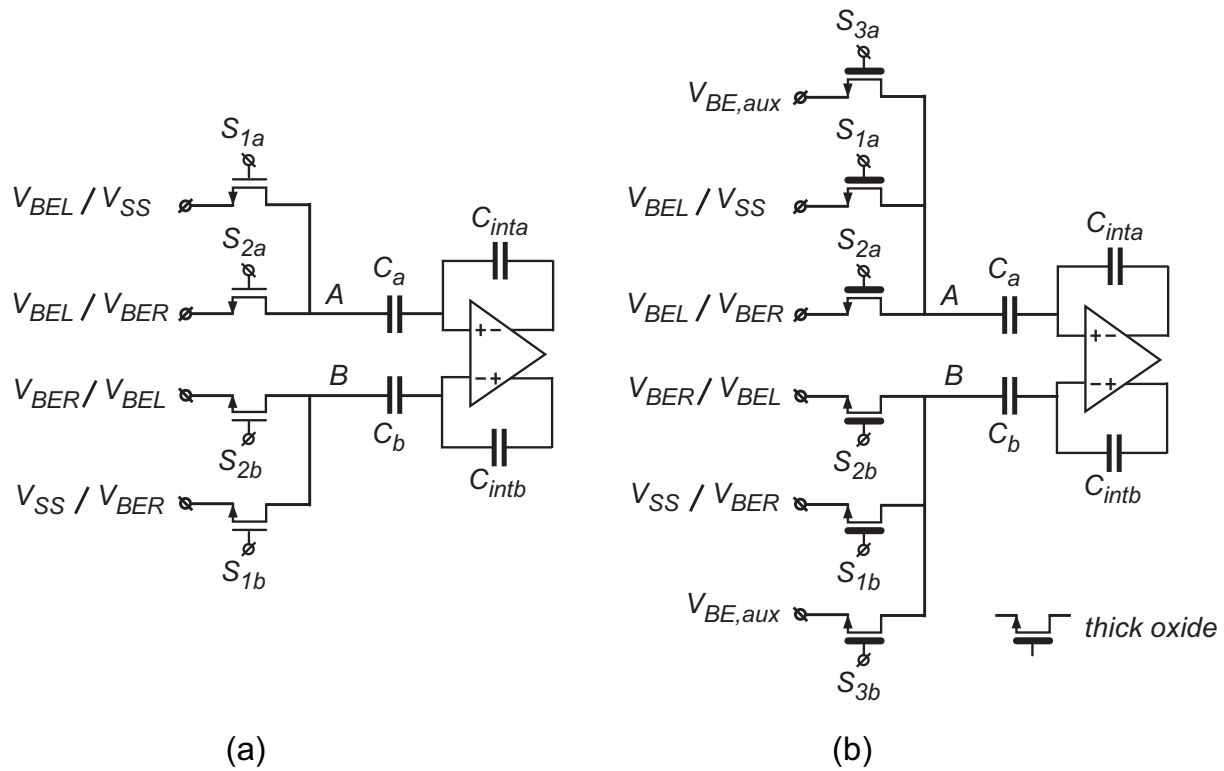


Figure 4.26: A typical implementation of unit elements in a capacitor DAC (a), and the modified switching scheme to cope with the leakage currents and  $R_{off}$  at high temperatures (b).

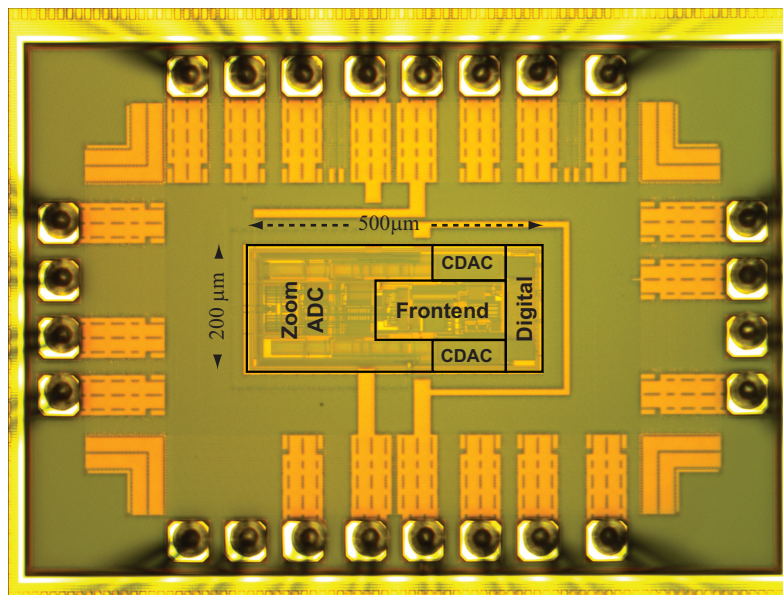


Figure 4.27: Chip micrograph of the sensor for high temperature sensing.

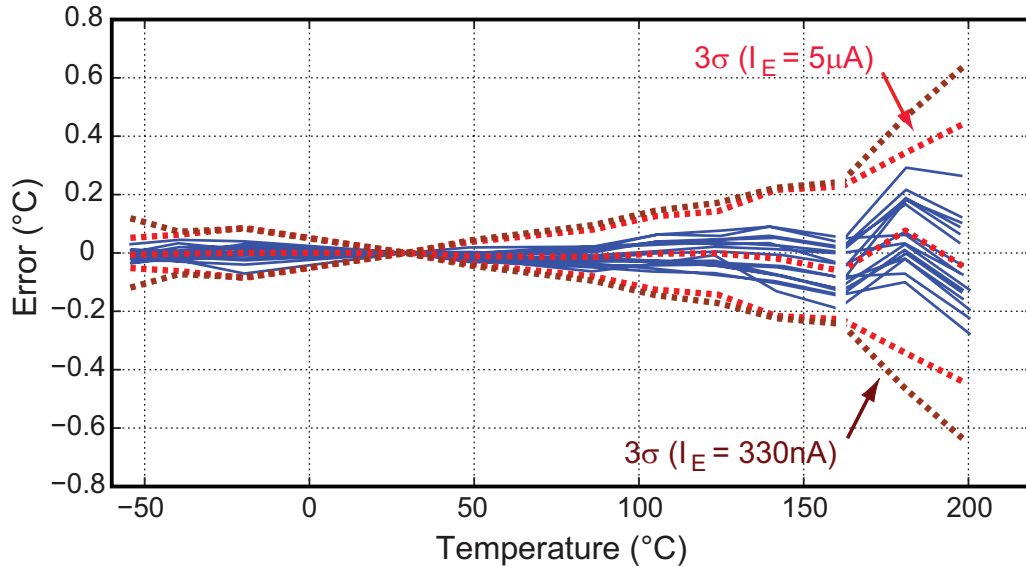


Figure 4.28: Temperature error of 16 devices after trimming at 30°C.

### 4.3.3 Measurement Results

The  $0.1\text{mm}^2$  (active area) sensor was realized in a  $0.16\mu\text{m}$  CMOS-SOI process (Figure 4.27), in a standard CMOS island, i.e. no SOI-specific features were used. For flexibility, the digital back-end, the control logic and the  $\Delta\Sigma$ -ADC's  $\text{sinc}^2$  decimation filter were implemented off-chip. 16 devices in ceramic DIL packages were characterized from  $-55^\circ\text{C}$  to  $200^\circ\text{C}$ . The sensor's accuracy was measured at three different values of  $I_b$ :  $330\text{nA}$ ,  $660\text{nA}$ , and  $1\mu\text{A}$ . Moreover, either  $Q_R$  or  $Q_L$ , which are biased at  $I_b$  and  $5 \times I_b$ , respectively, can be used to generate  $V_{BE}$  (Figure 4.25). The emitter current  $I_E$  at which  $V_{BE}$  is generated can, therefore, be programmed between  $330\text{nA}$  (min) and  $5\mu\text{A}$  (max). As expected, when  $V_{BE}$  is generated by  $I_E = 5\mu\text{A}$ , the impact of saturation and leakage currents is minimal, leading to an inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) after a one-point trim at  $30^\circ\text{C}$ , as shown in Figure 4.28. The inaccuracy at high temperatures gets worse as the current generating  $V_{BE}$  decreases. For a current of  $I_E = 330\text{nA}$ , the sensor's inaccuracy increases rapidly above  $160^\circ\text{C}$ , and is  $\pm 0.6^\circ\text{C}$  ( $3\sigma$ ) at  $200^\circ\text{C}$ , as shown in Figure 4.28. The  $180^\circ\text{C}$  and  $200^\circ\text{C}$  temperature points were measured in another setup with a different climate chamber, which may explain the discontinuity between the two sets of measurements. At  $25^\circ\text{C}$ , the sensor draws  $22\mu\text{A}$  from a  $1.6\text{V}$  to  $2\text{V}$  supply, and achieves a  $kT/C$  limited resolution of  $20\text{mK}$  (rms) in a conversion time of  $4.2\text{ms}$ . The sensor's performance is summarized in Table 4.5 and compared to that of other sensors capable of operating at high temperatures. Its inaccuracy is comparable to that of state-of-the-art TD sensor in [4.26], while dissipating much less power.



Parameter	This work ESSCIRC'13 [4.28]	ISSCC'12 [4.26]	LM95172 [4.25]	ADT7312 [4.24]
Chip area	0.1mm <sup>2</sup>	1mm <sup>2</sup>	-	-
Supply current (RT) <sup>†</sup>	22μA	520μA	500μA	245μA
Inaccuracy (trim points)	±0.4°C (3σ) (1)	±0.4°C (3σ) (1)	±0.4°C (trimmed)	±1°C (trimmed)
Temperature range	-55°C - 200°C	-70°C - 200°C	-40°C - 200°C	-55°C - 175°C
Resolution (T <sub>conv</sub> )	0.02°C (4.2msec)	0.075°C (1.4sec)	0.007°C (350msec)	0.007°C (240msec)

<sup>†</sup> Excluding the off-chip digital

Table 4.5: Performance summary and comparison with previous work.

## 4.4 Conclusions

In this chapter, various BJT-based sensor prototypes based on zoom-ADCs were demonstrated. By combining the benefits of SAR and  $\Delta\Sigma$ -ADCs, an accurate, energy-efficient temperature to digital conversion was obtained. By employing a 1<sup>st</sup>-order zoom-ADC in the first prototype, a low-power sensor was realized, drawing only 4.6μA. Since the ADC's fine conversion step was based on a slow, 1<sup>st</sup>-order  $\Delta\Sigma$  modulator, it required 100ms to achieve 15mK resolution, limiting its energy-efficiency. Moreover, a less power-efficient folded-cascode opamp was used to accommodate the large swing and gain requirements of the loop-filter. Employing a thermal calibration and a single  $\alpha$ -trim at 25°C, the sensor achieved an accuracy of ±0.2°C (3σ) over the temperature range from -30°C to 125°C.

To improve the sensor's energy-efficiency, a second prototype was realized which achieves similar resolution in about 16x less conversion time, while drawing 25% less supply current. This was achieved by using a 2<sup>nd</sup>-order zoom-ADC, combined with a new charge-balancing scheme. Simultaneous sampling of  $V_{BE}$  and  $\Delta V_{BE}$  also allowed the use of low-swing, and therefore low-power amplifiers, thus further improving the sensor's energy-efficiency. By using a 2<sup>nd</sup>-order modulator, the required loop-gain to maintain linearity between different fine conversion segments was readily achieved. The sensor's energy-efficiency was therefore improved by over 20x compared to the first prototype. Using thermal calibration and digital PTAT trimming at 30°C, the sensor achieved an

accuracy of  $\pm 0.15^\circ\text{C}$  ( $3\sigma$ ) over the military temperature range: from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ . To meet the extreme cost constraints on large volume sensor products, a voltage calibration technique based on electrical measurements was also explored. Compared to thermal calibration, it is significantly faster, requiring only two ADC conversions (200ms), while achieving comparable accuracy. Moreover, the impact of batch-to-batch spread and plastic packaging on sensor's accuracy was explored. As observed, both batch-to-batch spread and plastic packaging can cause temperature reading shifts in the order of  $0.4^\circ\text{C} \sim 0.5^\circ\text{C}$  over the military temperature range.

Last but not least, a BJT-based temperature sensor for very high temperature sensing was also implemented. To mitigate the impact of leakage and BJT's saturation current at high temperatures, the bias current and the emitter area of BJTs ( $A_E = 4\mu\text{m} \times 4\mu\text{m}$ ) were carefully optimized and robust circuit-level techniques were used. After a one-point trim at  $30^\circ\text{C}$ , the sensor achieved an accuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $200^\circ\text{C}$ , which is similar to that of state-of-the-art thermal diffusivity (TD) sensors, but it draws more than an order of magnitude less current.

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## Chapter 5

# All-CMOS Precision Temperature Sensors

In CMOS technology, BJT-based sensors are usually the temperature sensors of choice due to their decent accuracy after a single-temperature trimming, e.g.  $\pm 0.2^\circ\text{C}$  ( $3\sigma$ ) over the military temperature range:  $-55^\circ\text{C}$  to  $125^\circ\text{C}$  [5.1] - [5.3]. They also achieve low supply-sensitivity, typically in the order of  $0.1^\circ\text{C}/\text{V}$ . However, BJT-based sensors typically require supply voltages above 1V, since  $V_{BE}$  will be about 0.8V at  $-55^\circ\text{C}$  and some headroom is required for the current source (often cascoded) that biases the BJT. This restricts the use of such sensors in battery-powered systems, and also restricts the temperature range of implementations in nanometer CMOS [5.4]. However, from an energy-efficiency perspective, a lower  $V_{DD}$  value is preferred.

To achieve sub-1V operation, MOSFET-based temperature sensors have been proposed [5.5]. When biased in the sub-threshold region, the drain current  $I_D$  and the gate-source voltage  $V_{GS}$  of a MOSFET exhibit a temperature-dependent exponential relationship, similar to that between the collector current  $I_C$  and  $V_{BE}$  of a BJT. As a result, MOSFETs can also be used as temperature sensing elements. At  $-55^\circ\text{C}$ , however,  $V_{GS}$  will only be a few hundred millivolts, which makes sub-1V operation possible. Unfortunately, process spread affects two parameters of a MOSFET: the threshold voltage  $V_T$  and the charge mobility  $\mu$ . As a result, such sensors exhibit greater inaccuracy than BJT-based sensors, e.g.  $-1.8/+1^\circ\text{C}$  from  $10^\circ\text{C}$  to  $80^\circ\text{C}$  after a single-temperature trim [5.5]. The propagation delay of a CMOS inverter can also be used as a measure of temperature. Compared to BJT-based sensors, however, such sensors suffer from much greater supply sensitivity (about  $10^\circ\text{C}/\text{V}$ ) and greater inaccuracy ( $-0.4/+0.6^\circ\text{C}$ ) over a limited range from  $0^\circ\text{C}$  to  $90^\circ\text{C}$ , even after a more-costly two-temperature trim [5.6].

This chapter describes the design of two temperature sensor prototypes based on dynamic threshold MOS transistors (DTMOSTs) [5.7] - [5.9]. By using the DTMOST configuration, i.e. by connecting the body of a MOSFET to its gate, a near-ideal diode

characteristic can be realized. The resulting device can then replace the BJTs used in precision temperature sensors, e.g. sensors described in Chapter 4. In the following, the design of a first prototype DTMOST-based sensor is described. After a two-temperature trim, the sensor achieves an inaccuracy of  $\pm 0.1^\circ\text{C}$  ( $3\sigma$ ) over the military temperature range ( $-55^\circ\text{C}$  to  $125^\circ\text{C}$ ). This represents a 5x improvement in accuracy over previously reported MOSFET-based temperature sensors. Employing a one-point trim increases the inaccuracy to  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ), which is only a factor 2x worse compared to, similarly trimmed, BJT-based sensors. In a second prototype, the low-voltage capability of DTMOSTs is then exploited to realize a sub-1V, sub- $\mu\text{W}$  precision sensor.

## 5.1 DTMOSTs as Sensing Element <sup>1</sup>

A DTMOST is a standard MOSFET whose body and gate terminals are connected together. This causes its threshold voltage to vary dynamically, hence the name [5.7]. Compared to a diode-connected MOSFET, the  $V_{GS} - I_D$  characteristic of a diode-connected DTMOST, i.e. a DTMOST diode, is less sensitive to  $V_T$  spread [5.7] - [5.9]. In this section, a first DTMOST-based sensor prototype will be presented, which after a single-temperature trim, achieves an inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) over the military temperature range, and improves to  $\pm 0.1^\circ\text{C}$  ( $3\sigma$ ) after a two-temperature trim. Compared to state-of-the-art MOSFET-based sensors [5.5], [5.6], this represents a 5x improvement in accuracy.

### 5.1.1 Operating Principle

The use of a DTMOST as a temperature sensing element was described in Chapter 1. Like the base-emitter voltage  $V_{BE}$  of a BJT, the gate-source voltage  $V_{GS}$  of a DTMOST diode exhibits complementary-to-absolute temperature (CTAT) behavior, but with a smaller temperature coefficient:  $\approx -1\text{mV}/^\circ\text{C}$ . In a similar manner, the difference in the gate-source voltage  $\Delta V_{GS}$  of two DTMOST diodes is PTAT:  $\Delta V_{GS} = (kT/q) \cdot \ln(p)$ , where  $p$  is the bias current ratio and  $kT/q$  is the thermal voltage. Furthermore, a DTMOST has an effective bandgap voltage of roughly 0.6V [5.7], i.e. about half the bandgap voltage of a BJT. This suggests that DTMOSTs can be used to implement sub-1V bandgap voltage references [5.7], [5.10] and temperature sensors in standard CMOS. Since DTMOSTs are just alternatively connected MOSFETs, such circuits can be designed and optimized with the help of standard compact models. This is another advantage over parasitic BJTs,

<sup>1</sup>K. Souri, Y. Chae, Y. Ponomarev and K.A.A. Makinwa, "A Precision DTMOST-Based Temperature Sensor," in *Proc. ESSCIRC*, pp. 279 - 282, Sept. 2011.

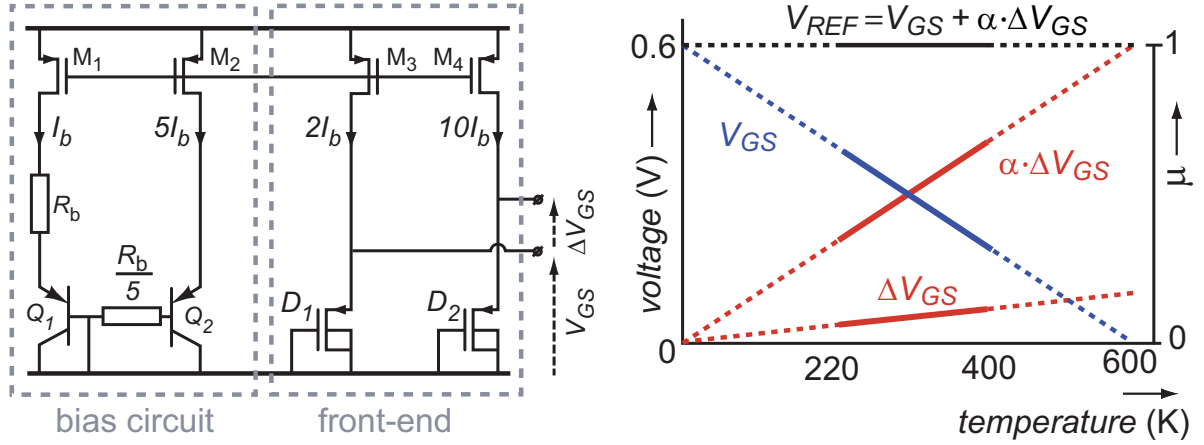


Figure 5.1: Temperature dependence of  $V_{GS}$  and  $\Delta V_{GS}$ , generated by two DTMOST diodes biased in sub-threshold region.

which require other models, and which are often not very accurately characterized. Last but not least, the gate-source voltage  $V_{GS}$  of a DTMOST exhibits less spread, about half that of a normal MOSFET [5.7], [5.8], which is a promising feature to realize accurate temperature sensors.

### 5.1.2 Temperature Sensor Design

Figure 5.1 shows the CTAT and PTAT behavior of the  $V_{GS}$  and  $\Delta V_{GS}$  voltages of a DTMOST diode. As in a bandgap voltage reference, a reference voltage can be derived from a linear combination of  $\Delta V_{GS}$  and  $V_{GS}$  [5.7]:  $V_{REF} = V_{GS} + \alpha \cdot \Delta V_{GS}$ , where  $\alpha$  is the fixed gain factor required to obtain a nominally zero temperature coefficient. As in a bandgap (BJT-based) temperature sensor [5.1], an analog-to-digital converter (ADC) can then digitize the PTAT ratio  $\mu$  between the PTAT voltage  $\alpha \cdot \Delta V_{GS}$  and  $V_{REF}$ , to obtain:

$$\mu = \frac{\alpha \cdot \Delta V_{GS}}{V_{GS} + \alpha \cdot \Delta V_{GS}}, \quad (5.1)$$

where  $\mu$  varies between 0 and 1. For the bias current ratio  $p = 5$  used in this design,  $\alpha \approx 8$ . A linear scaling of  $\mu$  then results in a digital output in degrees Celsius:  $D_{out} = A \cdot \mu + B$ , where  $A \approx 600\text{K}$ , and  $B \approx -273\text{K}$ , as presented in equation (2.9).

As shown in the previous chapter, the ratio  $X = V_{GS}/\Delta V_{GS}$  can now be used as a measure of temperature. For  $p = 5$ , the resulting non-linear ratio  $X = V_{GS}/\Delta V_{GS}$  ranges between 3 and 12 for temperatures ranging from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$ . This ratio can then be digitized by an ADC, and the PTAT ratio  $\mu$  is readily derived as:  $\mu = \alpha / (\alpha + X)$



(Figure 5.2). At  $-55^{\circ}\text{C}$  the  $V_{GS}$  of a DTMOST diode is about 400mV, about half that of a diode-connected BJT.

In the front-end of the proposed sensor (Figure 5.1), the voltages  $V_{GS}$  and  $\Delta V_{GS}$  are generated by two identical DTMOST diodes ( $W/L = 80\mu\text{m} / 0.65\mu\text{m}$ ) biased at  $I_b = 90\text{nA}$  (at  $25^{\circ}\text{C}$ ) and  $5 \cdot I_b$ . The required current ratio  $p = 5$  is established with the help of six unit current sources. As in [5.1], [5.3], dynamic element matching (DEM) is used to average out the mismatch between these current sources and between the two DTMOST diodes, as this would otherwise cause significant temperature-sensing errors. The biasing currents are generated by the same PNP-based bias circuit used in the BJT-based sensor of [5.3]. This choice allows a fair comparison to be made between BJT- and DTMOST-based temperature sensors, but also precludes sub-1V operation.

Another source of current-ratio error, especially at high temperatures, is the source-to-bulk leakage current of the DTMOST diodes. This was minimized by appropriate device sizing using standard device models [5.9]. The area of the optimized devices is then in the same order as that of the PNPs in [5.3]. It should be noted that this non-ideal leakage current is two orders of magnitude less than, the similarly non-ideal, base current of a parasitic PNP transistor in the same process. As a result, DTMOST-based sensors, unlike their PNP-based counterparts, do not require extra  $\beta$ -compensation-like circuitry [5.1], [5.3].

The ratio  $X = V_{GS}/\Delta V_{GS}$  is accurately digitized by a 16-bit zoom-ADC, which combines the advantages of a SAR-ADC and a 1<sup>st</sup>-order  $\Delta\Sigma$  converter in a two-step conversion scheme [5.3]. As described in section 3.3, a SAR algorithm is first employed to find the integer part of  $X$ , denoted by  $n$ , by comparing  $V_{GS}$  to integer multiples of  $\Delta V_{GS}$ . The fractional part  $\mu'$  is then determined by a 1<sup>st</sup>-order  $\Delta\Sigma$ -ADC, whose references are chosen so as to zoom into the region determined by the SAR algorithm, i.e. from  $n \cdot \Delta V_{GS}$  to  $(n + 1) \cdot \Delta V_{GS}$ . A modified 1<sup>st</sup>-order SC  $\Delta\Sigma$ -ADC is used to implement both conversion steps, resulting in a compact, energy-efficient ADC.

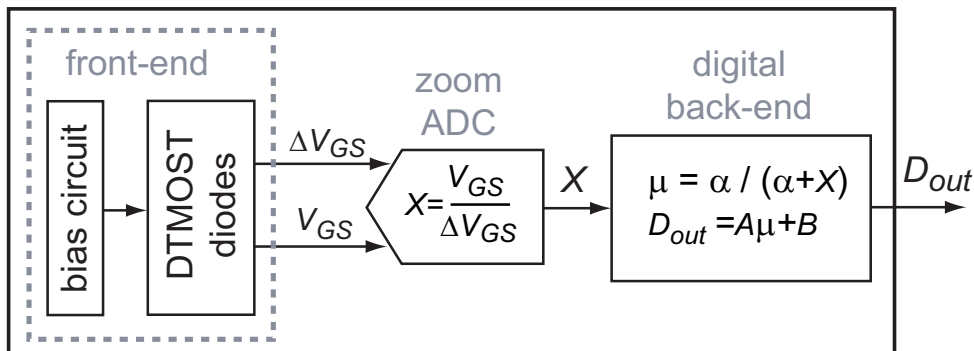


Figure 5.2: Block diagram of the smart temperature sensor.

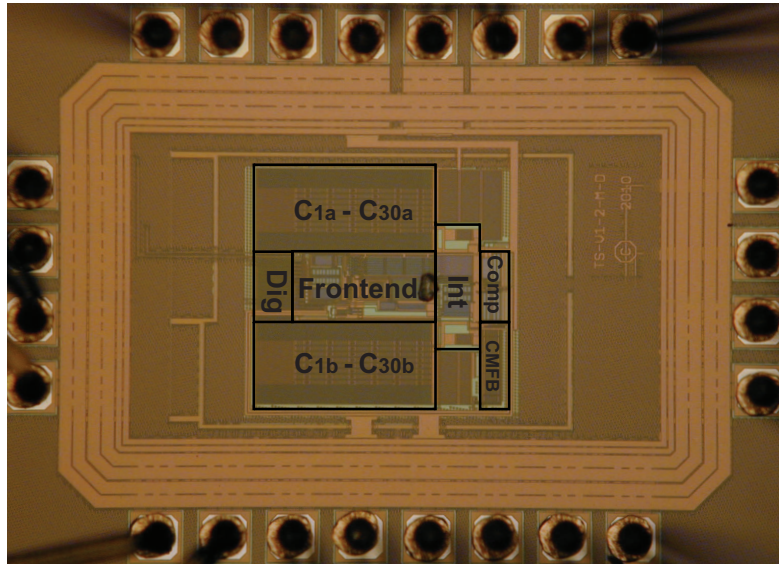


Figure 5.3: Chip micrograph of the first prototype DTMOST-based sensor employing a 1<sup>st</sup>-order zoom-ADC.

### 5.1.3 Measurement Results

The temperature sensor was realized in a standard 0.16 $\mu\text{m}$  CMOS process with five metal layers (Figure 5.3). The chip has an active area of 0.12mm<sup>2</sup>, and consumes 8.6 $\mu\text{W}$  from a 1.8V supply at 30°C. The sensor has a supply sensitivity of 0.1°C/V over a supply range of 1.6V-2V. For flexibility, the digital back-end, the control logic and the  $\Delta\Sigma$  modulator's sinc<sup>2</sup> decimation filter, were implemented off-chip.

Twenty devices from one batch were packaged in ceramic DIL packages, placed in a climate chamber and then characterized over the military temperature range: -55°C to 125°C. As shown in Figure 5.4, the sensor's batch-calibrated inaccuracy was about  $\pm 1.5^\circ\text{C}$  ( $3\sigma$ ), with a PTAT-like characteristic and a 3<sup>rd</sup>-order systematic non-linearity of about  $\pm 0.1^\circ\text{C}$ . After digital non-linearity compensation and a single offset-trim at 30°C, the sensor's inaccuracy was reduced to  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ), as shown in Figure 5.5. An  $\alpha$ -trim [5.3] resulted in slightly worse inaccuracy:  $\pm 0.5^\circ\text{C}$  ( $3\sigma$ ).

As shown in Figure 5.6, a two-temperature trim (commonly used by MOSFET-based temperature sensors [5.6]) at -10°C and 90°C, reduces the sensor's inaccuracy to  $\pm 0.1^\circ\text{C}$ . At 5 conversions/s (1024  $\Delta\Sigma$ -cycles), the sensor achieves a resolution of 33mK (rms), which is lower than that of BJT-based sensor in [5.3], which uses a similar readout topology. As previously discussed, for a similar 5 : 1 ratio of bias currents used in the front-end, the ratio  $X$  in a DTMOST-based sensor is about half that of a BJT-based sensor. The

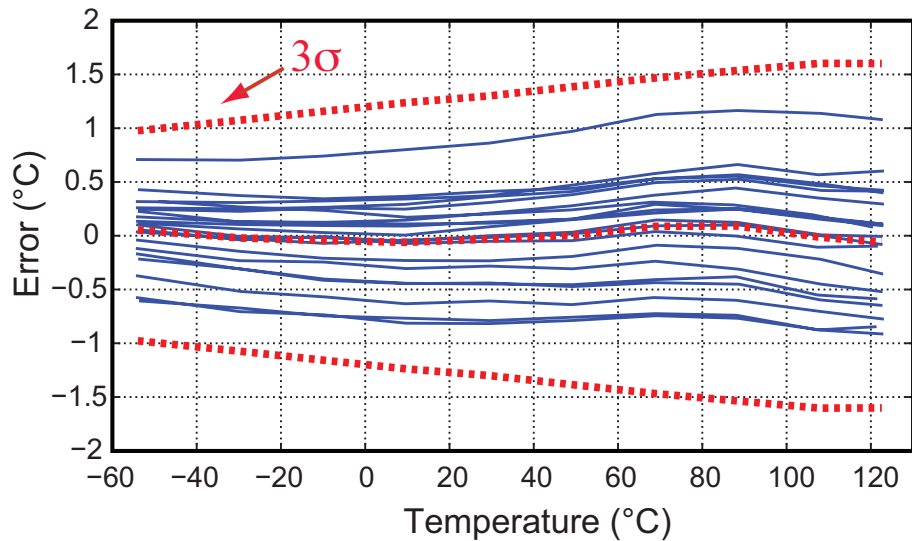


Figure 5.4: Temperature error of 20 measured DTMOST-based sensors before trimming (linear fit).

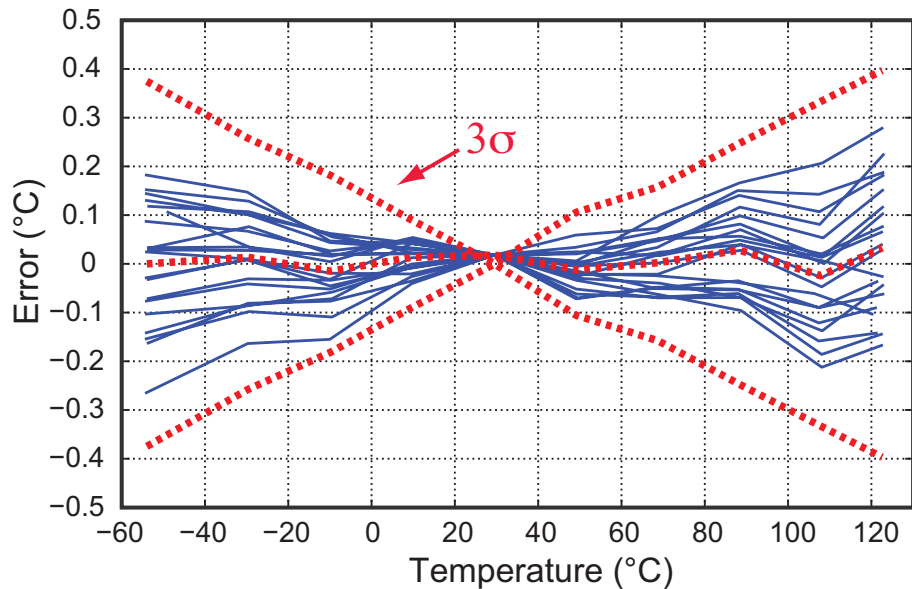


Figure 5.5: Temperature error of 20 measured DTMOST-based sensors after a single offset trim at 30°C.

number of  $\Delta V_{GS}$  sampling capacitors are, therefore, proportionally smaller and results in a larger  $kT/C$  noise. The sensor's performance is summarized in Table 5.1. Compared to the prior-art MOSFET-based temperature sensors [5.5], [5.6], this represents 5x more accuracy over a much wider temperature range and proves DTMOSTs as a promising sensing element.

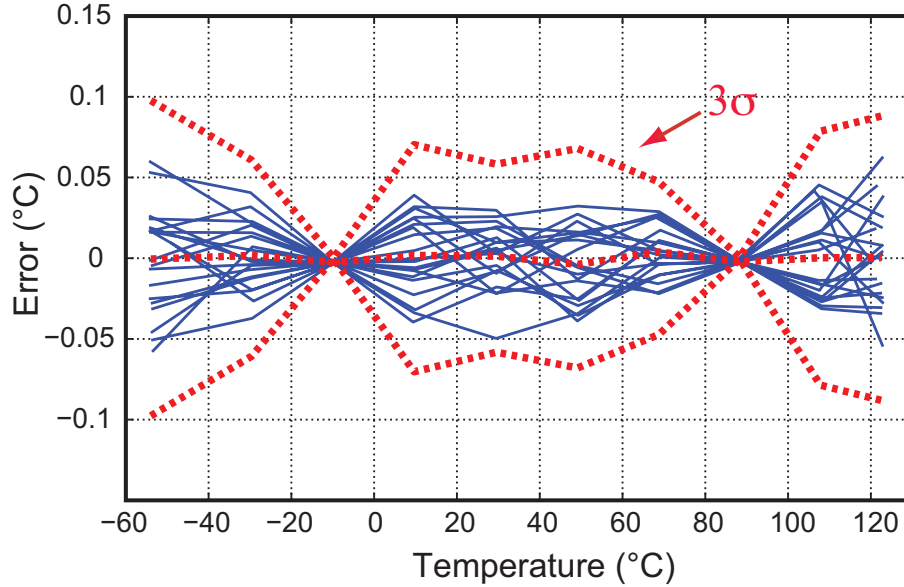


Figure 5.6: Temperature error of 20 measured DTMOST-based sensors after two-temperature trim at  $-10^{\circ}\text{C}$  and  $90^{\circ}\text{C}$ .

Parameter	This work ESSCIRC'11 [5.13]	S&A'11 [5.5]	JSSC'10 [5.6]
CMOS technology	0.16 $\mu\text{m}$	0.35 $\mu\text{m}$	0.35 $\mu\text{m}$
Chip area	0.12 $\text{mm}^2$	0.08 $\text{mm}^2$	0.6 $\text{mm}^2$
Supply current (RT) <sup>†</sup>	4.7 $\mu\text{A}$	4.5 $\mu\text{A}$	10.6mA
Supply sensitivity	0.1 $^{\circ}\text{C}/\text{V}$	2.5 $^{\circ}\text{C}/\text{V}$ (at 30 $^{\circ}\text{C}$ )	10 $^{\circ}\text{C}/\text{V}$ (at 30 $^{\circ}\text{C}$ )
Inaccuracy, (calibration points)	$\pm 0.1^{\circ}\text{C}$ ( $3\sigma$ ), (2) $\pm 0.4^{\circ}\text{C}$ ( $3\sigma$ ), (1)	-1.8 $^{\circ}\text{C}$ / +1 $^{\circ}\text{C}$ (max), (1)	-0.4 $^{\circ}\text{C}$ / +0.6 $^{\circ}\text{C}$ (max), (2)
Temperature range	-55 $^{\circ}\text{C}$ - 125 $^{\circ}\text{C}$	10 $^{\circ}\text{C}$ - 80 $^{\circ}\text{C}$	0 $^{\circ}\text{C}$ - 90 $^{\circ}\text{C}$
Resolution, ( $T_{\text{conv}}$ )	0.033 $^{\circ}\text{C}$ , (200msec)	0.1 $^{\circ}\text{C}$ , (10msec)	0.0918 $^{\circ}\text{C}$ , (90 $\mu\text{sec}$ )

<sup>†</sup> Excluding the off-chip digital

Table 5.1: Performance summary and comparison with previous work.

For comparison, a second chip was realized with essentially the same readout circuitry, but in which the DTMOST diodes were replaced by PNPs. For  $p = 5$ , however, the ratio  $X = V_{GS}/\Delta V_{GS}$  then spans a larger range: varying from 6 to 28 from  $-55^{\circ}\text{C}$  to  $125^{\circ}\text{C}$  [5.3]. To cover this range, the range of the zoom-ADC's SAR phase was extended, by proportionally increasing the number of elements in its capacitor DAC. As shown in Figure 5.7, the inaccuracy of the PNP-based sensor is  $\pm 0.2^{\circ}\text{C}$  ( $3\sigma$ ) after a single  $\alpha$ -trim:

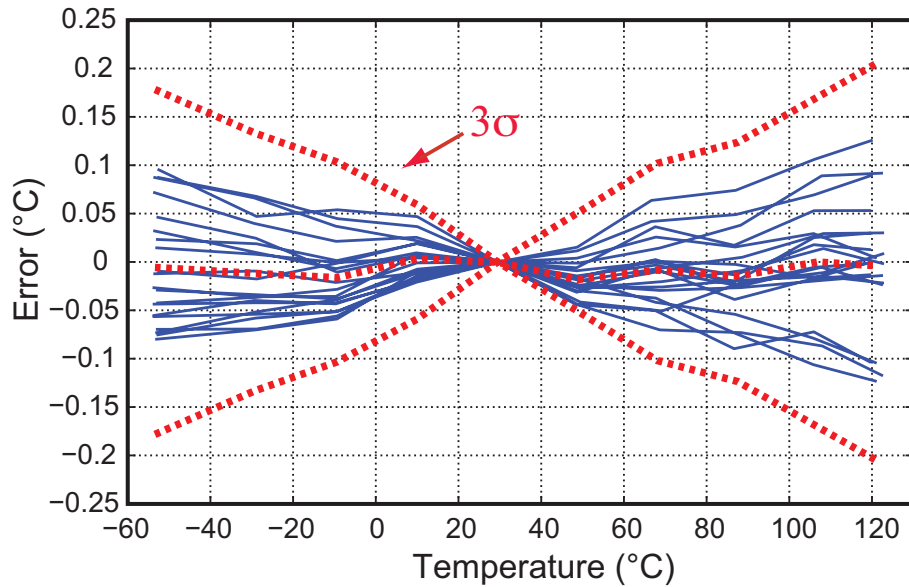


Figure 5.7: Temperature error of 20 measured PNP-based sensors after a single  $\alpha$ -trim at  $30^\circ\text{C}$ .

only 2x less than that of, a similarly trimmed, DTMOST-based sensor.

## 5.2 A Sub-1V All-CMOS Temperature Sensor <sup>2</sup>

The first DTMOST-based prototype presented in the previous section is not capable of operating at sub-1V supply voltages, due to the use of a BJT-based bias circuit. In this section a second prototype will be presented. It is an ultra-low-power all-CMOS temperature sensor, which operates from a 0.85V supply and draws only 700nA. To date, this is the only all-CMOS temperature sensor to achieve both sub-1V operation and a high accuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) over a wide temperature range ( $-40^\circ\text{C}$  to  $125^\circ\text{C}$ ), while employing only a single room-temperature trim.

### 5.2.1 Sensor Front-end

Consider the bias circuit shown in Figure 5.1. As already discussed in chapter 3, the circuit generates a PTAT,  $\beta_F$ -dependent bias current. It is then mirrored into the sensor

<sup>2</sup>K. Souri, Y. Chae, F. Thus and K.A.A. Makinwa, “A 0.85V 600nW All-CMOS Temperature Sensor with an Inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) from  $-40^\circ\text{C}$  to  $125^\circ\text{C}$ ,” in *Dig. Techn. Papers ISSCC*, pp. 222 - 223, Feb. 2014.

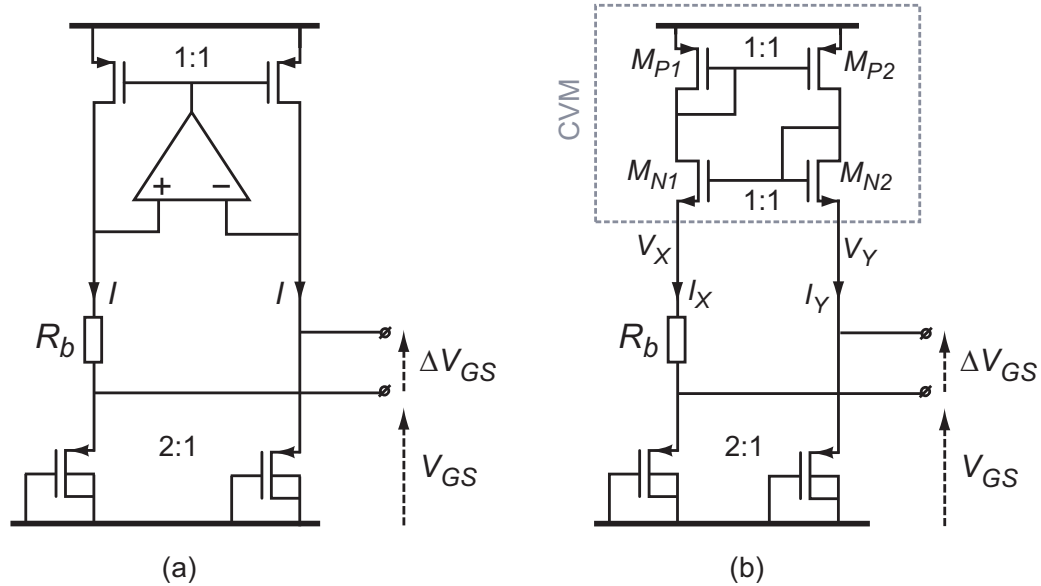


Figure 5.8: Circuit diagram of the DTMOST-based front-end using an amplifier (a), and a current-voltage mirror (b).

front-end to bias two identical DTMOSTs at a 5 : 1 current ratio. Unlike BJTs, however, DTMOSTs do not suffer from an equivalent to base current, thus ruling out the need for such  $\beta_F$ -compensation circuitry. In this design, therefore, the simple sensor front-end shown in Figure 5.8.a, has been adopted. A pair of DTMOST diodes with a 1 : 2 area ratio are biased with identical bias currents  $I = \Delta V_{GS}/R_b$ , which are enforced by a self-biased opamp configuration.

Since  $V_{GS} \approx 0.4V$  at  $-55^\circ C$ , and assuming some voltage headroom for the cascoded current sources (not shown), using a supply voltage below 1V is readily feasible. The challenge then would be to design a low-power opamp which operates from a sub-1V supply. A well-known, energy-efficient solution is the current-voltage mirror (CVM) circuit of Figure 5.8.b [5.11]. The top PMOS current mirrors  $M_{P1}$  and  $M_{P2}$  force  $I_X = I_Y$ , while  $M_{N1}$  and  $M_{N2}$  force  $V_X = V_Y$ , given their equal sizing and drain current values. Two diode-connected transistors  $M_{N1}$  and  $M_{P2}$  ensure an overall negative feedback around the loop. Although extremely compact and energy-efficient, this circuit exhibits poor regulation. Since the drain-source voltage of the PMOS current mirrors are different, channel length modulation makes  $I_X \neq I_Y$ , thereby resulting in a systematic offset between the  $V_{GS}$  voltages of  $M_{N1}$  and  $M_{N2}$ . This offset is subject to the supply and temperature changes and will result in an untrimmable error at the output of the sensor.

To solve the aforementioned issues, the circuit of Figure 5.9.a has been implemented in this work. The so-called symmetrically-matched current-voltage mirror (SM CVM) [5.12] forces the PMOS current mirrors to essentially have the same gate, drain, and

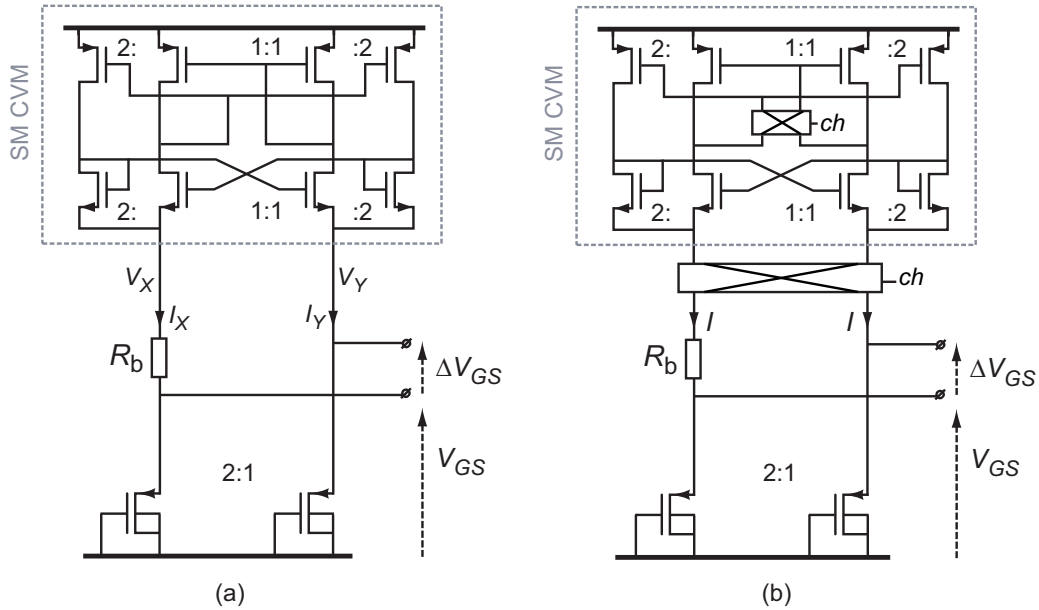


Figure 5.9: Circuit diagram of the DTMOST-based front-end using a SM CVM (a). Applying chopping to improve the accuracy of SM CVM (b).

source voltages, thereby forcing equal currents  $I_X = I_Y$ , and in turn, maintain  $V_X = V_Y$ . The voltage headroom for a CVM to operate is  $V_{GS,PMOS} + V_{DS,NMOS}$ . Given the low bias currents in each branch (90nA in this design), the top current mirrors are designed to operate in weak inversion, thus reducing their  $V_{GS}$  voltage. Assuming  $V_{GS,PMOS} = 0.3V$  and  $V_{DS,NMOS} = 100 - 150mV$ , a voltage headroom of  $\approx 450mV$  is sufficient, thus enabling sub-1V operation.

### 5.2.2 Accuracy Issues

A main source of error is the CVM's input-referred offset and  $1/f$  noise, which directly add to  $\Delta V_{GS}$  and thus impact the accuracy of the bias currents, and hence of both  $V_{GS}$  and  $\Delta V_{GS}$ . Such errors are mitigated by incorporating choppers into the CVM circuitry as shown in Figure 5.9.b. While the input chopper swaps the two inputs of the amplifier, thereby modulating the offset and  $1/f$  noise to chopping frequency  $f_{chop}$ , a second chopper inside the CVM sustains the negative feedback polarity. The succeeding  $\Delta\Sigma$  modulator and low-pass digital decimation filter then filter out the high-frequency components present at  $f_{chop}$ .

To minimize the effect of mismatch between the DTMOSTs, which would otherwise impact the accuracy of  $\Delta V_{GS}$ , the 1 : 2 area ratio is established by incorporating three unit DTMOSTs ( $W/L = 90\mu m / 0.7\mu m$  for each) in a dynamic element matching (DEM)

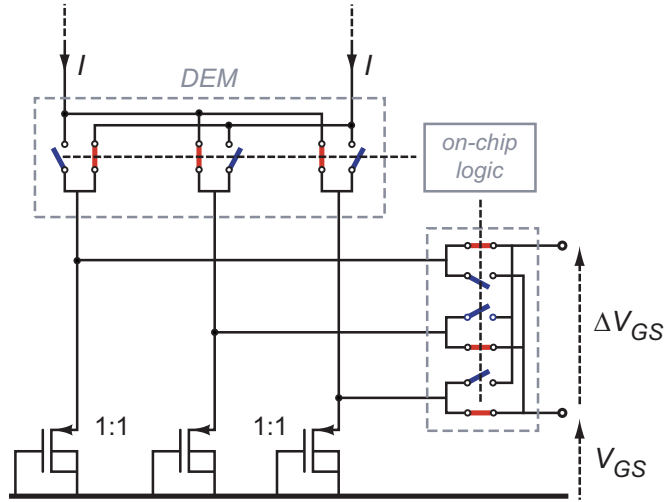


Figure 5.10: DEM and Kelvin connection used to improve the accuracy of the front-end.

scheme, as shown in Figure 5.10. Since the associated DEM switches carry bias current, Kelvin connections are used to accurately read out  $V_{GS}$  and  $\Delta V_{GS}$ . Unlike in [5.3] where off-chip DEM control was used, here, an on-chip DEM logic based on a 3-bit ring counter is employed.

### 5.2.3 System Diagram

The sensor's block diagram is shown in Figure 5.11.a. It consists of the DTMOST front-end of Figure 5.9.b, a  $2^{nd}$ -order inverter-based zoom-ADC, a voltage doubler and some control logic. As in [5.3], the zoom-ADC uses a power-efficient coarse/fine algorithm to convert the front-end's output voltages  $V_{GS}$  and  $\Delta V_{GS}$  into a temperature-dependent ratio  $X = V_{GS}/\Delta V_{GS}$ . For the choice of DTMOSTs with a 1 : 2 area ratio,  $X$  varies from 5 to 28 over the temperature range from  $-40^{\circ}\text{C}$  to  $125^{\circ}\text{C}$ . As before, an off-chip digital back-end then computes a PTAT ratio  $\mu = \alpha/(\alpha + X)$ , where  $\alpha$  is a gain factor, which could be also used as a trimming knob to compensate for  $V_{GS}$  spread.

Figure 5.11.b, shows the timing diagram of a full temperature-to-digital conversion. As shown, each conversion starts with a 5-bit SAR, which determines the integer part of  $X$ , i.e.  $n$ . It is then followed with a fine  $\Delta\Sigma$  conversion of length  $N_{\Delta\Sigma}$  to determine the fraction  $\mu'$ . Due to mismatch, the  $X$  value may change significantly during the two states of the chopped front-end, thus resulting in different coarse values  $n$ . This issue is solved as follows: first the chopping state is set to "0" and a coarse/fine conversion is performed, resulting in  $X_1 = n_1 + \mu'_1$ . The chopping state is then swapped and a second coarse/fine conversion is done to obtain  $X_2 = n_2 + \mu'_2$ . The final result is achieved by averaging the



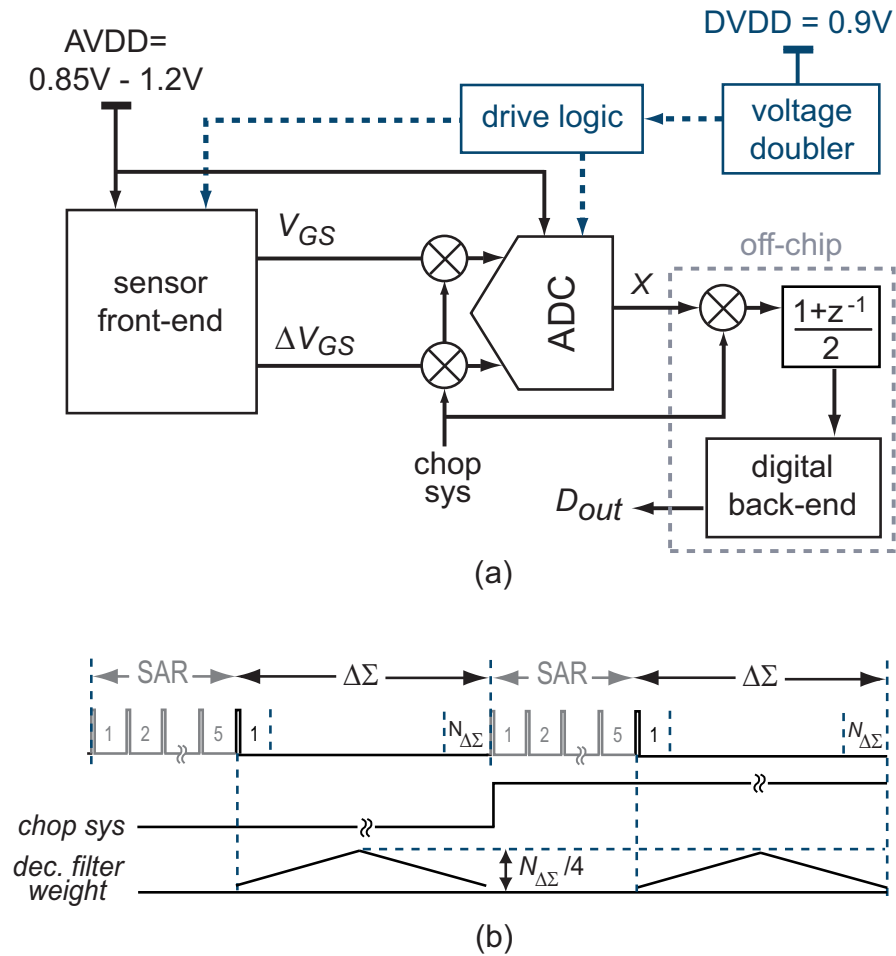


Figure 5.11: Top: The sensor’s block diagram. Bottom: Timing diagram of a temperature conversion.

results over two conversions, i.e.  $X = (X_1 + X_2)/2$ . To minimize the ADC’s own residual offset, system-level chopping with similar timing is employed, as shown in Figure 5.11.b.

## 5.2.4 Power Domains

The sensor has two supply voltages: an analog supply AVDD, which powers the front-end and the ADC, and a digital supply DVDD, which powers the voltage doubler. AVDD can range between 0.85V - 1.2V, while DVDD = 0.9V. The voltage doubler provides the drive logic block with  $\approx 1.8V$  that, in turn, drives the gates of NMOS switches that sample  $V_{GS}$  and  $\Delta V_{GS}$ , thus facilitating the use of sub-1V supply voltages. The voltage doubler employs a charge-pump topology as in Figure 5.12. Running at 80kHz, the two anti-phase clocks  $\phi_1$ ,  $\phi_2$  with the amplitude of  $V_{clk}$ , along with the input voltage  $V_{in}$  are used

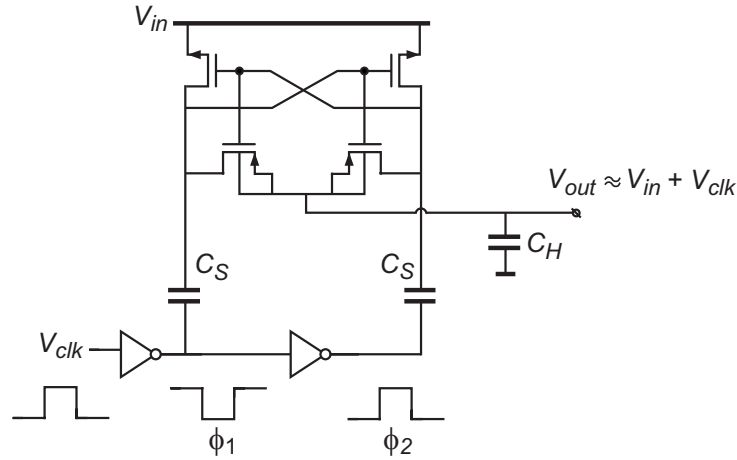


Figure 5.12: Circuit diagram of the switched-capacitor voltage doubler.

to periodically charge the sampling capacitors  $C_S$ , and thus delivering an output voltage  $V_{out} \approx V_{clk} + V_{in}$ . The holding capacitor  $C_H = 1\text{pF}$  is used to reduce the output ripples to less than 30mV. Since the power dissipated in the drive logic is small, the area/power overhead due to the use of the charge-pump block is minimal.

### 5.2.5 Inverter-Based Zoom ADC

The heart of the zoom-ADC is a feed-forward  $2^{nd}$ -order switched-capacitor  $\Delta\Sigma$ -ADC, as shown in Figure 5.13. At its input, a cap-DAC with 30 unit elements (each 60fF) samples either  $V_{GS}$  or  $k \cdot \Delta V_{GS}$ , where  $k = [1 \dots 30]$ . The two capacitors  $C_G = 0.5 \cdot C_{Sk}$  are used to implement guard-banding.

In contrast to [5.3], both integrators are built around pseudo-differential inverter-based amplifiers, thus fully exploiting the reduced integrator swing conferred by zooming. The first integrator draws 135nA while the second, less critical, integrator draws only 66nA. Figure 5.14 shows the implementation details of the  $1^{st}$  integrator's OTA. The  $2^{nd}$  integrator essentially uses the same topology, except that its current is scaled by half. To decrease the inverter's sensitivity to power supply and process spread, a dynamic current-biasing technique is used. During the auto-zeroing phase  $\phi_1$ ,  $M_{N1}$  and  $M_{P1}$  are diode-connected and biased with two current sources. The left plates of the offset storing capacitors  $C_{OS}$  are connected to the signal ground in this phase, thus storing the operating bias voltages of  $M_{N1}$  and  $M_{P1}$ . The bias voltages  $V_{b1}$  and  $V_{b2}$  are chosen such that  $M_{N2}$ ,  $M_{P2}$  are essentially off during  $\phi_1$ . The two bias currents are mirrored from the front-end's precision bias circuit to ensure robustness to supply and process variations. In  $\phi_2$ , the two biasing current sources are disconnected and  $M_{N1}$  and  $M_{P1}$  are configured as common-

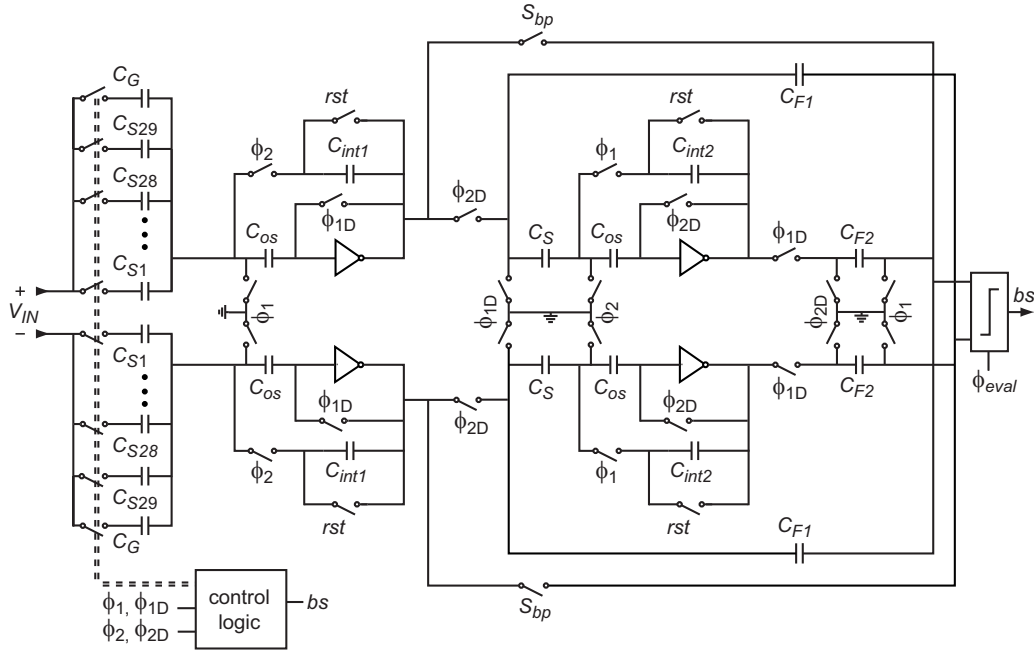


Figure 5.13: Circuit diagram of the 2<sup>nd</sup>-order inverter-based zoom-ADC.

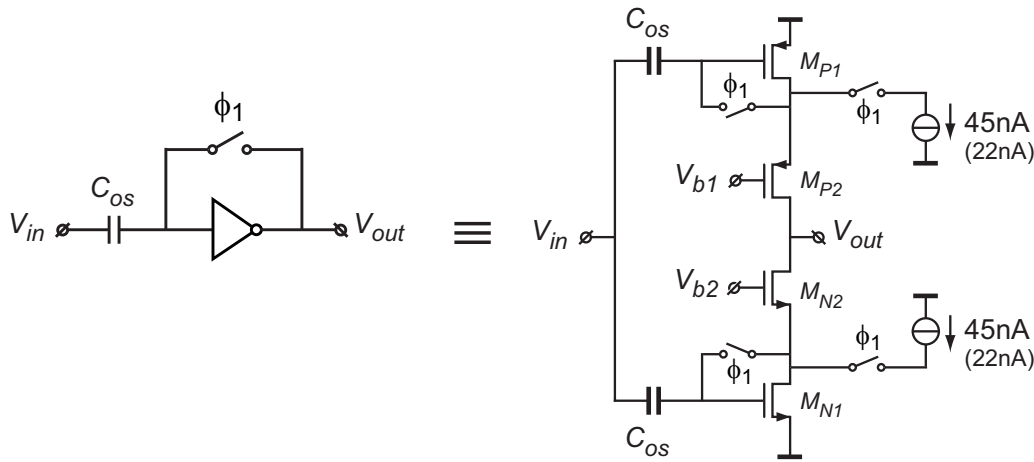


Figure 5.14: Circuit diagram of the inverter-based OTA. A PTAT bias current of 45nA and 22nA (at room temperature) is used in the first and second OTAs, respectively.

source amplifiers and together form a class-AB amplifier, with a virtual ground at  $V_{in}$ . Since the output voltage swing requirement is reduced to about  $\pm 200\text{mV}$  by the prior coarse conversion step, cascoding of  $M_{N1}$  and  $M_{P1}$  is readily possible, thus enhancing the inverter's output resistance, and hence its DC gain. A passive summation network at the input of quantizer combines the output of 2<sup>nd</sup> integrator with that of the 1<sup>st</sup> integrator via the feed-forward capacitor  $C_{F1}$ , as shown in Figure 5.13. During the coarse conversion, the first integrator computes  $V_{GS} - k \cdot \Delta V_{GS}$ , while its output is connected directly to the

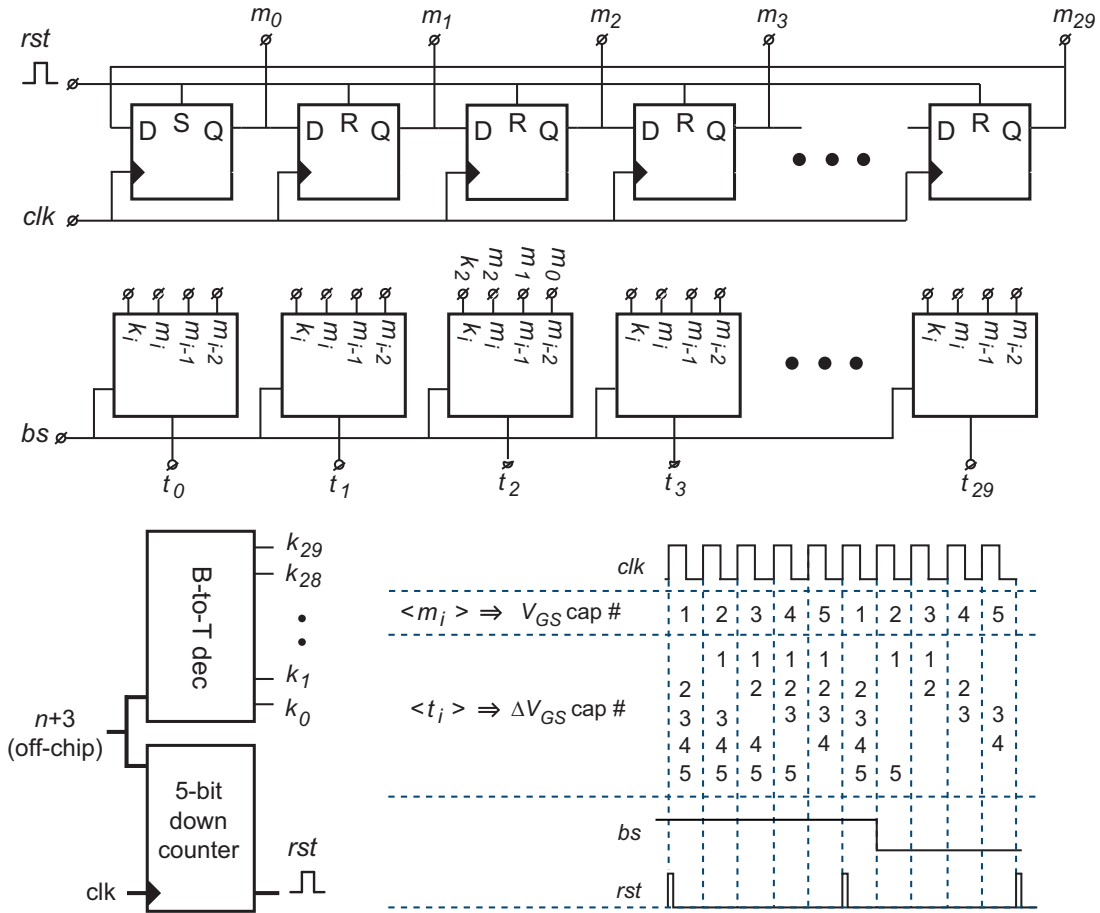


Figure 5.15: Top: On-chip DEM logic for the zoom-ADC’s Cap-DAC. Bottom: Timing diagram for a 5-element Cap-DAC as an example.

comparator via the switch  $S_{bp}$ . Off-chip logic then implements the SAR algorithm by applying trial values  $k$  to the chip and monitoring the comparator’s output.

During the fine conversion, the mismatch between the unit elements of the cap-DAC is mitigated by the use of DEM. In contrast to [5.3], the required DEM logic is implemented on-chip. Figure 5.15 shows the circuit diagram of the implemented on-chip DEM logic. It consists of a 30-bit circular shift-register (SR), with an effective length of  $n + 3$  bits, where  $n$  is the result of the coarse conversion. The SR length is defined by a periodic reset signal, which is generated by means of an on-chip 5-bit down-counter, that is preloaded with  $n + 3$  at every falling edge of the resulting reset signal. Resetting the SR loads it with a single logic “1” appearing at  $m_0$  output, which then circulates on every succeeding clock pulse. This bit (via the  $m_i$  outputs) is then used to select the single capacitor that samples  $V_{GS}$ .

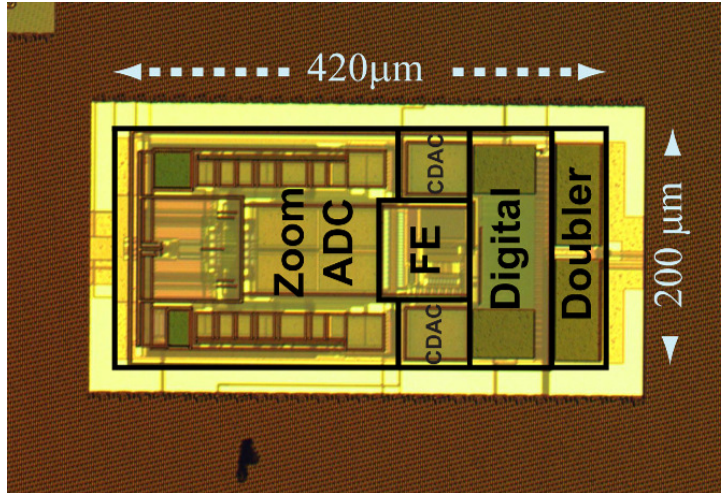


Figure 5.16: Chip micrograph of the second prototype DTMOST-based sensor employing a  $2^{nd}$ -order zoom-ADC.

During the fine conversion, either  $n$  or  $n+2$  capacitors will be used to sample  $\Delta V_{GS}$ , depending on the bitstream output  $bs$ . An on-chip decoder is used to convert the binary code  $n+3$  into a thermometer code. The resulting thermometer code  $[k_0 k_1 \dots k_{29}]$ , the  $V_{GS}$  sampling control  $[m_0 m_1 \dots m_{29}]$ , and the bitstream output  $bs$  are then applied to a simple combinational logic to generate the  $\Delta V_{GS}$  sampling control  $[t_0 t_1 \dots t_{29}]$ . During the coarse conversion, the SR is reset, so that the same capacitors are always used for the SAR conversion.

### 5.2.6 Prototype and Measurement Results

The second prototype was also realized in a standard  $0.16\mu\text{m}$  CMOS process. Figure 5.16 shows the chip micrograph of the sensor. It occupies  $0.085\text{mm}^2$ , and draws  $700\text{nA}$  from a  $0.85\text{V}$  supply. The front-end and ADC draw  $560\text{nA}$ , while the voltage doubler and the rest of the on-chip digital circuitry draw  $140\text{nA}$ . For flexibility, the SAR logic and the  $\text{sinc}^2$  decimation filter were implemented off-chip. However, simulations show that implementing them on-chip would only incur an extra  $10\text{nW}$  per conversion. With DVDD fixed at  $0.9\text{V}$ , AVDD was varied from  $0.85\text{V}$  to  $1.2\text{V}$ . The corresponding supply sensitivity of the front-end and ADC was  $0.45^\circ\text{C}/\text{V}$ .

A total of 16 devices in ceramic DIL packages were characterized over the temperature range from  $-40^\circ\text{C}$  to  $125^\circ\text{C}$ . As shown in Figure 5.17 (top), their batch-calibrated inaccu-

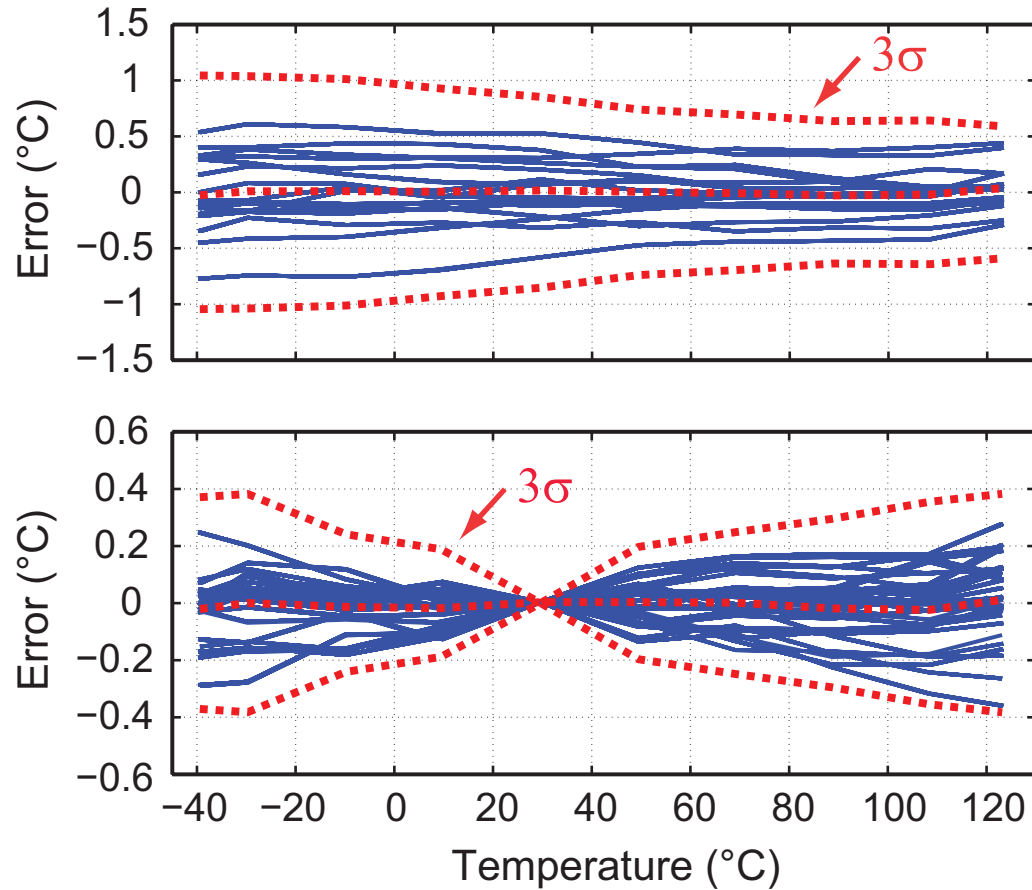


Figure 5.17: Measured temperature error of 16 samples before trimming (top) and after a single  $\alpha$ -trim at 30°C (bottom).

racy was  $\pm 1^\circ\text{C}$  ( $3\sigma$ , 16 devices), with a residual curvature of only  $0.03^\circ\text{C}$ . Compared to the first prototype (Figure 5.4), this design exhibits better untrimmed accuracy, but with a CTAT rather than PTAT profile. At cold, the bias current in the sensor front-end is minimum, while the number of sampling capacitors is maximum. Therefore, any settling errors will be larger at cold, which could explain the CTAT profile of the untrimmed inaccuracy shown in Figure 5.17 (top). Furthermore, the residual offset of the SM CVM will appear in series with  $\Delta V_{GS}$ , and thus directly impacts the sensor's accuracy. This effect is more severe at cold, when  $\Delta V_{GS}$  is smaller. After an  $\alpha$ -trim at 30°C, the inaccuracy improves to  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ), as shown in Figure 5.17 (bottom), which is similar to the first prototype. Offset trimming is slightly worse, resulting in an inaccuracy of  $\pm 0.5^\circ\text{C}$  ( $3\sigma$ ). These results show that DTMOSTs, like BJTs, can be effectively trimmed at a single temperature.

Parameter	2 <sup>nd</sup> Prototype ISSCC'14 [5.18]	1 <sup>st</sup> Prototype ESSCIRC'11 [5.13]	JSSC'10 [5.15]	CICC'08 [5.16]	VLSI'11 [5.17]
Sensor type	DTMOST	DTMOST	MOSFET	MOSFET	Resistor
CMOS technology	0.16 $\mu$ m	0.16 $\mu$ m	0.18 $\mu$ m	0.18 $\mu$ m	0.18 $\mu$ m
Chip area	0.085mm <sup>2</sup>	0.12mm <sup>2</sup>	0.042mm <sup>2</sup>	0.05mm <sup>2</sup>	0.18mm <sup>2</sup>
Supply current (RT) <sup>†</sup>	700nA	4.6 $\mu$ A	190nA	220nA	20 $\mu$ A
Supply voltage	0.85V - 1.2V	1.5V - 2V	0.5V (sensor) 1.0V (digital)	1V	1.2V - 2V
Supply sensitivity	0.45°C/V	0.1°C/V	-	Supply referenced	0.625°C/V
Inaccuracy (calibration points)	$\pm 0.15^\circ\text{C}$ (3 $\sigma$ ) (1)	$\pm 0.2^\circ\text{C}$ (3 $\sigma$ ) (1)	-0.8°C / +1°C (2)	-1.6°C / +3°C (2)	$\pm 0.5^\circ\text{C}$ (max) (1)
Temperature range	-55°C - 125°C	-30°C - 125°C	-10°C - 30°C	0°C - 100°C	0°C - 100°C
Relative inaccuracy <sup>††</sup>	0.48%	0.44%	4.5%	4.6%	1%
Resolution (T <sub>conv</sub> )	0.063°C (6msec)	0.033°C (200msec)	0.2°C (30msec)	0.1°C (100msec)	0.25°C (12.5 $\mu$ sec)
Resolution FoM	14.1pJ°C <sup>2</sup>	1535pJ°C <sup>2</sup>	140pJ°C <sup>2</sup>	220pJ°C <sup>2</sup>	19pJ°C <sup>2</sup>

<sup>†</sup> Excluding the off-chip digital

<sup>††</sup> Relative inaccuracy (%) = 100 · (maximum error / specified temperature range)

Table 5.2: Performance summary and comparison with previous work.

While running at a clock frequency of 25kHz, the sensor requires only 3.6nJ to achieve a  $kT/C$  limited resolution of 63mK (rms) in a conversion time of 6ms. This corresponds to a resolution FoM [5.14] of 14.1pJK<sup>2</sup>, which is in line with the state-of-the-art BJT-based temperature sensor presented in previous chapter [5.3]. The performance of the two DTMOST-based sensor prototypes presented in this chapter are summarized in Table 5.2 and compared to that of other state-of-the-art low-voltage designs. It can be seen that the DTMOST-based designs improve the relative accuracy by at least 2x, while the second prototype achieves the best energy-efficiency, evidenced by a resolution FoM of 14.1pJ°C<sup>2</sup>. Compared to the second BJT-based prototype presented in Chapter 3, this work achieves similar energy-efficiency, but at supply voltages as low as 0.85V. The price to pay will be  $\sim 2x$  less accuracy, which is acceptable in many applications.

## 5.3 Conclusions

BJT-based temperature sensors, although capable of achieving good energy-efficiency and accuracy over wide temperature ranges, are not capable of operating at low supply voltages, which is an important requirement in battery-powered applications and nano-scale CMOS processes. A fully CMOS compatible device, a DTMOST diode, was shown to be a promising alternative to parasitic BJTs as a temperature sensing element. When operated in weak inversion, a DTMOST enables sub-1V operation, while achieving a reasonably high accuracy over a wide temperature range. In this chapter two sensor prototypes based on such sensing elements were demonstrated. The first prototype enabled an apples-to-apples comparison with BJTs, and resulted in the conclusion that DTMOSTs are indeed only a factor 2x less accurate in the chosen 160nm process. Employing a fully inverter-based 2<sup>nd</sup>-order zoom-ADC, the second prototype achieved a FoM of 14.1pJ°K<sup>2</sup>, which was in line with state-of-the-art BJT-based sensors, back in 2014. When operating at a supply voltage of 0.85V and after a single-temperature trim, the sensor achieved an inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) from  $-40^\circ\text{C}$  to  $125^\circ\text{C}$ . These results prove that DTMOSTs could be considered as the temperature sensors of choice when sub-1V, high accuracy, and energy-efficiency are key requirements.



## 5.4 References

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## Chapter 6

# Conclusions

In this thesis, the development of energy-efficient, accurate smart temperature sensors for wireless temperature sensing applications has been investigated. It has been shown that the existing temperature sensors prior to the start of this research were ill-suited for use in such applications, where energy-efficiency and low-cost are critical requirements. In the following, first the main findings of this research are discussed. The other applications of the developed techniques are then presented, followed by some proposals for the future improvements.

### 6.1 Main Findings

The main findings of this thesis are as follows:

- After a single temperature trim, a BJT-based temperature sensor achieved an inaccuracy of  $\pm 0.2^\circ\text{C}$  ( $3\sigma$ ) from  $-55^\circ\text{C}$  to  $125^\circ\text{C}$  across different lots of a  $0.16\mu\text{m}$  CMOS process. In older CMOS processes and using a similar sensing approach, slightly better accuracy was achieved, e.g.  $\pm 0.1^\circ\text{C}$  ( $3\sigma$ ) in a  $0.7\mu\text{m}$  CMOS process. Down to the  $0.16\mu\text{m}$  node, therefore, CMOS scaling does not significantly impact the accuracy of BJT-based temperature sensors (Chapter 4).
- Of the various CMOS-compatible temperature sensing elements, substrate PNPs are best suited for use in wireless temperature sensing applications. This is because they can be operated at bias currents down to tens of nA with minimal impact on their accuracy. Moreover, they exhibit a well-defined process spread, which can be effectively trimmed at a single-temperature (Chapter 1).

- A two-step zoom-ADC was developed, which combines the speed of a coarse SAR-ADC, with the high resolution/accuracy of a fine  $\Delta\Sigma$ -ADC. In contrast to conventional  $\Delta\Sigma$ -ADCs, the full-scale range of the fine converter is considerably reduced, which notably relaxes various key requirements such as the number of  $\Delta\Sigma$ -cycles and the DC-gain and swing of the loop filter. In this architecture, both conversion time and power-efficiency can be improved, resulting in a substantial improvement in the energy-efficiency. The fact that dynamic correction techniques can still be used in the fine conversion phase, ensures that the accuracy of the zoom-ADC can be as good as that of conventional  $\Delta\Sigma$ -ADC architectures (Chapter 3).
- BJT-based temperature sensors are also applicable to automotive applications, where high temperature operation ( $>150^\circ\text{C}$ ) is desired. It was shown that by optimizing the emitter area and bias current of substrate PNPs, the impact of their saturation current  $I_S$  at high temperatures can be well mitigated. Furthermore, robust circuit techniques were employed to cope with the various leakage currents at such temperatures. The sensor achieved an accuracy similar to that of state-of-the-art thermal diffusivity (TD) sensors, while drawing more than an order of magnitude less current (Chapter 4).
- DTMOSTs were shown to be a promising alternative to parasitic BJTs for temperature sensing at low supply voltages. When operated in weak inversion, a DTMOST enables sub-1V operation, while achieving a reasonably high accuracy over a wide temperature range. In a  $0.16\mu\text{m}$  process, after a single-point trim, an inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) was achieved over the military temperature range, which is only a factor 2x worse compared to, similarly trimmed, BJT-based sensors (Chapter 5).

Figure 6.1 plots the dissipated energy per conversion versus sensor resolution of various smart temperature sensors [6.1]. The different sensor prototypes developed during this research are denoted by the bold symbols in this plot. As shown, the efficiency of smart sensors published prior to the start of this research (in 2009) was limited to about  $1\text{nJ}^\circ\text{C}^2$ , denoted by the solid FoM line. Since then, the work described in this thesis (among others) has improved the energy efficiency of BJT-based temperature sensors by a remarkable two orders of magnitude, represented by the dashed FoM line of  $10\text{pJ}^\circ\text{C}^2$ . One of the designs in this work, achieved a resolution FoM of  $11\text{pJ}^\circ\text{C}^2$ , which defined the state-of-the-art when it was published in 2012 [6.2]. Furthermore, this research presented the first DTMOST-based smart temperature sensors, which achieved the highest relative accuracy in the class of MOSFET-based temperature sensors, after only a single-temperature calibration [6.3], [6.4]. To date, the work presented in [6.4] is the only all-CMOS temperature sensor to achieve both sub-1V operation, and a high accuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) over a wide temperature range ( $-40^\circ\text{C}$  to  $125^\circ\text{C}$ ), while employing

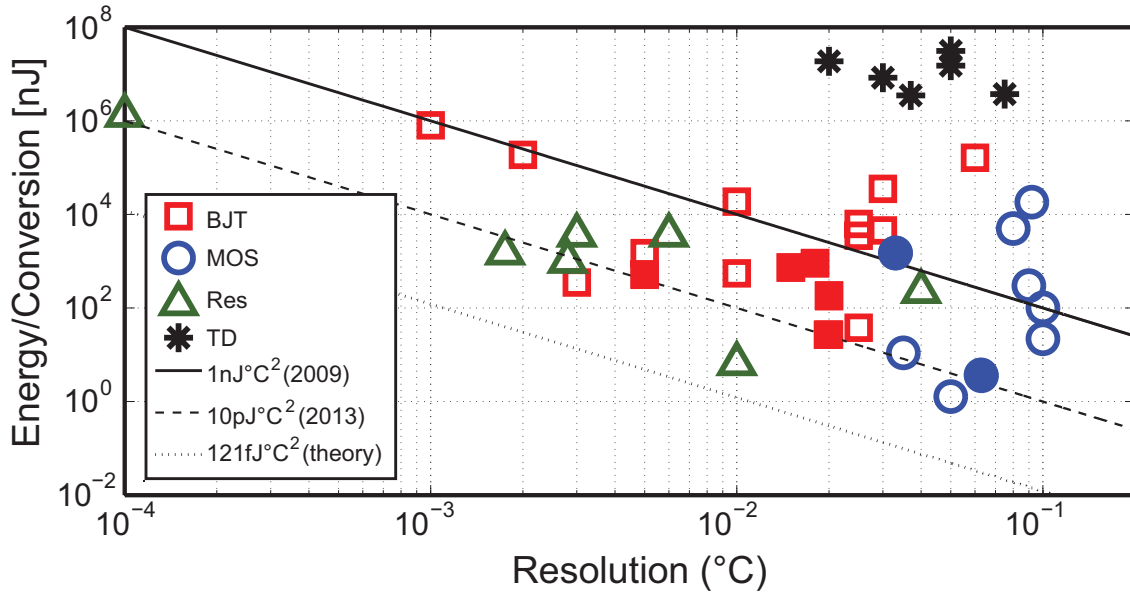


Figure 6.1: Energy per conversion versus resolution for various smart temperature sensors [6.1]. The solid and dashed FoM lines represent the state-of-the-art in 2009 and 2013, respectively. The solid symbols indicate the various sensors realized during this research. The theoretical FoM line of a BJT-based front-end using the (worst-case) bias current scaling ( $p = 2$ ) is illustrated by the dotted line (section 2.4.1).

only a single room-temperature trim. It also achieves good energy-efficiency: with a FoM of  $14\text{pJ}^\circ\text{C}^2$ .

Last but not least, different flavors of the first BJT-based sensor prototype (Chapter 4) were productized by NXP Semiconductors, namely the PCT2075 [6.5] and the PCT2202UK [6.6]. Moreover, the first prototype was also integrated into a multi-sensor platform for smart RFID tag applications, during the EU-funded Pasteur project led by NXP Semiconductors. The tag is used to predict the expiration date of perishable goods, based on the history of conditions (including temperature), in which the items are stored and transported [6.7]. Future products based on the  $2^{\text{nd}}$ -order zoom-ADC are in the pipeline.

## 6.2 Other Applications of this Work

The zoom-ADC technique described in this thesis can also be used to improve the energy-efficiency of the general-purpose ADCs. However, one drawback of the proposed architecture is that the coarse and fine conversions are performed sequentially. Therefore, the references of the fine  $\Delta\Sigma$ -ADC are set based on the results of the coarse SAR conversion,

and thus are fixed throughout the fine conversion. This, in turn, imposes a limitation on the input signal bandwidth, which must remain within certain limits during the fine conversion step, to avoid out-of-range errors. This property matches the requirements of most instrumentation and sensor readout applications, in which high resolution/accuracy and energy-efficiency are key requirements, while the input signal bandwidth is usually low, e.g. less than 100Hz. To demonstrate this, a general-purpose incremental zoom-ADC was developed [6.8]. It combines a 6-bit coarse SAR conversion with a 15-bit fine  $\Delta\Sigma$ -ADC, and employs a fully inverter-based structure. In a conversion time of 40ms, it achieves a resolution of 20 bits, while drawing  $6\mu\text{A}$  from a 1.8V supply. By the use of dynamic error correction techniques, it achieved an offset of  $1\mu\text{V}$  and an INL of 6ppm. It also obtained a Schreier FoM of 182.7dB, which defined the state-of-the-art in the class of incremental ADCs, when it was published in 2013 [6.8].

In a recent work [6.9], parallel operation of the coarse and fine conversions in the zoom-ADC is presented, which allows the input signal bandwidth to be increased without causing out-of-range errors. Since the coarse SAR-ADC is continuously running, the references of the fine  $\Delta\Sigma$ -ADC are dynamically updated, and thus rapidly track the input signal. The resulting zoom-ADC achieves a dynamic range of 107.5dB in a signal bandwidth of 20kHz, which meets the requirements of audio codecs. By drawing 0.92mA from a 1.8V supply, it achieves a competitive Schreier FoM of 178.3dB, but occupies significantly smaller area ( $0.16\text{mm}^2$ ), when compared to state-of-the-art ADCs with similar performance. This is due to the advantages offered by the zoom-ADC architecture, which enables a simple circuit design and a compact silicon realization.

### 6.3 Future Work

As presented, the BJT-based second prototype in Chapter 4 achieves a FoM of  $11\text{pJ}^\circ\text{C}^2$ . As described in Chapter 2, the theoretical energy-efficiency limit for the BJT-based sensors is  $121\text{fJ}^\circ\text{C}^2$  (see Figure 6.1), which is more than one order of magnitude smaller than that of state-of-the-art sensors. This is partially because in the theoretical FoM calculations of Chapter 2, the power dissipated in the front-end's bias circuit and ADC were neglected. After applying the same assumption to our sensor prototype, however, results in a FoM of  $3.5\text{pJ}^\circ\text{C}^2$ , which is still much larger than the theoretical limit. The main reason for this large gap, is the fact that the sensor's resolution is dominated by the  $kT/C$  noise of its switched-capacitor readout. Considering the switched-capacitor  $2^{\text{nd}}$ -order zoom-ADC presented in Chapter 4, and assuming that the output noise is dominated by the sampling

of  $\Delta V_{BE}$ , the associated voltage noise during the fine conversion can be approximated by:

$$v_{n,\Delta V_{BE}}^2 \approx \frac{kT}{(n+1) \cdot C_U \cdot N_{\Delta\Sigma}}, \quad (6.1)$$

where  $k$  and  $T$  represent the Boltzmann constant and temperature in Kelvin, respectively. The parameter  $n$  is the integer part of  $X$ , as presented in equation (3.7),  $C_U$  is the DAC's unit sampling capacitor, and  $N_{\Delta\Sigma}$  is the number of  $\Delta\Sigma$  cycles during the fine conversion step. For the second prototype presented in Chapter 4,  $n = 14$  at room temperature,  $C_U = 120\text{fF}$ , and  $N_{\Delta\Sigma} = 128$ . By applying these values to equation (6.1) and after multiplying by the sensitivity  $S_{\Delta V_{BE}}^{D_{out}}(T)$  (see equation 2.19), an output noise of  $\approx 16\text{mK}$  can be calculated, which is in-line with the measurement results presented in section 4.2.4. Therefore, even by neglecting the power dissipation in the bias circuit and the ADC, it is not possible to achieve a FoM better than  $2.2\text{pJ}^\circ\text{C}^2$ . Increasing the size of  $C_U$  to reduce noise will not improve things, as it will require a proportionally longer settling time, and thus more energy dissipation. This observation suggests that in order to further improve energy-efficiency, a continuous-time (CT) readout should be used, thus avoiding the noise aliasing inherent to the use of switched-capacitor circuits. The challenge will then be to tackle the resulting sensitivity to sampling clock jitter, and how to implement the dynamic element matching needed for accuracy.

As discussed in Chapter 4, and can be seen from equation (2.9), the lot-to-lot spread observed in the fitting parameters  $A$  and  $B$  corresponds to a major loss of accuracy, e.g. a shift of  $0.5^\circ\text{C}$  was observed between two lots, at room temperature. Over the full temperature range and across multiple process lots, this error is expected to be even larger. To overcome this, the optimum  $A$  and  $B$  for each batch should be determined via a batch calibration, i.e. by characterizing a few samples from each batch over the full temperature range. Better accuracy can be obtained by an individual calibration, preferably only at room-temperature. Voltage calibration, as described in [6.10], can be used to significantly speed up such room-temperature calibration. Nevertheless, characterization over temperature, as required for the batch calibration, will still dictate the use of climate chambers or oil-baths. The associated calibration time, and therefore extra cost, is a major limiting factor on the accuracy of commercial products. The further development of calibration techniques based on electrical measurements, in order to eliminate the need for over-temperature characterization will be extremely valuable. One can think of combining voltage calibration with integrated heaters to generate on-chip thermal variations, and thereby extracting the parameters  $A$  and  $B$ .

Although voltage calibration is an effective alternative to thermal calibration, it still requires accurate knowledge of the externally applied reference voltage, e.g. a  $\pm 200\mu\text{V}$



error corresponds to a calibration error of  $\pm 0.1^\circ\text{C}$  [6.10]. Such accurate voltage measurement is not always possible in standard high-volume test equipment, and may call for customized test setups, which is not desirable. Since thermal-diffusivity (TD) sensors are apparently able to achieve high accuracy without calibration [6.11], an alternative method would be to use them as on-die reference sensors during the calibration process. This would eliminate the need for precision voltage measurements by the tester. The TD sensor's relatively large operating power would then only be dissipated during the calibration step. The challenge, however, will be to design a single readout circuit which can be efficiently shared between the TD and the main sensor, e.g. a BJT- or DTMOST-based sensor, with minimal power and area overhead. It should be noted, however, that the uncalibrated accuracy of TD sensors over different process lots has not yet been demonstrated, and requires further investigations.

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# Summary

Nowadays, smart temperature sensors, i.e. sensors with digital outputs are widely used in various systems. Integrating smart sensors into wireless systems such as RFID tags or wireless sensor networks (WSNs) enables wireless temperature sensing, which in turn opens up a wide range of new applications. This thesis describes the requirements, design, and implementation of smart temperature sensors for use in wireless temperature sensing.

In Chapter 1, an introduction to wireless temperature sensing and its requirements is given. Typically, a wireless node is either powered by a battery or scavenges its energy from the environment, e.g. from an external RF magnetic field. Due to the limited amount of energy available, energy-efficiency of the integrated sensor either restricts the battery's lifetime, or the operating range of the wireless node. On the other hand, mass production imposes stringent requirements on the cost, which calls for CMOS compatible sensors. To obtain sufficient accuracy, however, CMOS sensors often require time-consuming (and thus costly) *calibration*: a process in which the sensor's output is compared with that of a reference sensor at a number of known temperatures. The information obtained during calibration is then used to *trim* the sensor, thereby improving its accuracy. A short survey of various CMOS compatible choices is presented. It is shown that substrate PNP's are suitable candidates for wireless temperature sensing. They are power-efficient and exhibit a well-defined process spread, which can be effectively trimmed at a single-temperature. However, they require supply voltages greater than  $\approx 1.2V$ , making them ill-suited for low-voltage applications and nano-scale CMOS processes. A promising alternative is to bias a MOSFET in the sub-threshold region, while its body and gate terminals are shorted. This so-called DTMOST configuration enables sub-1V operation while exhibiting less spread when compared to the bulk configuration. Finally, to facilitate the comparison between energy-efficiency of various temperature sensors, a single figure of merit (FoM) is presented.

In Chapter 2, the operating principle of BJT-based smart temperature sensors is presented. Using the parasitic BJTs available in CMOS, a complementary-to-absolute-

temperature (CTAT) voltage  $V_{BE}$  and a proportional-to-absolute-temperature (PTAT) voltage  $\Delta V_{BE}$  can be generated. By properly scaling  $\Delta V_{BE}$  (with a scalar  $\alpha$ ) and combining it with  $V_{BE}$ , a reference voltage  $V_{REF}$  can then be obtained. In a generic BJT readout, the ratio of  $\alpha \cdot \Delta V_{BE}$  and  $V_{REF}$  is digitized by means of an analog-to-digital converter (ADC) to generate a PTAT ratio  $\mu$ . The resolution requirement of the ADC is also discussed. It is shown that almost  $\frac{2}{3}$  of the ADS's dynamic range is wasted with this approach. To identify the energy-efficiency of existing sensors prior to the start of this research, a study of energy-efficiency limits in BJT-based sensors is presented. In this analysis, the ultimate energy-efficiency of a BJT-based sensor front-end is calculated and the theoretical limits are defined. Two different approaches based on bias-current and emitter-area scaling are considered. Based on this analysis, a significant energy-efficiency gap, over four orders of magnitude, is observed between the prior-art and theoretical limits. The study of various sensor architectures reveals that, in fact, the reason behind this gap lies in the employed readout circuits, which mostly include  $\Delta\Sigma$ - or SAR-ADCs. They either suffer from long conversion times and poor power-efficiency, or are not capable of providing the target resolution or accuracy. To bridge this efficiency gap, a new readout architecture is clearly required.

In Chapter 3, different BJT-based sensor architectures based on digitizing non-linear ratios between  $\Delta V_{BE}$  and  $V_{BE}$  (or their combinations) are explored. The required linearization to calculate the PTAT ratio  $\mu$  is then performed in the digital back-end. Since the coefficient  $\alpha$  is digitally implemented, it can also be used for trimming. The employed ADC architectures in these examples, however, often result in more waste of dynamic range than in the generic approach, exacerbating the lack of energy-efficiency. To address this issue, a new readout topology based on digitizing the ratio  $X = V_{BE}/\Delta V_{BE}$  is proposed. Since temperature changes are rather slow, the ratio  $X$  is accurately digitized by a two-step *zoom*-ADC. As  $X$  is typically greater than one, it can be expressed as  $X = n + \mu'$ , where  $n$  and  $\mu'$  correspond to the integer and fractional parts, respectively. First, a full-range SAR conversion obtains the integer  $n$  by performing a binary search algorithm, comparing  $V_{BE}$  to integer multiples of  $\Delta V_{BE}$ . This is then followed by a low-range fine  $\Delta\Sigma$  converter, whose references are set to  $n$  and  $n+1$ . In this manner, the ratio  $\mu'$  can be accurately digitized with high resolution. In contrast to the conventional  $\Delta\Sigma$ -ADCs, the full-scale range of the fine converter in the zoom-ADC is considerably reduced, which notably relaxes various key requirements such as the number of  $\Delta\Sigma$ -cycles and the DC-gain and swing of the loop-filter. In this architecture, both conversion time and power-efficiency can be improved, which results in a substantial improvement in energy-efficiency. The fact that dynamic correction techniques can be used in the fine conversion phase, ensures that the accuracy of the zoom-ADC can be as good as that of conventional  $\Delta\Sigma$ -ADC architectures.

In Chapter 4, a low-power BJT-based sensor prototype based on a 1<sup>st</sup>-order switched-capacitor (SC) zoom-ADC is presented. It achieves a resolution of 15mK in a conversion time of 100ms while dissipating only 4.6 $\mu$ A. After a single  $\alpha$ -trim at 25°C, the sensor obtains an inaccuracy of  $\pm 0.2^\circ\text{C}$  ( $3\sigma$ ) from -30°C to 125°C. This result shows 11x energy-efficiency improvement when compared to sensors with similar accuracy, back in 2011. However, its fine conversion step employs a slow, 1<sup>st</sup>-order  $\Delta\Sigma$  modulator, limiting its energy-efficiency. Moreover, each  $\Delta\Sigma$  cycle requires two full clock periods, since  $V_{BE}$  and  $\Delta V_{BE}$  are separately sampled/integrated. To further improve the sensor's energy-efficiency, a second prototype is realized which achieves similar resolution in about 16x less conversion time, while drawing 25% less supply current. This is achieved by using a 2<sup>nd</sup>-order zoom-ADC, combined with a new charge-balancing scheme, whose operation is based on *simultaneous* sampling of  $V_{BE}$  and  $\Delta V_{BE}$ . This allows the use of low-swing, low-power amplifiers. The sensor's energy-efficiency is therefore improved by over 20x compared to the first prototype. Using a thermal calibration and digital PTAT trimming at 30°C, the sensor achieves an inaccuracy of  $\pm 0.15^\circ\text{C}$  ( $3\sigma$ ) from -55°C to 125°C. Moreover, a voltage calibration technique based on electrical measurements is also explored, which is significantly faster (only requires 200ms), while achieving comparable accuracy. The impact of batch-to-batch spread and plastic packaging on sensor's accuracy is investigated as well. As observed, both of them can cause temperature reading shifts in the order of  $0.4^\circ\text{C} \sim 0.5^\circ\text{C}$  from -55°C to 125°C.

In the last part of Chapter 4, a BJT-based sensor prototype for sensing high temperatures ( $>150^\circ\text{C}$ ) is also demonstrated. It is shown that by optimizing the emitter area and bias current of a substrate PNP, the impact of saturation current  $I_S$  at high temperatures can be mitigated. Furthermore, robust circuit techniques are employed to cope with the various leakage currents at such temperatures, which would otherwise impact the accuracy of  $V_{BE}$  and  $\Delta V_{BE}$ , and thus the sensor output. It achieves an inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) from -55°C to 200°C, which is similar to that of state-of-the-art sensors capable of operating over such temperature ranges. However, it draws only 22 $\mu$ A, which is more than an order of magnitude less.

In Chapter 5, the use of DTMOSTs as temperature sensing elements is demonstrated. When operated in weak inversion, the gate-source voltage  $V_{GS}$  of a DTMOST is almost half of the base-emitter voltage  $V_{BE}$ , thus enabling sub-1V operations. Moreover, compared to a diode-connected MOSFET, the  $V_{GS} - I_D$  characteristic of a diode-connected DTMOST is less sensitive to the spread in threshold voltage  $V_T$ , making it a promising candidate for realizing accurate temperature sensors. Two sensor prototypes based on such sensing elements are demonstrated in a chosen 160nm CMOS process. After a single-temperature trim, the first prototype achieves an inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) from -55°C to 125°C, and enables an apples-to-apples comparison with BJTs, proving that

DTMOSTs are indeed only a factor 2x less accurate. In the second prototype, the low-voltage capability of DTMOSTs is then exploited to realize a sub-1V, sub- $\mu$ W precision sensor. Employing fully inverter-based SC integrators, a 2<sup>nd</sup>-order zoom-ADC is realized in the second prototype. It can operate at supply voltages as low as 0.85V, while drawing only 700nA. It also maintains the same inaccuracy of  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) from  $-40^\circ\text{C}$  to  $125^\circ\text{C}$ , after a single-temperature trim. These results prove that DTMOSTs could be considered as the temperature sensors of choice when sub-1V, high accuracy, and energy-efficiency are key requirements.

In Chapter 6, the main findings of this work are summarized. These include the development of the zoom-ADC and its application in energy-efficient smart temperature sensors. The final prototype BJT-based sensor achieved a resolution FoM of  $11\text{pJ}^\circ\text{C}^2$  and improved state-of-the-art by a factor of 15x (in 2012). Another key finding was the fact that DTMOST sensors enable low voltage operations while being only 2x less accurate than BJT-based sensors. The final prototype achieved a fairly good energy-efficiency, evidenced by a FoM of  $14\text{pJ}^\circ\text{C}^2$ . The chapter also contains some suggestions for future work: to further improve the energy-efficiency, continuous-time (CT) readouts could be considered as promising alternatives to the switched-capacitor circuits. Furthermore, to reduce the cost of over-temperature characterizations, a combination of voltage calibration with integrated heaters could be used to quickly extract the global calibration parameters. Another alternative could be to exploit the high accuracy of thermal-diffusivity (TD) sensors as on-die references during the calibration process. The chapter ends with a discussion of the potential use of the zoom-ADC technique to realize general-purpose ADCs with high energy-efficiency.

# Samenvatting

*Smart* temperatuursensoren, met digitale uitgangen, worden tegenwoordig veel toegepast in allerlei systemen. Door smart sensors te integreren in draadloze systemen, zoals RFID tags of draadloze sensor netwerken, wordt draadloze temperatuurmeting mogelijk. Hierdoor ontstaat een breed scala aan nieuwe toepassingen. Dit proefschrift beschrijft de vereisten, het ontwerp en de implementatie van smart temperatuursensoren, toegepast in draadloze temperatuurmeting.

Hoofdstuk 1 geeft een inleiding in draadloze temperatuurmeting en de vereisten daarvoor. Meestal wordt daarbij gebruik gemaakt van een draadloze node, gevoed door een batterij, of door energie die uit de omgeving wordt opgevangen, bijvoorbeeld uit een RF elektromagnetisch veld. De geringe hoeveelheid beschikbare energie beperkt de levensduur van de batterij, of het bereik van de draadloze node. Aan de andere kant stelt massaproductie hoge eisen aan de kosten, zodat CMOS compatibele sensors de aangewezen oplossing lijken te zijn. Voor voldoende nauwkeurigheid vergen CMOS sensors veelal tijdrovende en dus dure *kalibratie*: daarbij worden de outputs van de sensor en een referentiesensor vergeleken bij een aantal bekende temperaturen. De informatie die tijdens kalibratie wordt verkregen wordt benut om de sensor af te regelen en zo zijn nauwkeurigheid te vergroten. In een kort overzicht worden diverse CMOS-compatibele keuzes gepresenteerd. Substraat PNP's blijken bruikbare kandidaten voor draadloze temperatuurmeting. Ze zijn energiezuinig en hebben een goed gedefinieerde processpreiding, die bij een enkele temperatuur afgeregeld kan worden. Ze hebben echter een voedingsspanning van meer  $\approx 1.2V$  of meer nodig, waardoor ze slecht geschikt zijn voor laagspanningstoepassingen en CMOS processen op nanoschaal. Een veelbelovend alternatief is het gebruik van een MOSFET in sub-threshold, met kortgesloten body en gate. Deze zogeheten DTMOST configuratie is geschikt voor voedingsspanningen onder 1V. Die configuratie vertoont minder spreiding dan de bulk configuratie. Tenslotte wordt een single figure of merit (FoM) gepresenteerd, teneinde een vergelijking tussen diverse temperatuursensoren mogelijk te maken.

Hoofdstuk 2 beschrijft het werkingsprincipe van op BJT-gebaseerd smart temperatu-



ursensoren. Een CTAT (complementary-to-absolute-temperature) spanning  $V_{BE}$  en een PTAT (proportional-to-absolute-temperature) spanning  $\Delta V_{BE}$  worden opgewekt door gebruik te maken van de parasitaire BJTs die in CMOS beschikbaar zijn. Door  $\Delta V_{BE}$  te schalen met constante  $\alpha$  en het resultaat te combineren met  $V_{BE}$  wordt een referentiespanning  $V_{REF}$  verkregen. In een generieke BJT uitlezing wordt de verhouding van  $\alpha \cdot \Delta V_{BE}$  en  $V_{REF}$  gedigitaliseerd met een ADC (analog-to-digital converter) om daarmee een PTAT verhouding  $\mu$  te bepalen. De vereiste ADC resolutie wordt ook besproken. Het blijkt dat bijna  $\frac{2}{3}$  van het dynamisch bereik van de ADC met bovenstaande benaderingswijze ongebruikt blijft. Een studie naar de grenswaarden van het energierendement bij sensoren die op BJTs gebaseerd zijn en die reeds voor deze studie bestonden wordt gepresenteerd. Deze analyse berekent het uiteindelijke energierendement en de theoretische limieten van een op BJT-gebaseerd sensor. Twee benaderingswijzen op basis van instelstroom en emitter oppervlakteschaling worden beschouwd. Op basis van deze analyse wordt een significant verschil van vier grootteordes tussen de bestaande technieken en de theoretische limieten geconstateerd. De studie van verscheidene sensor architecturen laat zien dat de reden voor dit significante verschil in de gebruikte uitleescircuits ligt, die meestal  $\Delta\Sigma$ - of SAR-ADCs bevatten. Deze hebben ofwel een slecht energierendement en een lange conversietijd ofwel ze zijn niet in staat de vereiste resolutie of nauwkeurigheid te leveren. Om dit gemis aan rendement goed te maken is duidelijk een nieuwe architectuur nodig voor het uitleescircuit.

Hoofdstuk 3 verkent diverse BJT-gebaseerde sensorarchitecturen die gebruik maken van gedigitaliseerde niet lineaire verhoudingen tussen (combinaties van)  $\Delta V_{BE}$  en  $V_{BE}$ . De benodigde linearisering voor de bepaling van PTAT verhouding  $\mu$  wordt uitgevoerd in het digitale back-end. Omdat de constante  $\alpha$  digitaal wordt geïmplementeerd, kan deze ook voor afregeldoelinden gebruikt worden. De ADC architecturen die in deze voorbeelden worden toegepast resulteren echter in een grotere verspilling van dynamisch bereik dan de generieke benaderingswijze, waardoor het gemis aan energierendement verergerd wordt. Om dit probleem aan te pakken, wordt een nieuwe topologie voor de uitlezing voorgesteld, die is gebaseerd op het digitaliseren van de verhouding  $X = V_{BE} / \Delta V_{BE}$ . Aangezien de temperatuur tamelijk langzaam verandert, wordt de verhouding  $X$  nauwkeurig gedigitaliseerd door een twee-staps zoom-ADC.  $X$  is typisch groter dan een. Daarom kan  $X$  uitgedrukt worden als  $X = n + \mu'$ , waarbij  $n$  en  $\mu'$  respectievelijk corresponderen met het gehele en het breukdeel. Ten eerste bepaalt een full-range SAR conversie het gehele deel  $n$ , met behulp van een binair zoekalgoritme, waarbij  $V_{BE}$  vergeleken wordt met veelvoud van  $\Delta V_{BE}$ . Hierop volgt dan een converter met klein bereik en hoge resolutie, waarvan de referenties ingesteld worden op  $n$  en  $n+1$ . Op die manier wordt de verhouding  $\mu'$  nauwkeurig gedigitaliseerd met een hoge resolutie. Anders dan bij conventionele  $\Delta\Sigma$ -ADCs wordt het volle schaalbereik van de hoge resolutie converter in de zoom-ADC

aanzienlijk verkleind. Dit betekent een aanzienlijke versoepeling van de vereisten, zoals het aantal  $\Delta\Sigma$ -cycli, DC-gain en signaalzwaai van het lusfilter. Met deze architectuur kunnen zowel conversietijd als vermogensrendement verbeterd worden, wat resulteert in een substantiële verbetering van het energierendement. Doordat dynamische correctietechnieken kunnen worden gebruikt in de grote resolutie conversiefase, kan de nauwkeurigheid van de zoom-ADC net zo goed zijn als in conventionele  $\Delta\Sigma$ -ADC architecturen.

Hoofdstuk 4 beschrijft een laagvermogen BJT-gebaseerd sensorprototype, gebaseerd op een eerste-orde switched capacitor (SC) zoom-ADC. Deze sensor heeft een resolutie van 15mK bij een conversietijd van 100ms, met een stroomverbruik van slechts  $4.6\mu\text{A}$ . Na een enkele -afregeling bij  $25^\circ\text{C}$ , heeft de sensor een onnauwkeurigheid van  $\pm 0.2^\circ\text{C}$  ( $3\sigma$ ) bij een temperatuurbereik van  $-30^\circ\text{C}$  tot  $125^\circ\text{C}$ . Dit resulteert in een 11x beter energierendement in vergelijking met sensoren met vergelijkbare nauwkeurigheid, zoals die in 2011 bestonden. De sensor heeft echter een hoge resolutie conversie stap, die een langzame eerste-orde  $\Delta\Sigma$  modulator gebruikt, wat het energierendement verlaagt. Daar komt bij dat elke  $\Delta\Sigma$  cyclus twee hele klokperiodes nodig heeft omdat  $V_{BE}$  en  $\Delta V_{BE}$  elk apart worden bemonsterd en geïntegreerd. Om het energierendement van de sensor verder te verbeteren is een tweede prototype gemaakt, met een vergelijkbare resolutie maar met een 16x kleinere conversietijd, dat 25% minder voedingsstroom nodig heeft. Dit wordt bereikt door een tweede orde zoom-ADC te gebruiken in combinatie met een nieuwe wijze van ladingsbalancering, gebaseerd op *gelijktijdig* bemonsteren van  $V_{BE}$  en  $\Delta V_{BE}$ . Hierdoor kunnen low-power versterkers met een kleine uitgangszwaai worden gebruikt. Zo wordt het energierendement van de sensor 20x beter dan het eerste prototype. Met een thermische kalibratie en digitale PTAT afregeling bij  $30^\circ\text{C}$ , heeft de sensor een onnauwkeurigheid van  $\pm 0.15^\circ\text{C}$  ( $3\sigma$ ) bij een temperatuurbereik van  $-55^\circ\text{C}$  tot  $125^\circ\text{C}$ . Verder wordt een op elektrische metingen gebaseerde spanningskalibratietechniek verkend, die aanzienlijk sneller is (200ms) en toch een vergelijkbare nauwkeurigheid bereikt. De invloed van processpreiding tussen batches en de invloed van een plastic behuizing op de nauwkeurigheid van de sensor wordt tevens onderzocht. Waargenomen is dat elk van deze twee invloeden verschuivingen in de afgelezen temperatuur kan veroorzaken in de orde van  $0.4^\circ\text{C} \sim 0.5^\circ\text{C}$  bij een temperatuurbereik van  $-55^\circ\text{C}$  tot  $125^\circ\text{C}$ .

In het laatste gedeelte van hoofdstuk 4 wordt BJT-gebaseerd sensorprototype voor hoge temperatuur meting ( $>150^\circ\text{C}$ ) getoond. Het blijkt dat optimalisatie van emitter oppervalk en biasstroom van een substraat PNP het effect van verzadigingsstroom  $I_S$  bij hoge temperaturen kan verminderen. Vervolgens worden robuuste circuit ontwerpstechnieken gebruikt om diverse lekstromen bij zulke hoge temperaturen te hanteren, die anders de nauwkeurigheid van  $V_{BE}$  en  $\Delta V_{BE}$  zouden aantasten en dus ook de sensor aflezing. Zo wordt een nauwkeurigheid van  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) bij een temperatuurbereik van  $-55^\circ\text{C}$  tot  $200^\circ\text{C}$  bereikt. Zulke waarden worden ook bereikt met state-of-the-art sensors die over

dergelijke temperatuur bereiken kunnen werken. De gepresenteerde sensor gebruikt echter slechts  $22\mu\text{A}$ , wat meer dan een grootteorde beter is dan de state-of-the-art sensoren.

Hoofdstuk 5 toont de toepassing van DTMOSTs als temperatuurgevoelige elementen. In zwakke inversie is de gate-source spanning  $V_{GS}$  van een DTMOST ongeveer de helft van basis-emitter spanning  $V_{BE}$ , waardoor het sub-1V bereik haalbaar wordt. Wanneer we de  $V_{GS} - I_D$  karakteristiek van een als diode geschakelde MOSFET vergelijken met een als diode geschakelde DTMOST, is de DTMOST minder gevoelig voor spreiding in drempelspanning  $V_T$ , waardoor de DTMOST een veelbelovende kandidaat wordt voor het realiseren van nauwkeurige temperatuursensoren. Twee prototypes die op DTMOST gebaseerd zijn, zijn gerealiseerd in een 160nm CMOS proces. Na afregeling bij een enkele temperatuur heeft het eerste prototype een onnauwkeurigheid van  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) bij een temperatuurbereik van  $-55^\circ\text{C}$  tot  $125^\circ\text{C}$ . Dit prototype maakt een eerlijke vergelijking met BJTs mogelijk, die aantoont dat DTMOSTs inderdaad 2x minder nauwkeurig zijn. Het tweede prototype realiseert een sub-1V, sub- $\mu\text{W}$  precisie sensor door gebruik te maken van DTMOSTs die immers een lage spanning nodig hebben. Met gebruikmaking van geheel op inverters gebaseerde SC integratoren wordt in het tweede prototype een tweede-orde zoom-ADC gerealiseerd. Dit prototype werkt op voedingsspanningen van 0.85V en hoger en gebruikt slechts 700nA. De nauwkeurigheid is weer  $\pm 0.4^\circ\text{C}$  ( $3\sigma$ ) bij een temperatuurbereik van  $-40^\circ\text{C}$  tot  $125^\circ\text{C}$ , na een afregeling bij een enkele temperatuur. Deze resultaten bewijzen dat DTMOSTs de voorkeur hebben wanneer sub-1V, grote nauwkeurigheid en hoog energierendement hoofdvereisten zijn.

Hoofdstuk 6 vat de belangrijkste bevindingen van dit werk samen. Deze betreffen de ontwikkeling van de zoom-ADC en toepassing daarvan in energiezuinige smart temperatuursensoren. Het uiteindelijke prototype van een BJT-gebaseerde sensor heeft een resolutie FoM van  $11\text{pJ}^\circ\text{C}^2$  en overtreft state-of-the-art (in 2012) met een faktor 15x. Een andere bevinding is dat DTMOST sensoren werken bij lage spanningen terwijl ze slechts tweemaal minder nauwkeurig zijn dan op BJT-gebaseerde sensoren. Het uiteindelijke DTMOST prototype heeft een redelijk goed energierendement met een FoM van  $14\text{pJ}^\circ\text{C}^2$ . Hoofdstuk 6 geeft voorts een aantal suggesties voor verder onderzoek: om het energierendement verder te verbeteren zou tijd-continue (CT) uitlezing overwogen kunnen worden. Dit is een veelbelovend alternatief voor switched-capacitor circuits. Verder zouden de kosten van karakterisering over temperatuur omlaag gebracht kunnen worden door toepassing van spanningskalibratie en geïntegreerde warmtebronnen om zo snel de globale kalibratie parameters te bepalen. Een verder alternatief kan het gebruik van de grote nauwkeurigheid thermal-diffusivity (TD) sensors zijn, als on-die referenties tijdens het kalibratieproces. Het hoofdstuk eindigt met een bespreking van het mogelijk gebruik van de zoom-ADC techniek om general-purpose ADCs met een hoog energierendement te maken.

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# List of Publications

## Journal Papers

- [1] Y. Chae, **K. Souri** and K. A. A. Makinwa, "A  $6.3\mu\text{W}$  20bit incremental zoom-ADC with 6ppm INL and  $1\mu\text{V}$  offset," *IEEE Journal of Solid-State Circuits*, vol. 48, No. 12, pp. 3019 - 3027, Dec. 2013.
- [2] **K. Souri**, Y. Chae and K. A. A. Makinwa, "A CMOS temperature sensor with a voltage-calibrated inaccuracy of  $\pm 0.15^\circ\text{C}$  ( $3\sigma$ ) from  $-55$  to  $125^\circ\text{C}$ ," *IEEE Journal of Solid-State Circuits*, vol. 48, No. 1, pp. 292 - 301, Jan. 2013.
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## Conference Papers

- [1] **K. Souri**, Y. Chae, F. Thus and K. A. A. Makinwa, "A 0.85V 600nW all-CMOS temperature sensor with an inaccuracy of  $\pm 0.4^{\circ}\text{C}$  ( $3\sigma$ ) from  $-40$  to  $125^{\circ}\text{C}$ ," in *Dig. Techn. Papers ISSCC*, pp. 222 - 223, Feb. 2014.
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## About the Author



**Kamran Souri** was born in Tabriz, Iran in 1980. He received his B.Sc. in Electronics and M.Sc. in Telecommunication Systems from Amirkabir University of Technology, Iran, in 2001 and 2004, respectively. From 2001 to 2007, he worked at PSP-Ltd, Tehran, Iran, designing embedded systems for use in high-quality audio/video systems and KVM switches. In Sept. 2007 he joined the Electronic Instrumentation Laboratory (EI-Lab), TU-Delft, where he received the M.Sc. degree in Microelectronics (cum laude) in 2009. Between 2007 and 2009 he was the recipient of a Top Talent Fellowship from the same faculty. From 2008 to 2009, he was an intern at NXP-Semiconductors, Eindhoven, designing energy-efficient temperature sensors for use in RFID tags. Between 2009 and 2014 he was a Ph.D. candidate at EI-Lab, TU-Delft, working on energy-efficient smart temperature sensors in CMOS technology. He was the recipient of the IEEE Solid-State Circuits Society Predoctoral Achievement Award in 2013. He is currently a Principal Circuit Design Engineer at SiTime Corp., Sunnyvale, CA.