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THE WEYL CALCULUS FOR GROUP GENERATORS SATISFYING THE CANONICAL COMMUTATION RELATIONS

JAN VAN NEERVEN and PIERRE PORTAL

*In memory of Alan McIntosh (1942-2016), celebrating his friendship with
José Enrique Moyal (1910-1998)*

Communicated by Florian-Horia Vasilescu

ABSTRACT. We generalise the classical Weyl pseudo-differential calculus on \mathbb{R}^d to the setting of two d -tuples of operators $A = (A_1, \dots, A_d)$ and $B = (B_1, \dots, B_d)$ acting on a Banach space generating bounded C_0 -groups satisfying the Weyl canonical commutation relations. We show that the resulting Weyl calculus extends to symbols in the standard symbol class S^0 provided appropriate bounds can be established. Using transference techniques we obtain boundedness of the H^∞ -functional calculus (and even the Hörmander calculus), for the abstract harmonic oscillator $L = \frac{1}{2} \sum_{j=1}^d (A_j^2 + B_j^2) - \frac{1}{2}d$.

KEYWORDS: *Weyl pairs, canonical commutation relations, pseudo-differential calculus, twisted convolution, transference of C_0 -groups, UMD spaces, H^∞ -functional calculus, spectral multipliers.*

MSC (2010): Primary 47A60; Secondary 35S05, 47A13, 47D03, 47G30, 81S05.

1. INTRODUCTION

In the early 1980's, Alan McIntosh introduced the H^∞ -functional calculus as a refined version of the Dunford holomorphic functional calculus for unbounded sectorial operators (see the original paper [43] in the Hilbert space setting, their extensions to Banach spaces [15], [35], and the monographs [25], [33]). This calculus is meant to be an operator-theoretic abstraction of the calculus of Fourier multipliers, which it recovers when applied to constant coefficient differential operators such as the Laplacian on $L^2(\mathbb{R}^d)$. One of its roles is to provide a framework for perturbation theory: deriving properties of the functional calculus of differential operators with varying coefficients from its constant coefficient counterpart. The quintessential example of such an application is given in [6], where perturbation is first understood in the operator-theoretic sense, then in the harmonic

analytic sense of (an extension of) Calderón–Zygmund theory. The combination of both perspectives leads to a striking boundedness result for the H^∞ -functional calculus of Dirac operators that includes the solution of the celebrated Kato’s square root problem (originally obtained in [5]).

In the present paper, we introduce an operator-theoretic framework which aims to generalise pseudo-differential calculus in the same way that the McIntosh H^∞ -functional calculus generalises Fourier multiplier calculus. Our starting point is the Weyl calculus of standard position and momentum operators $Q_j f(x) = x_j f(x)$ and $P_j f(x) = i\partial_j f(x)$, $j = 1, \dots, d$, acting on their natural domains in $L^2(\mathbb{R}^d)$. For Schwartz functions $a \in \mathcal{S}(\mathbb{R}^{2d})$ one can define a bounded operator $a(Q, P)$ acting on $L^2(\mathbb{R}^d)$ by

$$(1.1) \quad a(Q, P)f = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{a}(u, v) e^{i(uQ+vP)} f \, d u \, d v, \quad f \in L^2(\mathbb{R}^d).$$

Here, $e^{i(uQ+vP)}$ is understood as the Schrödinger representation

$$(1.2) \quad e^{i(uQ+vP)} f(x) := e^{(1/2)iu v} e^{iuQ} e^{ivP} f(x) = e^{(1/2)iu v + iux} f(x + v)$$

which unitarily represents the canonical commutation relations for the position and momentum operators on $L^2(\mathbb{R}^d)$; the first identity is suggested by the Baker–Campbell–Hausdorff formula, noting that all higher commutators of P and Q vanish. As shown in Proposition 1, p. 554 of [50], (1.1) encodes the standard pseudo-differential calculus, in the sense that for every $a \in \mathcal{S}(\mathbb{R}^{2d})$ there exists a unique $b \in \mathcal{S}(\mathbb{R}^{2d})$ such that

$$(1.3) \quad a(Q, P)f(x) = \frac{1}{(2\pi)^{d/2}} \int_{\mathbb{R}^d} b(x, \xi) \widehat{f}(\xi) e^{ix\xi} \, d\xi,$$

the map $a \mapsto b$ being continuous with respect to various relevant topologies. The advantage of (1.1) over (1.3) is that the former makes sense for generators of bounded groups on an arbitrary Banach spaces, whereas a representation such as (1.3) is restricted to function spaces on which an appropriate Fourier transform can be defined. We thus take (1.1) as our starting point for the development of a general theory.

We work with general Weyl pairs (see Section 3 for precise definition), i.e., two d -tuples $A = (A_1, \dots, A_d)$ and $B = (B_1, \dots, B_d)$ acting on a Banach space X such that iA_1, \dots, iA_d and iB_1, \dots, iB_d generate bounded C_0 -groups satisfying the canonical (integrated) commutation relations

$$(1.4) \quad e^{isA_j} e^{itA_k} = e^{itA_k} e^{isA_j}, \quad e^{isB_j} e^{itB_k} = e^{itB_k} e^{isB_j}, \quad e^{isA_j} e^{itB_k} = e^{-ist\delta_{jk}} e^{itB_k} e^{isA_j}.$$

In this context, (1.2) is replaced by

$$e^{i(uA+vB)} := e^{(1/2)iu v} e^{iuA} e^{ivB} := e^{(1/2)iu v} \prod_{j=1}^d e^{iu_j A_j} \prod_{k=1}^d e^{iv_k B_k}.$$

The analogue of (1.1),

$$(1.5) \quad a(A, B)f = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{a}(u, v) e^{i(uA+vB)} f \, du \, dv, \quad f \in X,$$

defines an algebra homomorphism between $\mathcal{S}(\mathbb{R}^{2d})$ endowed with the (non-commutative) Moyal product

$$(1.6) \quad (a \# b)(x, \xi) = \frac{1}{\pi^{2d}} \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} a(x+u, \xi+u') b(x+v, \xi+v') e^{-2i(vu'-uv')} \, du \, du' \, dv \, dv'$$

into the space of bounded linear operators $\mathcal{L}(X)$.

This calculus was already studied by Anderson [3], [4] (who assumed the groups to consist of isometries; but this can be achieved by a simple renorming of X). The latter paper contains a generalisation of the Moyal product (1.6) in the case where the operators $iA_1, \dots, iA_n, iB_1, \dots, iB_n$ are the generators of a strongly continuous unitary representation of an arbitrary nilpotent Lie group, equation (1.6) being the special case corresponding to the Schrödinger representation of the reduced Heisenberg group (cf. Section 4).

The Moyal product defined in (1.6) is used in pseudo-differential operator theory to deal with composition of symbols. In Section 3, we show that if the algebra homomorphism defined by (1.5) is continuous from $\mathcal{S}(\mathbb{R}^{2d})$ endowed with the topology of the standard pseudo-differential class of symbols S^0 to $\mathcal{L}(X)$, then the calculus can be meaningfully extended from $\mathcal{S}(\mathbb{R}^{2d})$ to S^0 . This is an analogue of the fundamental convergence lemma in the theory of H^∞ -functional calculus (see, e.g. Proposition 10.2.11 of [33]), and is proved using asymptotic expansions of the Moyal product, typical of pseudo-differential calculus. Having such a convergence lemma shows that a pseudo-differential calculus for (A, B) can be defined as soon as appropriate bounds on the operators defined in (1.5) are obtained.

This is the main point of our paper. While previous studies, such as Anderson's [3], [4], have focused on algebraic properties of the calculus considered as a map from $\mathcal{S}(\mathbb{R}^{2d})$ to $\mathcal{L}(X)$, we focus on the analytic question of extending the calculus to a continuous map from a large set of functions to $\mathcal{L}(X)$. This is the operator theoretic analogue of the harmonic analytic problem of identifying classes of symbols for which pseudo-differential operators are bounded on L^p .

One of the applications of pseudo-differential calculus is to study Schrödinger operators such as the harmonic oscillator defined by $\frac{1}{2}\Delta f(x) - \frac{1}{2}|x|^2 f(x)$ on $L^2(\mathbb{R}^d)$. In our abstract situation, we show that it is possible to express, in Section 5, the semigroup generated by

$$(1.7) \quad -L := \frac{1}{2}d - \frac{1}{2} \sum_{j=1}^d (A_j^2 + B_j^2)$$

in terms of the Weyl calculus as

$$(1.8) \quad e^{-tL} = a_t(A, B),$$

where $a_t \in \mathcal{S}(\mathbb{R}^{2d})$ is the function

$$a_t(x, \xi) := \left(1 + \frac{1 - e^{-t}}{1 + e^{-t}}\right)^d \exp\left(-\frac{1 - e^{-t}}{1 + e^{-t}}(|x|^2 + |\xi|^2)\right).$$

For the pair of position and momentum operators associated with the Ornstein–Uhlenbeck operator (see Example 3.4), (1.8) is a well known formula for the Ornstein–Uhlenbeck semigroup which goes back, at least, to [52]; see also [45], where this formula was rediscovered by a reduction to Mehler’s formula. Here we show, with a different proof, that it generally holds for the operators L associated with Weyl pairs through (1.7). As such, (1.8) can be thought of as an abstract analogue of Mehler’s formula for Weyl pairs.

To obtain useful bounds for various functions of L we use, in Section 6, the idea of transference to derive bounds for $a(A, B)$ acting on X from corresponding bounds on the twisted convolution with \widehat{a} , viewed as an operator acting on $L^p(\mathbb{R}^d; X)$. This idea can be traced back to Coifman and Weiss [14] and the form used here is inspired by the work of Hieber and Prüss [30], Haase [26], and Haase and Rozendaal [28]. They have shown that bounds on the Phillips functional calculus defined, for a generator iG of a bounded C_0 -group acting on a Banach space X , by

$$a(G)f = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} \widehat{a}(u)e^{iuG} f du,$$

can be obtained from bounds on convolution operators acting on $L^2(\mathbb{R}; X)$. The latter can then be proven using, for instance, Bourgain’s UMD-valued Fourier multiplier theorem [11], or its analogue for operator-valued kernels proven by Weis in [53].

For twisted convolutions, however, no UMD-valued theory is yet available. Developing such a theory is bound to be difficult, given that the (scalar-valued) L^p -theory of twisted convolutions, as developed by Mauceri in [41], is already subtle (see also [42]). For applications to spectral multipliers theorems for L , fortunately, we only need to handle highly specific twisted convolutions that can effectively be “untwisted”. This is shown in Section 7, where we prove R -sectoriality for the operator L defined by (1.7) in UMD lattices X . In Section 8, we use this result to deduce the boundedness of the H^∞ -calculus of L on UMD lattices X from the boundedness of the Weyl calculus of (A, B) . We also show that the angle of this calculus is best possible (namely 0). Going even further, we apply the recent Kriegler–Weis approach to spectral multipliers developed in [37], [38] to show that this H^∞ -calculus can in fact be extended to a Hörmander class of sufficiently smooth but not necessarily analytic functions. This is possible because the estimates obtained in Section 7 are precise enough for us to check the assumption of [38].

The present paper provides a foundation for a pseudo-differential operator theory generalised in at least three directions: Witten pseudo-differential calculus, global pseudo-differential calculus on Lie groups, and rough pseudo-differential calculus. In the Witten pseudo-differential calculus, one is interested in pairs (A, B) acting on $L^p(\mathbb{R}^d, e^{-\phi(x)} dx)$, such that, informally, the “Witten Laplacian” L is of the form $h(A, B)$ for an appropriate “Hamiltonian” h which is chosen so that the measure $e^{-\phi(x)} dx$ is an invariant measure for L . We started such a theory in [45] in the most classical case where the choice $\phi(x) = \frac{1}{2}|x|^2$ brings us back to the Gaussian setting and L reduces to the Ornstein–Uhlenbeck operator. In work in progress, some of the results proven in the present paper are applied to extend the functional calculus theory of the Ornstein–Uhlenbeck operator in [23].

From the Lie group point of view, the present paper can be seen as an approach to (sub)pseudo-differential calculus on $L^p(H)$, where H is the Heisenberg group. The prefix “sub” here indicates that we consider a pseudo-differential calculus that extends the Fourier multiplier calculus given by the functional calculus of the sub-Laplacian (removing this prefix by extending the present paper to add ∂_t to the joint functional calculus of the Weyl pair (X, Y) , in the spirit of [51], would be interesting). In the setting where X is an L^p -space, a Lie group representation approach to some of the results in Section 5 has already been pursued in [17], [18] for more general higher-order commutator relations; see also [19], [20]. Building on earlier work in [19], in the setting of L^p -spaces the boundedness of the H^∞ -calculus of $\varepsilon + L$ for $\varepsilon > 0$ has been proved in [49] by more direct transference arguments. The present operator-theoretic perspective could help construct global pseudo-differential calculi on nilpotent Lie groups. Such a theory is currently being developed by Ruzhansky, Fischer, and their collaborators (see, in particular, [22]). The theory of pseudo-differential operators on the Heisenberg group is developed in [7].

Last but not least, we aim to perturb the Weyl calculus, both from an operator theoretic and a harmonic analytic perspective, to eventually treat pairs of the form

$$Q_{B,j}f(x) = \frac{1}{2} \left(\partial_j f(x) + x_j f(x) - \sum_{k=1}^d \beta_{kj}(x) (\partial_k f(x) - x_k f(x)) \right),$$

$$P_{B,j}f(x) = \frac{1}{2i} \left(\partial_j f(x) + x_j f(x) + \sum_{k=1}^d \beta_{kj}(x) (\partial_k f(x) - x_k f(x)) \right),$$

where both the matrix $B = (\beta_{kj})_{k,j=1}^d$ and its inverse have bounded measurable coefficients. Notice that we recover the standard pair with $B = I$. These are analogues of the perturbations of Dirac operators considered in [6]. Since the latter can be interpreted as a rough Fourier multiplier theory, a corresponding theory for (Q_B, P_B) could be interpreted as a rough pseudo-differential calculus.

NOTATION AND CONVENTIONS. All vector spaces are complex unless the contrary is stated. To be in line with standard notation in pseudo-differential calculus, we reserve the notation (x, ξ) for the general point in $\mathbb{R}^{2d} = \mathbb{R}^d \times \mathbb{R}^d$. Because most applications are concerned with function spaces anyway, general elements in a Banach space X will be denoted by f, g, \dots . For $\xi \in \mathbb{R}^d$ we write $\langle \xi \rangle = (1 + |\xi|^2)^{1/2}$. Standard multi-index notation is used. We let $\mathbb{N} = \{0, 1, 2, \dots\}$.

When $A = (A_1, \dots, A_d)$ and $B = (B_1, \dots, B_d)$ are d -tuples of linear operators with domains $D(A_j)$ and $D(B_j)$ respectively, we set $D(A) = \bigcap_{j=1}^d D(A_j)$ and

$D(B) = \bigcap_{j=1}^d D(B_j)$. For $u, v \in \mathbb{R}^d$ we write $uv := \sum_{j=1}^d u_j v_j$ and define the operators uA and vB , with domains $D(A)$ and $D(B)$ respectively, by

$$uA = \sum_{j=1}^d u_j A_j, \quad vB = \sum_{j=1}^d v_j B_j.$$

We write $a \lesssim_{p_1, p_2, \dots} b$ to express that there exists a constant C , depending on the data p_1, p_2, \dots , but not on any other relevant data, such that $a \leq Cb$. If the constant is independent of all relevant data we write $a \lesssim b$.

2. PRELIMINARIES

We assume familiarity with the basic theory of pseudo-differential operators and semigroup theory. Good sources for our purposes are [1], [21] and [50]. Here we collect some terminology and results concerning UMD Banach spaces, R -boundedness, and the H^∞ -calculus of sectorial operators. Our main references are [32], [33]; other sources for these notions are, respectively, [16], [25], [27], [39], [46].

2.1. UMD SPACES. A Banach space X is said to have the UMD_p property, where $1 < p < \infty$, if there exists a finite constant $C \geq 0$ such that whenever $(m_n)_{n=1}^N$ is a finite X -valued martingale (defined on a measure space which may vary from case to case and whose length N may vary as well) and $(\epsilon_n)_{n=1}^N$ is a sequence of scalars of modulus one, we have

$$\mathbb{E} \left\| \sum_{n=1}^N \epsilon_n m_n \right\|^p \leq C^p \mathbb{E} \left\| \sum_{n=1}^N m_n \right\|^p.$$

It can be shown that if X has the UMD_p property for some $1 < p < \infty$, then it has this property for all $1 < p < \infty$. Accordingly it makes sense to call a Banach space a *UMD space* if it has the UMD_p property for some (equivalently, for all) $1 < p < \infty$.

In some treatments only scalars $\epsilon_n = \pm 1$ are used. This leads to an equivalent definition, the only difference being that the numerical value of the constant may change (see Proposition 4.2.10 of [32]).

The importance of the class of UMD spaces derives from a celebrated theorem due to Burkholder and Bourgain [10], [12] which characterises it as precisely the class of Banach spaces X for which the Hilbert transform extends to a bounded operator on $L^p(\mathbb{R}; X)$ for some (equivalently, for all) $1 < p < \infty$. This, in turn, allows one to prove the boundedness in $L^p(\mathbb{R}^d; X)$ of very general classes of singular integral operators. For some of the sharpest results presently available see [31]. In particular every Calderón–Zygmund operator with a kernel satisfying the so-called “standard estimates” is bounded on $L^p(\mathbb{R}^d; X)$ for all UMD spaces X and exponents $1 < p < \infty$.

Examples of UMD spaces include Hilbert spaces, the L^p -spaces with $1 < p < \infty$, and the Schatten classes \mathcal{C}_p with $1 < p < \infty$. The class of UMD spaces is stable under passing to equivalent norms and taking closed subspaces, quotients, and ℓ^p -direct sums. If X is UMD and $1 < p < \infty$, then also $L^p(M, \mu; X)$ is UMD, for any measure space (M, μ) . As a consequence, all “classical” function spaces used in analysis such as Sobolev spaces, Besov spaces, and Triebel–Lizorkin spaces are UMD as long as the exponents in their definitions are within the reflexive range. UMD spaces are reflexive, and therefore spaces such as $c_0, \ell^1, \ell^\infty, C(K), L^1(M, \mu), L^\infty(M, \mu)$ are not UMD (with exception of the trivial cases when the latter three are finite-dimensional).

2.2. R -BOUNDEDNESS. A *Rademacher sequence* is a sequence of independent random variables $(\epsilon_n)_{n=1}^\infty$, defined on some probability space, the values of which are uniformly distributed in the set of scalars of modulus one. Thus if the scalar field is real, Rademacher variables take values ± 1 with equal probability $\frac{1}{2}$, and if the scalar field is complex their values are uniformly distributed in the unit circle in the complex plane.

Let X and Y be Banach spaces and let $\mathcal{L}(X, Y)$ denote the space of all bounded linear operators from X into Y . A subset \mathcal{T} of $\mathcal{L}(X, Y)$ is said to be R_p -bounded, where $0 < p < \infty$, if there exists a finite constant $C \geq 0$ such that for all finite sequences $T_1, \dots, T_N \in \mathcal{T}$ and $x_1, \dots, x_N \in X$ (where N may vary) one has

$$\mathbb{E} \left\| \sum_{n=1}^N \epsilon_n T_n x_n \right\|^p \leq C^p \mathbb{E} \left\| \sum_{n=1}^N \epsilon_n x_n \right\|^p.$$

The least admissible constant C is called the R_p -bound of \mathcal{T} and is denoted by $\mathcal{R}_p(\mathcal{T})$.

By the Kahane–Khintchine inequality (see Theorem 6.2.4 of [33]), if \mathcal{T} is R_p -bounded for some $0 < p < \infty$, then it is R_p -bounded for all $0 < p < \infty$, and for all $0 < p < \infty$ we have

$$\mathcal{R}_p(\mathcal{T}) \approx_p \mathcal{R}(\mathcal{T}),$$

where by default we write $\mathcal{R}(\mathcal{T}) := \mathcal{B}_2(\mathcal{T})$. Accordingly it makes sense to call \mathcal{T} *R-bounded* if it is *R*-bounded for some (equivalently, for all) $0 < p < \infty$.

In some treatments real-valued Rademacher variables (random variables taking the values ± 1 with equal probability) are used. This leads to an equivalent definition, the only difference being that the numerical value of the *R*-bounds may change (see Proposition 6.1.9 of [33]). Upon replacing the role of Rademacher variables by Gaussian variables, one arrives at the notion of *γ -boundedness*. Every *R*-bounded set of operators is γ -bounded (by a simple randomisation argument, see Theorem 8.1.3 of [33]), and every γ -bounded set is uniformly bounded (take $N = 1$). If X has finite cotype, every γ -bounded family in $\mathcal{L}(X, Y)$ is *R*-bounded, and if X has cotype 2 and Y has type 2 (in particular, if X and Y are isomorphic to Hilbert spaces), then every uniformly bounded family in $\mathcal{L}(X, Y)$ is *R*-bounded (see Theorem 8.1.3 of [33]). The Kahane contraction principle (see Theorem 6.1.13 of [33]) implies that bounded subsets of the scalar field, viewed as bounded operators on a Banach space X through scalar multiplication, are *R*-bounded. *R*-Bounded sets enjoy many permanence properties; in particular they are closed under taking convex hulls and weak operator closure (see Sections 8.1.e, 8.4.a, 8.5.a of [33]).

The notion of *R*-boundedness originates from harmonic analysis, where it captures the essence of so-called “square function estimates”. As such it goes back to the works [8], [11]; its first systematic study is [13]. Rather than explaining this aspect in full detail (for this we refer to Chapter 8 of [33]) we mention (see Proposition 6.3.3 of [33]) that if $X = L^q(M, \mu)$ with $1 \leq q < \infty$, then for all $0 < p < \infty$ one has the equivalence of norms

$$\left(\mathbb{E} \left\| \sum_{n=1}^N \epsilon_n f_n \right\|_{L^q(M, \mu)}^p \right)^{1/p} \approx_{p,q} \left\| \left(\sum_{n=1}^N |f_n|^2 \right)^{1/2} \right\|_{L^q(M, \mu)}$$

with implied constants that depend only on p and q . Thus, in the context of L^q -spaces, *R*-boundedness reduces to a square function estimate.

2.3. *H[∞]-CALCULUS*. Let X be a Banach space and let $0 < \sigma < \pi$. A closed operator $L : D(L) \supseteq X \rightarrow X$ (with $D(L)$ the *domain* of L) is said to be *σ -sectorial* if its spectrum is contained in the closure of the sector

$$\Sigma_\sigma = \{z \in \mathbb{C} : z \neq 0, |\arg z| < \sigma\}$$

(arguments are taken in $(-\pi, -\pi)$) and satisfies

$$(2.1) \quad \|R(z, L)\| \leq \frac{M}{|z|}$$

on the complement of $\bar{\Sigma}_\sigma$, for some finite constant $M \geq 0$. Here, $R(z, L) := (z - L)^{-1}$ is the resolvent operator. An operator is said to be *sectorial* if it is σ -sectorial for some $0 < \sigma < \pi$. The number

$$\omega(L) := \inf\{\sigma \in (0, \pi) : L \text{ is } \sigma\text{-sectorial}\}$$

is called the *angle of sectoriality of L*.

For $0 < \theta < \pi$ let $H^1(\Sigma_\theta)$ be the Banach space of all holomorphic functions $\phi : \Sigma_\theta \rightarrow \mathbb{C}$ satisfying

$$\|\phi\|_{H^1(\Sigma_\theta)} := \sup_{0 < \nu < \theta} \frac{1}{2\pi} \int_{\partial\Sigma_\nu} |\phi(z)| \frac{|dz|}{|z|} < \infty.$$

If L is σ -sectorial, then for any $\phi \in H^1(\Sigma_\theta)$ with $\sigma < \theta < \pi$ we may define

$$(2.2) \quad \phi(L) := \frac{1}{2\pi i} \int_{\partial\Sigma_\nu} \phi(z)R(z, L) dz,$$

taking $\sigma < \nu < \theta$ with the understanding that $\partial\Sigma_\nu$ is downwards oriented. This integral converges absolutely and defines a bounded operator of norm at most $M\|\phi\|_{H^1(\Sigma_\theta)}$, where M is the constant of (2.1). It is a consequence of Cauchy’s theorem and Proposition H.2.5 of [33] that the definition of $\phi(L)$ is independent of the choice of the angle ν .

If we were to replace the role of $H^1(\Sigma_\theta)$ by the space $H^\infty(\Sigma_\theta)$ of all bounded holomorphic functions on Σ_θ , we would run into the difficulty that the corresponding Dunford integral in (2.2) becomes singular at both the origin and at infinity. To handle this situation a sectorial operator L is said to have a *bounded $H^\infty(\Sigma_\sigma)$ -calculus*, where $\omega(L) < \sigma < \pi$, if there exists a finite constant $K \geq 0$ such that

$$\|\phi(L)\| \leq K\|\phi\|_{H^\infty(\Sigma_\sigma)}$$

for all $\phi \in H^1(\Sigma_\sigma) \cap H^\infty(\Sigma_\sigma)$. A sectorial operator L is said to have a *bounded H^∞ -calculus* if it has a bounded $H^\infty(\Sigma_\sigma)$ -calculus for some $\omega(L) < \sigma < \pi$. The number

$$\omega_{H^\infty}(L) := \inf\{\sigma \in (\omega(L), \pi) : L \text{ has a bounded } H^\infty(\Sigma_\sigma)\text{-calculus}\}$$

is called the *angle of the H^∞ -calculus of L*.

If L is densely defined, has dense range and a bounded $H^\infty(\Sigma_\sigma)$ -calculus, the *McIntosh convergence lemma* [43] (see also Theorem 10.2.13 of [33]) allows one to uniquely define, for every $\phi \in H^\infty(\Sigma_\sigma)$, a bounded operator $\phi(L)$ by

$$\phi(L)f := \lim_{n \rightarrow \infty} \phi_n(L)f, \quad f \in X,$$

where $(\phi_n)_{n \geq 1}$ is any sequence in $H^1(\Sigma_\sigma) \cap H^\infty(\Sigma_\sigma)$ that is uniformly bounded and converges to ϕ pointwise on Σ_σ .

The prime example of a sectorial operator with a bounded H^∞ -calculus (of angle 0) is the negative Laplacian $L = -\Delta$ on $L^p(\mathbb{R}^d; X)$ for any UMD space X and $1 < p < \infty$. More generally, under minor regularity assumptions on the coefficients, uniformly elliptic operators on sufficiently regular domains D in \mathbb{R}^d have bounded H^∞ -calculi on $L^p(D; X)$ under various boundary conditions. Examples with precise formulations are reviewed in [16], [33], [39].

There is an interesting interplay between R -boundedness and H^∞ -calculus. Let us say that a closed operator L is σ - R -sectorial if $\sigma(L)$ is contained in $\overline{\Sigma}_\sigma$ and the set

$$\{zR(z, L) : z \in \mathbb{C}\overline{\Sigma}_\sigma\}$$

is R -bounded. Since R -boundedness implies boundedness, every σ - R -sectorial is σ -sectorial. The operator L is said to be R -sectorial if it is σ - R -sectorial for some $0 < \sigma < \pi$. The infimum

$$\omega_R(L) := \inf\{\sigma \in (\omega(L), \pi) : L \text{ is } \sigma\text{-}R\text{-sectorial}\}$$

is called the *angle of R -sectoriality* of L . It was shown by Kalton and Weis [35] (see also Corollary 10.4.10 of [33]) that if L is a sectorial operator with a bounded H^∞ -calculus on a UMD Banach space X (actually a slightly weaker assumption will do for this purpose, but this is not relevant to us here), then L is R -sectorial and we have

$$(2.3) \quad \omega_R(L) = \omega_{H^\infty}(L).$$

In this context it is interesting to observe that for R -sectorial operators L it may happen that $\omega_R(L) > \omega(L)$; see [34].

3. WEYL PAIRS

Let $A = (A_1, \dots, A_d)$ and $B = (B_1, \dots, B_d)$ be two d -tuples of closed and densely defined operators acting in a complex Banach space X . We assume that each of the operators iA_j and iB_j generates a uniformly bounded C_0 -group on X . We denote these groups by $(e^{itA_j})_{t \in \mathbb{R}}$ and $(e^{itB_j})_{t \in \mathbb{R}}$, respectively.

DEFINITION 3.1. Under the above assumptions, the pair (A, B) is called a *Weyl pair* of dimension d if the (integrated) *canonical commutation relations* hold for all $s, t \in \mathbb{R}$ and $1 \leq j, k \leq d$:

$$(3.1) \quad e^{isA_j} e^{itA_k} = e^{itA_k} e^{isA_j}, \quad e^{isB_j} e^{itB_k} = e^{itB_k} e^{isB_j}, \quad e^{isA_j} e^{itB_k} = e^{-ist\delta_{jk}} e^{itB_k} e^{isA_j},$$

where δ_{jk} is the usual Kronecker symbol.

REMARK 3.2. In essence, Weyl pairs are bounded strongly continuous representations on X of the reduced Heisenberg group. This point of view will not play a role in this paper, for a reason that is best illustrated by considering the particular case of the Schrödinger representation corresponding to the classical Weyl pair of position and momentum operators. The two most studied (and unitarily equivalent) realisations of this representation, corresponding respectively to the harmonic oscillator (where one takes $X = L^2(\mathbb{R}^d)$) and the Ornstein–Uhlenbeck operator (where one takes $X = L^2(\mathbb{R}^d, \gamma_d)$ with γ_d the standard Gaussian distribution in d dimensions) become very different once we depart from their Hilbertian settings and consider their extensions to $L^p(\mathbb{R}^d)$ and $L^p(\mathbb{R}^d, \gamma_d)$, respectively.

This is elaborated in the Examples 3.3 and 3.4 below. It shows that the underlying algebraic structure of the Heisenberg group has only a limited bearing on the properties of general Weyl pairs.

Being a Weyl pair is an isomorphic notion, in that it is insensitive to changing to an equivalent norm. In particular, by switching to the equivalent norm

$$\|x\| := \max_{j=1,\dots,d} \max \left\{ \sup_{t \in \mathbb{R}} \|e^{itA_j}x\|, \sup_{t \in \mathbb{R}} \|e^{itB_j}x\| \right\}$$

we could have equivalently assumed that the operators iA_j and iB_j generate C_0 -groups of isometries on X . If (A, B) is a Weyl pair on X and $T : X \rightarrow Y$ is an isomorphism of Banach spaces, then (TAT^{-1}, TBT^{-1}) is a Weyl pair on Y . This is of course trivial, but it is of some interest in connection with the next example, for on Hilbert spaces it easily provides examples of non-selfadjoint Weyl pairs.

EXAMPLE 3.3 (Standard position/momentum pair). On $L^p(\mathbb{R}^d)$, $1 \leq p < \infty$, the position and momentum operators Q_j and P_j , $1 \leq j \leq d$, are defined by

$$Q_j f(x) = x_j f(x), \quad P_j f(x) = \frac{1}{i} \partial_j f(x), \quad x \in \mathbb{R}^d.$$

With their natural domains, it is easily checked that they define a Weyl pair (Q, P) . Indeed, iQ_j generates the multiplication group on $L^p(\mathbb{R}^d)$ given by

$$e^{itQ_j} g(x) = e^{itx_j} g(x), \quad x \in \mathbb{R}^d, t \in \mathbb{R},$$

and iP_j generates the translation group on $L^p(\mathbb{R}^d)$ given by

$$e^{itP_j} g(x) = g(x + te_j), \quad x \in \mathbb{R}^d, t \in \mathbb{R},$$

with e_j the j -th unit vector of \mathbb{R}^d . The commutation relations are easily checked.

The position/momentum pair is sometimes referred to as the *standard pair* and provides the main example of a Weyl pair. A well-known uniqueness result of Stone and von Neumann (see, e.g., Chapter 14 of [29] or Section 4.3 of [48]) asserts that every Weyl pair of dimension d of self-adjoint operators in a Hilbert space is unitarily equivalent to a direct sum of copies of standard pairs on $L^2(\mathbb{R}^d)$.

EXAMPLE 3.4 (Gaussian position/momentum pair). Let us denote by γ the standard Gaussian measure on \mathbb{R}^d . On $L^p(\mathbb{R}^d, \gamma)$, $1 \leq p < \infty$, we consider the position and momentum pair (Q^γ, P^γ) given by $Q^\gamma = (Q_1^\gamma, \dots, Q_d^\gamma)$ and $P^\gamma = (P_1^\gamma, \dots, P_d^\gamma)$ defined by

$$Q_j^\gamma := \frac{1}{\sqrt{2}}(a_j + a_j^\dagger), \quad P_j^\gamma := \frac{1}{i\sqrt{2}}(a_j - a_j^\dagger),$$

where the annihilation and creation operators a_j and a_j^\dagger are defined by

$$a_j = \partial_j, \quad a_j^\dagger = -\partial_j + x_j.$$

Thus, for $f \in C_c^1(\mathbb{R}^d)$,

$$Q_j^\gamma f(x) = \frac{1}{\sqrt{2}} x_j f(x), \quad P_j^\gamma f(x) = \frac{1}{i\sqrt{2}} (2\partial_j - x_j) f(x).$$

It is readily verified that the pair (Q^γ, P^γ) satisfies the canonical commutation relations. As we will explain in a moment, for $p = 2$ this pair is unitarily equivalent to the standard pair.

It is clear that the operators iQ_j^γ generate C_0 -contraction groups of multiplication operators on $L^p(\mathbb{R}^d, \gamma)$ for all $1 \leq p < \infty$. On the other hand, the operators iP_j^γ generate bounded C_0 -groups on $L^p(\mathbb{R}^d, \gamma)$ if and only if $p = 2$. Thus (Q^γ, P^γ) is a Weyl pair on $L^p(\mathbb{R}^d, \gamma)$ if and only if $p = 2$. This can be deduced from Theorem 5.2 below as follows. By a result of [45], in $L^2(\mathbb{R}^d, \gamma)$ the operator $\frac{1}{2}((Q^\gamma)^2 + (P^\gamma)^2) - \frac{1}{2}d$ considered in Theorem 5.2 is the Ornstein–Uhlenbeck operator. If (Q^γ, P^γ) were to be a Weyl pair in $L^p(\mathbb{R}^d, \gamma)$ for certain $p \in (1, \infty) \setminus \{2\}$, the theorem would imply that the Ornstein–Uhlenbeck semigroup extends holomorphically to the right half-plane $\{\Re z > 0\}$, and this is well known to be false. In fact the optimal angle θ_p of holomorphy for the Ornstein–Uhlenbeck semigroup on $L^p(\mathbb{R}^d, \gamma)$ is known to be $\cos \theta_p = \frac{|p-2|}{2\sqrt{p-1}}$ (see [23]).

The failure of iP_j^γ to generate a bounded C_0 -group on $L^p(\mathbb{R}^d, \gamma)$ for $p \neq 2$ can also be easily checked by hand. Let $m(dx) = \frac{dx}{(2\pi)^{d/2}}$ denote the normalised Lebesgue measure on \mathbb{R}^d . On $L^2(\mathbb{R}^d, \gamma)$ the group generated by iP_j^γ is given by $e^{itP_j^\gamma} = U^{-1}T_j(t)U$, where $T_j(t)$ is the translation group on $L^2(\mathbb{R}^d, m)$ in the j -th direction and $U : L^2(\mathbb{R}^d, \gamma) \rightarrow L^2(\mathbb{R}^d, m)$ is the unitary mapping given by $U = \delta \circ E$ with

$$Ef(x) = e^{-(1/4)|x|^2} f(x), \quad \delta f(x) := (\sqrt{2})^d f(\sqrt{2}x).$$

An easy computation shows that, in $L^2(\mathbb{R}^d, \gamma)$, the operators $e^{itP_j^\gamma}$ are given by

$$e^{itP_j^\gamma} f(x) = e^{(1/4)|x|^2 - (1/2)(x/\sqrt{2}-t)^2} f(x + t\sqrt{2}).$$

Then, after an integration and change of variable,

$$\|e^{itP_j^\gamma} f\|_p^p = \frac{1}{(2\pi)^{d/2}} \int_{\mathbb{R}^d} e^{((1/2)-(p/4))(2\sqrt{2}xt-2t^2)} |f(x)|^p d\gamma(x).$$

For $p \in [1, 2)$ it follows that $e^{itP_j^\gamma}$ fails to extend to a bounded operator acting on $L^p(\mathbb{R}^d, \gamma)$ for all $t > 0$, and for $p \in (2, \infty)$ the operators $e^{itP_j^\gamma}$ are bounded on $L^p(\mathbb{R}^d, \gamma)$, but not uniformly bounded as a function of $t > 0$.

EXAMPLE 3.5 (Modified Gaussian position/momentum pair). It is of some interest to note that the pair (Q^γ, P^γ) of the previous example does form a Weyl pair on $L^p(\mathbb{R}^d, \gamma_{2/p})$ for all $p \in [1, 2]$, where $\gamma_\tau(dx) = (2\pi\tau)^{-d/2} e^{-|x|^2/2\tau} dx$.

This is simply because with this scaling of the measure the mapping U considered above defines an isomorphism from $L^p(\mathbb{R}^d, \gamma_{2/p})$ onto $L^p(\mathbb{R}^d, m)$. Then each iP_j^γ generates a bounded C_0 -group on $L^p(\mathbb{R}^d, \gamma_{2/p})$ which, under U , conjugates with the translation group in the j -th direction on $L^p(\mathbb{R}^d, m)$.

EXAMPLE 3.6 (Duality). If (A, B) is a Weyl pair in X , then the pair of adjoint operators (B^*, A^*) is a Weyl pair in X^* provided the operators A_j^* and B_j^* are densely defined (by a classical result in semigroup theory (see Proposition I.5.14 of [21]) this is always the case if X is reflexive).

EXAMPLE 3.7 (Additive commuting perturbations). If (A, B) is a Weyl pair and C is a bounded operator resolvent commuting with A , then $(A, B + C)$ is a Weyl pair whenever the group generated by $i(B + C)$ is bounded. Indeed, the assumption implies that C commutes with the operators e^{itA} , and the commutation relations (3.1) follow from this by going through the standard proof of the variation of constants formula for perturbed (semi)groups using Picard iteration. The simplest example is obtained by taking $C = \omega I$ with $\omega \in \mathbb{R}$. This amounts to frequency modulating the group generated by iB . More generally one could take C to be any densely defined closed operator such that iC generates a bounded group commuting with the group generated by iA .

Similarly, if (A, B) is a Weyl pair and C is a bounded operator commuting with the resolvent of B , then $(A + C, B)$ is a Weyl pair whenever the group generated by $i(A + C)$ is bounded.

EXAMPLE 3.8 (Skew transforms). If (A, B) is a Weyl pair, then for every $\lambda \in \mathbb{R}$ the pair $(A, \lambda A + B)$ is a Weyl pair. Some care has to be taken with the interpretation of $\lambda A + B$; we interpret it as the generator of the C_0 -group given by

$$e^{it(\lambda A + B)} := e^{(1/2)i\lambda t^2} e^{i\lambda t A} e^{itB}$$

(this idea will be further developed in a moment). Similarly, if (A, B) is a Weyl pair, then for every $\lambda \in \mathbb{R}$ the pair $(A + \lambda B, B)$ is a Weyl pair.

EXAMPLE 3.9. Let $((Q_1, Q_2), (P_1, P_2))$ be the standard pair of dimension $2d$ on $L^2(\mathbb{R}^{2d})$, i.e.,

$$\begin{aligned} Q_{1,j}f(x, \xi) &= x_j f(x, \xi), & Q_{2,j}f(x, \xi) &= \xi_j f(x, \xi), \\ P_{1,j}f(x, \xi) &= \frac{1}{i} \frac{\partial f}{\partial x_j}(x, \xi), & P_{2,j}f(x, \xi) &= \frac{1}{i} \frac{\partial f}{\partial \xi_j}(x, \xi), \end{aligned}$$

for $1 \leq j \leq d$. Reasoning as in the preceding examples, we see that $(-\frac{1}{2}Q_2 - P_1, \frac{1}{2}Q_1 - P_2)$ is a Weyl pair of dimension d on $L^2(\mathbb{R}^{2d})$. As we show in Lemma 6.4, the Weyl calculus of this pair encodes twisted convolutions. Many variations on twisted convolutions can be considered through the Weyl calculus of twisted standard pairs obtained from different twists than the one above.

EXAMPLE 3.10 (Quantum variables). In [24], González-Pérez, Junge, and Parcet introduce a (non-commutative) Fourier transform, as well as position and momentum operators, associated with certain von Neumann algebras called quantum euclidean spaces (or Moyal deformations, or CCR algebras). Their construction allows them to define non-commutative analogues of the key notions of Calderón–Zygmund theory, including off-diagonal kernel estimates and Hörmander symbol classes, and then to prove analogues of the main theorems in singular integral operator theory. We cannot describe their construction in detail here, but note that their quantum variables $(x_{\Theta,j})_{j=1,\dots,2d}$ are Weyl pairs (for the appropriate choice of Θ) acting on some non-commutative L^p -spaces (see Proposition 1.9 of [24]).

We now collect some easy properties of Weyl pairs which will be useful later on. For $d = 1$ they are due to Kato [36] (see also Section 4.9 of [48]) and the proofs given there extend without difficulty to the present case. The main observation is that, upon taking Laplace transforms, the third commutation relation in (3.1) implies the identities

$$(3.2) \quad R(\lambda, iA_j)e^{itB_j} = e^{itB_j}R(\lambda + it, iA_j), \quad R(\lambda, iB_j)e^{itA_j} = e^{itA_j}R(\lambda - it, iB_j),$$

for all $t \in \mathbb{R}$, $\Re\lambda \neq 0$, and $1 \leq j \leq d$. It follows that e^{itB_j} leaves $D(A_j)$ invariant, e^{itA_j} leaves $D(B_j)$ invariant, and

$$(3.3) \quad A_j e^{itB_j} f = e^{itB_j} (A_j - t) f \quad f \in D(A_j), \quad B_j e^{itA_j} f = e^{itA_j} (B_j + t) f \quad f \in D(B_j).$$

The same argument applies to the remaining combinations of A_j and B_k , but no shifts over $\pm t$ occur when the operators commute. Thus we obtain the following lemma.

LEMMA 3.11. *Let (A, B) be a Weyl pair. The operators e^{itA_j} and e^{itB_j} leave both $D(A) := \bigcap_{k=1}^d D(A_k)$ and $D(B) := \bigcap_{k=1}^d D(B_k)$ invariant. For $j \neq k$ we have*

$$(3.4) \quad A_j e^{itB_k} f = e^{itB_k} A_j f \quad f \in D(A_j), \quad B_k e^{itA_j} f = e^{itA_j} B_k f \quad f \in D(B_k),$$

while for $j = k$ the identities (3.3) hold.

Differentiating (3.2) at $t = 0$ gives

$$(3.5) \quad \begin{aligned} R(\lambda, iB_j)R(\mu, iA_j) &= R(\mu, iA_j)R(\lambda, iB_j)[I - iR(\lambda, iB_j)R(\mu, iA_j)], \\ R(\lambda, iA_j)R(\mu, iB_j) &= R(\mu, iB_j)R(\lambda, iA_j)[I + iR(\lambda, iA_j)R(\mu, iB_j)]. \end{aligned}$$

If

$$g = \prod_{j,j'=1}^d R(\lambda_j, iA_j)R(\lambda_{j'}, iA_{j'}) \prod_{k,k'=1}^d R(\mu_k, iB_k)R(\mu_{k'}, iB_{k'}) f$$

with $f \in X$, then (3.5) may be used to rewrite g , for any pair $1 \leq j, k \leq d$, as

$$g = R(\lambda, iA_j)R(\mu, iB_k)C_{jk}f = R(\mu, iB_k)R(\lambda, iA_j)D_{jk}f = R(\mu, iB_j)R(\lambda, iB_k)E_{jk}f$$

for suitable bounded operators C_{jk}, D_{jk}, E_{jk} . From this we see that g belongs to $\bigcap_{1 \leq j, k \leq d} (D(A_j A_k) \cap D(A_j B_k) \cap D(B_k A_j)) \cap D(B_j B_k)$. Since, for all $f \in X$,

$$\lim_{\lambda, \mu \rightarrow \infty} \prod_{j, j', k, k'=1}^d \lambda_j \lambda_{j'} R(\lambda_j, iA_j) R(\lambda_{j'}, iA_{j'}) \mu_k \mu_{k'} R(\mu_k, iB_k) R(\mu_{k'}, iB_{k'}) f = f,$$

the limit being taken in any order for $\lambda_1, \dots, \lambda_d, \lambda'_1, \dots, \lambda'_d, \mu_1, \dots, \mu_d, \mu'_1, \dots, \mu'_d$ tending to ∞ , this subspace is dense in X . The identity (3.5) also gives the identity $A_j B_j g - B_j A_j g = ig$ for g of the above form. The same argument gives commutation for the remaining combinations of A_j and B_k . Thus we obtain the following lemma.

LEMMA 3.12. *Let (A, B) be a Weyl pair of dimension d . The subspace*

$$\bigcap_{1 \leq j, k \leq d} (D(A_j A_k) \cap D(A_j B_k) \cap D(B_k A_j) \cap D(B_j B_k))$$

is dense, and on this subspace we have $A_j A_k = A_k A_j, B_j B_k = B_k B_j$, and $A_j B_k - B_k A_j = \delta_{jk} iI$.

Let (A, B) be a Weyl pair of dimension d . Consider, for $t \in \mathbb{R}$ and $u, v \in \mathbb{R}^d$, the bounded operators

$$T_{u,v}(t) := e^{(1/2)it^2 uv} e^{ituA} e^{itvB} = e^{-(1/2)it^2 uv} e^{itvB} e^{ituA}.$$

PROPOSITION 3.13. *The family $(T_{u,v})_{t \in \mathbb{R}}$ is a bounded C_0 -group on $X, D(A) \cap D(B)$ is a core for its generator $G_{u,v}$, and, on this core, the generator is given by*

$$G_{u,v} f = iuA f + ivB f, \quad f \in D(A) \cap D(B).$$

Proof. The identity $T_{u,v}(0) = I$ is trivial. The group property $T_{u,v}(t_0) \circ T_{u,v}(t_1) = T_{u,v}(t_0 + t_1)$ follows straightforwardly from the commutation relations (3.1). Strong continuity is also clear.

It follows from the general properties of Weyl pairs mentioned earlier that each operator $T_{u,v}(t)$ maps the subspace $D(A) \cap D(B)$ into itself.

Moreover, $D(A) \cap D(B)$ is dense in X . By Lemma 3.14 below, every $f \in D(A) \cap D(B)$ belongs to $D(G_{u,v})$ and differentiation gives

$$G_{u,v} f = \frac{d}{dt} \Big|_{t=0} T_{u,v}(t) f = iuA f + ivB f, \quad f \in D(A) \cap D(B).$$

A general result in semigroup theory (see, e.g., Proposition II.1.7 of [21]) now implies that $D(A) \cap D(B)$ is a core for $G_{u,v}$. ■

The proof of Proposition 3.13 is completed by the following observation, which we leave as an easy exercise to the reader.

LEMMA 3.14. *Let $(S(t))_{t \in \mathbb{R}}$ and $(T(t))_{t \in \mathbb{R}}$ be strongly continuous families of operators, and let $f \in X$ be fixed. If*

$$(i) \ t \mapsto S(t)f \text{ is differentiable at } t = 0, \text{ with derivative } S'(0)f := \frac{d}{dt} \Big|_{t=0} S(t)f,$$

(ii) $t \mapsto T(t)S(0)f$ is differentiable at $t = 0$, with derivative

$$T'(0)S(0)f := \left. \frac{d}{dt} \right|_{t=0} T(t)S(0)f,$$

then $t \mapsto T(t)S(t)f$ is differentiable at $t = 0$, with derivative

$$\left. \frac{d}{dt} \right|_{t=0} T(t)S(t)f = T'(0)S(0)f + T(0)S'(0)f.$$

4. THE WEYL CALCULUS

Let (A, B) be a Weyl pair of dimension d on a Banach space X . For $(x, \zeta) \in \mathbb{R}^{2d}$ we consider the bounded operators

$$(4.1) \quad e^{i(uA+vB)} := e^{(1/2)iu v} e^{iuA} e^{ivB}.$$

This notation is justified by Proposition 3.13.

EXAMPLE 4.1. For the standard pair (Q, P) on $L^2(\mathbb{R}^d)$, (4.1) reduces to the Schrödinger representation: the operators $e^{i(uQ+vP)}$ are unitary on $L^2(\mathbb{R}^d)$ and given by

$$e^{i(uQ+vP)} f(x) = e^{(1/2)iu v + iux} f(x + v).$$

DEFINITION 4.2 (Weyl calculus). Let (A, B) be a Weyl pair of dimension d . For functions $a \in \mathcal{S}(\mathbb{R}^{2d})$ we define

$$a(A, B)f := \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{a}(u, v) e^{i(uA+vB)} f \, du \, dv, \quad f \in X,$$

where

$$\widehat{a}(u, v) = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} a(x, \zeta) e^{-i(xu+\zeta v)} \, dx \, d\zeta$$

is the Fourier–Plancherel transform of a . The mapping $a \mapsto a(A, B)$ from $\mathcal{S}(\mathbb{R}^{2d})$ to $\mathcal{L}(X)$ is called the Weyl calculus of (A, B) .

An easy computation based on the identity

$$(4.2) \quad e^{i(uA+vB)} \circ e^{i(u'A+v'B)} = e^{(1/2)i(u'v-uv')} e^{i(u+u')A+(v+v')B},$$

which follows from the commutation relations (3.1), gives the following analogue of the multiplicativity property of the functional calculus of a single operator: for all $a, b \in \mathcal{S}(\mathbb{R}^{2d})$ we have

$$a(A, B) \circ b(A, B) = (a \# b)(A, B),$$

where $a \# b$ is the Moyal product of a and b , given by (see Section XII.3.3 of [50])

$$(a \# b)(x, \zeta) = \frac{1}{\pi^{2d}} \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} a(x + u, \zeta + u') b(x + v, \zeta + v') e^{-2i(vu' - uv')} \, du \, du' \, dv \, dv'.$$

DEFINITION 4.3. Let $N, m \in \mathbb{N}$. A Weyl pair (A, B) is said to admit a *bounded Weyl calculus of type $(-N, m)$* if, for all $a \in \mathcal{S}(\mathbb{R}^{2d})$, we have

$$\|a(A, B)\| \lesssim \max_{|\alpha|, |\beta| \leq m} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{N+|\alpha|} |\partial_\xi^\alpha \partial_x^\beta a(x, \xi)|,$$

with a constant independent of a . The pair (A, B) is said to admit a *bounded Weyl calculus of type $-N$* if it admits a bounded Weyl calculus of type $(-N, m)$ for some $m \in \mathbb{N}$.

In Subsection 4.1 we will prove that if X is a UMD space and $1 < p < \infty$, then the standard pair (Q, P) has a bounded Weyl calculus of type 0 on $L^p(\mathbb{R}^d; X)$.

The convergence lemma for the Dunford calculus for sectorial operators (see, e.g., Theorem 10.2.2 of [33]) has the following analogue for the Weyl calculus.

LEMMA 4.4 (Convergence lemma). *Let $(a_n)_{n \in \mathbb{N}}$ be a sequence of Schwartz functions defined on \mathbb{R}^{2d} and let $N \in \mathbb{N}$. There exist $m = m(d, N) \in \mathbb{N}$ and $M = M(d, N) \in \mathbb{N}$, both depending only on d and N , such that the following holds: if (A, B) is a Weyl pair with a bounded Weyl calculus of type $(-N - 1, m)$, and if*

(i) *for all multi-indices $\gamma \in \mathbb{N}^d$ with $|\gamma| \leq M$ we have $\lim_{n \rightarrow \infty} \partial^\gamma a_n = 0$ uniformly on compact sets,*

(ii) $\sup_{n \in \mathbb{N}} \|a_n(A, B)\| < \infty,$

then $\lim_{n \rightarrow \infty} a_n(A, B)f = 0$ for all $f \in X$.

Admittedly the formulation of this lemma is a bit awkward; the point here is that we need (A, B) to be of type $(-N - 1, m)$ for all $m \geq m_0$, where m_0 may depend on N and d . The proof of the lemma is based on an asymptotic expansion representation for Moyal products of Schwartz functions.

LEMMA 4.5. *There exists a sequence $(c_\alpha)_{\alpha \in \mathbb{N}^{2d}}$ of complex numbers such that, for all $a, b \in \mathcal{S}(\mathbb{R}^{2d})$ and all integers $M \in \mathbb{N}$, there exists a function $r_{a,b;M+1} \in \mathcal{S}(\mathbb{R}^{2d})$ such that*

$$a(A, B)b(A, B) = \sum_{\alpha \in \mathbb{N}^{2d}, |\alpha|_\infty \leq M} c_\alpha \partial^\alpha (ab)(A, B) + r_{a,b;M+1}(A, B)$$

whenever (A, B) is a Weyl pair. Moreover, there exists an $m \in \mathbb{N}$, depending only on d and M , such that if (A, B) has a bounded Weyl calculus of type $(-M - 1, m)$, then

$$\begin{aligned} &\|r_{a,b;M+1}(A, B)\| \\ &\lesssim \max_{\alpha', \beta', \alpha'', \beta'' \in \mathbb{N}^d, |\alpha'|, |\beta'|, |\alpha''|, |\beta''| \leq m} \sup_{x, \xi \in \mathbb{R}^{2d}} \langle \xi \rangle^{\min(|\alpha'|, |\alpha''|)} |\partial_\xi^{\alpha'} \partial_x^{\beta'} a(x, \xi) \partial_\xi^{\alpha''} \partial_x^{\beta''} b(x, \xi)|. \end{aligned}$$

Proof. Let $a, b \in \mathcal{S}(\mathbb{R}^{2d})$. Recall that $a(A, B)b(A, B) = (a\#b)(A, B)$, where $a\#b$ is the Moyal product of a and b . By Theorem 3.16 of [1], for any $M \geq 0$ there

exists a function $r_{a,b;M+1} \in \mathcal{S}(\mathbb{R}^{2d})$ such that

$$(4.3) \quad a\#b(x, \xi) = \sum_{\alpha \in \mathbb{N}^d, |\alpha| \leq M} \frac{1}{\alpha!} \frac{1}{i^\alpha} \partial_\xi^\alpha a(x, \xi) \partial_x^\alpha b(x, \xi) + r_{a,b;M+1}(x, \xi).$$

This gives the formula in the first part of the theorem (with many coefficients c_α equal to 0).

Suppose next that (A, B) has a bounded Weyl calculus of type $(-M - 1, m)$ for some $M \in \mathbb{N}$, where $m \in \mathbb{N}$ is arbitrary for the moment but will be fixed later. Then, by assumption, the remainder $r_{a,b;M+1}(A, B)$ in the expansion (4.3) for this particular value of M satisfies the estimate

$$\|r_{a,b;M+1}(A, B)\| \lesssim \max_{|\gamma|, |\delta| \leq m} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{M+1+|\gamma|} |\partial_\xi^\gamma \partial_x^\delta r_{a,b;M+1}(x, \xi)|$$

with a constant only depending on M, m and the pair (A, B) . By Theorem 3.15 of [1], $r_{a,b;M+1}(x, \xi)$ is given by a finite linear combination, extending over all multi-indices satisfying $|\alpha| = M + 1$, of terms of the form

$$R_{\alpha,a,b}(x, \xi) := \int_{\mathbb{R}^{2d}} e^{-ix' \xi'} (\xi')^\alpha \int_0^1 \partial_\xi^\alpha p(x, \xi + \theta \xi', x + x', \xi) (1 - \theta)^M d\theta dx' d\xi'$$

for $p(x, \xi, x', \xi') = a(x, \xi) b(x', \xi')$.

As in the proof of Theorem 3.15 in [1] (see, in particular, (3.20) on page 54 and (3.10) on page 47), there exists $m(d, M) \in \mathbb{N}$, depending only on d and M , such that for all multi-indices γ, δ satisfying $|\gamma|, |\delta| \leq m(d, M)$ we have

$$|\partial_\xi^\gamma \partial_x^\delta R_{\alpha,a,b}(x, \xi)| \lesssim \langle \xi \rangle^{-(|\alpha|+|\gamma|)} = \langle \xi \rangle^{-(M+1+|\gamma|)},$$

with constant depending linearly on

$$\max_{|\alpha'|, |\beta'|, |\alpha''|, |\beta''| \leq m(d, M)} \sup_{x, \xi \in \mathbb{R}^{2d}} \langle \xi \rangle^{\min(|\alpha'|, |\alpha''|)} |\partial_\xi^{\alpha'} \partial_x^{\beta'} a(x, \xi) \partial_\xi^{\alpha''} \partial_x^{\beta''} b(x, \xi)|.$$

If we fix the integer m to be this $m(d, M)$, the second part of the lemma follows by collecting estimates. ■

The proof of the convergence lemma requires one further auxiliary result. Given a function $\eta : \mathbb{R}^{2d} \rightarrow \mathbb{C}$ and a real number $\delta > 0$ we set $\eta_\delta(x, \xi) := \eta(\delta x, \delta \xi)$.

LEMMA 4.6. *For all $\eta \in C_c^\infty(\mathbb{R}^{2d})$ with $\eta(0, 0) = 1$, and $f \in X$, we have*

$$\lim_{k \rightarrow \infty} \eta_{1/k}(A, B) f = f.$$

Proof. For all $f \in X$ we have

$$\eta_{1/k}(A, B) f = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{\eta}_{1/k}(u, v) e^{i(uA+vB)} f \, dudv$$

$$\begin{aligned}
 &= \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} k^{2d} \widehat{\eta}(ku, kv) e^{i(uA+vB)} f \, du \, dv \\
 &= \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{\eta}(u, v) e^{i((u/k)A+(v/k)B)} f \, du \, dv \xrightarrow{k \rightarrow \infty} \eta(0,0) f = f. \quad \blacksquare
 \end{aligned}$$

Proof of Lemma 4.4. Fix $N \in \mathbb{N}$, let $m = m(d, N)$ be as in Lemma 4.5 (where we take $M = N$), and suppose (A, B) has a bounded Weyl calculus of type $(-N, m)$. Let $(a_n)_{n \geq 1}$ be a sequence of Schwartz functions satisfying the assumptions (i) and (ii) in the statement of the lemma. Let $\eta \in C_c^\infty(\mathbb{R}^{2d})$ be supported in $B(0, 2)$ and identically 1 on $B(0, 1)$. Fixing $f \in X$ and $\varepsilon > 0$, by Lemma 4.6 and the uniform boundedness of the operators $a_n(A, B)$ we may choose a large enough integer k so that

$$(4.4) \quad \limsup_{n \rightarrow \infty} \|a_n(A, B)f\| \leq \limsup_{n \rightarrow \infty} \|a_n(A, B)\eta_{1/k}(A, B)f\| + \varepsilon.$$

Fix $n \geq 1$ for the moment. By Lemma 4.5,

$$\begin{aligned}
 &\|a_n(A, B)\eta_{1/k}(A, B)\| \\
 &\lesssim \left\| \sum_{|\alpha|_\infty \leq N} c_\alpha \partial^\alpha (a_n \eta_{1/k})(A, B) \right\| + \|r_{a_n, \eta_{1/k}; N+1}(A, B)\| \\
 (4.5) \quad &\lesssim \max_{\alpha \in \mathbb{N}^{2d}, |\alpha|_\infty \leq N} \|\partial^\alpha (a_n \eta_{1/k})(A, B)\| + \max_{\alpha', \beta' \in \mathbb{N}^d, |\alpha'|, |\beta'| \leq m} \sup_{(x, \xi) \in B(0, 2k)} \langle \xi \rangle^{|\alpha'|} |\partial_\xi^{\alpha'} \partial_x^{\beta'} a_n(x, \xi)|.
 \end{aligned}$$

with constants independent of n . For later reference (we do not need this here) we observe that the constants are also uniform in k , as is evident from the proof of Lemma 4.5.

The first term on the right-hand side of (4.5) can be estimated as follows:

$$\begin{aligned}
 \max_{|\alpha|_\infty \leq N} \|\partial^\alpha (a_n \eta_{1/k})(A, B)\| &\lesssim \max_{|\alpha|_\infty \leq N} \|\partial^\alpha \widehat{(a_n \eta_{1/k})}\|_1 \\
 &\lesssim \max_{|\alpha|_\infty \leq N} \|(u, v) \mapsto \langle (u, v) \rangle^{2d+1} \partial^\alpha \widehat{(a_n \eta_{1/k})}\|_\infty \\
 &\lesssim \max_{|\beta|_\infty \leq N+2d+1} \|\partial^\beta (a_n \eta_{1/k})\|_1 \\
 &\lesssim \max_{|\beta|_\infty \leq N+2d+1} \|\partial^\beta a_n\|_{L^\infty(B(0, 2k))},
 \end{aligned}$$

with constants independent of n . This results in the estimate

$$\begin{aligned}
 &\|a_n(A, B)\eta_{1/k}(A, B)\| \\
 &\lesssim \max_{|\beta|_\infty \leq N+2d+1} \|\partial^\beta a_n\|_{L^\infty(B(0, 2k))} + \max_{|\alpha'|, |\beta'| \leq m} \sup_{(x, \xi) \in B(0, 2k)} \langle \xi \rangle^{|\alpha'|} |\partial_\xi^{\alpha'} \partial_x^{\beta'} a_n(x, \xi)|
 \end{aligned}$$

with constants independent of n .

Set $M := \max(dN + d + 2d^2, 2m)$; the extra factor d in the first term in the maximum comes from $|\alpha| \leq d|\alpha|_\infty$. If all partial derivatives up to order M tend to 0 uniformly on $B(0, 2k)$, it follows that $\lim_{n \rightarrow \infty} \|a_n(A, B)f\| = 0$. ■

DEFINITION 4.7. A function $a \in C^\infty(\mathbb{R}^{2d})$ is said to belong to the standard symbol class S^{-N} , with $N \in \mathbb{Z}$, if

$$\sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{N+|\alpha|} |\partial_\xi^\alpha \partial_x^\beta a(x, \xi)| < \infty$$

for all multi-indices $\alpha, \beta \in \mathbb{N}^d$.

The Schwartz class is included in S^0 , and if $N \geq M$ then $S^{-N} \subseteq S^{-M}$. The class S^{-N} for $N > 0$ plays a key role in estimating error terms that arise from the difference between the pointwise product of functions and their Moyal product. In particular, we use the fact that, for any $N > 0$ and $r \in S^{-N}$, we may write

$$(4.6) \quad T_r f(x) = \int_{\mathbb{R}^d} K_r(x, x - y) f(y) dy,$$

with

$$(4.7) \quad \begin{aligned} & \sup_{x \in \mathbb{R}^d} \int_{\mathbb{R}^d} |K_r(x, x - y)| dy + \sup_{y \in \mathbb{R}^d} \int_{\mathbb{R}^d} |K_r(x, x - y)| dx \\ & \lesssim \max_{|\alpha|, |\beta| \leq 2d+1} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha|} |\partial_\xi^\alpha \partial_x^\beta r(x, \xi)|. \end{aligned}$$

This is proven by combining Proposition 1 p. 554 of [50] and Theorem 5.12 of [1] (see also Theorem 5.15, Corollary 5.16 of [1]).

We are now ready to state and prove the main result of this section. It asserts that the calculus of a Weyl pair with bounded calculus of type $(-N, m)$ extends continuously to symbols in the class S^{-N} .

THEOREM 4.8. *Let $N \in \mathbb{N}$. If (A, B) has a bounded Weyl calculus of type $(-N, m)$, where $m = m(d, N)$ is as in Lemma 4.5, then the Weyl calculus $a \mapsto a(A, B)$ extends continuously to functions $a \in S^{-N}$. More precisely, if $a \in S^{-N}$ is given and $(a_n)_{n \in \mathbb{N}}$ is sequence in $\mathcal{S}(\mathbb{R}^{2d})$ such that for all multi-indices $\gamma \in \mathbb{N}^{2d}$ we have*

$$\partial^\gamma a_n \rightarrow \partial^\gamma a$$

uniformly on compact sets as $n \rightarrow \infty$, then the limit

$$a(A, B) := \lim_{n \rightarrow \infty} a_n(A, B)$$

exists in the strong operator topology of $\mathcal{L}(X)$ and is independent of the approximating sequence. Furthermore, for all $a \in S^{-N}$ we have

$$\|a(A, B)\| \lesssim \max_{|\alpha|, |\beta| \leq m+N} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{N+|\alpha|} |\partial_\xi^\alpha \partial_x^\beta a(x, \xi)|.$$

Proof. The existence and uniqueness of the strong operator limits follows from what we have already proved. As pointed out in Section 1.4, p. 232 of [50], it is possible to approximate functions $a \in S^{-N}$ by Schwartz functions in the way stated, by taking $a_n(x, \xi) = a(x, \xi)\eta(\frac{x}{n}, \frac{\xi}{n}) = a(x, \xi)\eta_{1/n}(x, \xi)$ for some $\eta \in C_c^\infty(\mathbb{R}^{2d})$ such that $\eta(0, 0) = 1$.

It remains to prove the bound for the norm of $a(A, B)$. For this we return to (4.4) and (4.5), both of which also hold if we replace a_n by a . For a given $\varepsilon > 0$, and a large enough k , this gives

$$\begin{aligned} \|a(A, B)\| &\leq \| (a\eta_{1/k})(A, B) \| + 2\varepsilon \\ &\lesssim \max_{\alpha \in \mathbb{N}^{2d}, |\alpha|_\infty \leq N} \|\partial^\alpha (a\eta_{1/k})(A, B)\| \\ &\quad + \max_{\alpha', \beta' \in \mathbb{N}^d, |\alpha'|, |\beta'| \leq m} \sup_{(x, \xi) \in B(0, 2k)} \langle \xi \rangle^{|\alpha'|} |\partial_\xi^{\alpha'} \partial_x^{\beta'} a(x, \xi)| + 2\varepsilon \end{aligned}$$

with estimates uniform in $\varepsilon > 0$ and $k \geq 1$ (note that the sup norms of the derivatives of η_k are uniform in $k \geq 1$).

Each expression in the first term on the right-hand side can be estimated using the type $(-N, m)$ of the Weyl calculus of (A, B) :

$$\begin{aligned} \|\partial^\alpha (a\eta_{1/k})(A, B)\| &\lesssim \max_{\gamma, \delta \in \mathbb{N}^d, |\gamma|, |\delta| \leq m} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{N+|\gamma|} |\partial_\xi^\gamma \partial_x^\delta \partial^\alpha (a\eta_{1/k})(x, \xi)| \\ &\lesssim \max_{\alpha', \beta' \in \mathbb{N}^d, |\alpha'|, |\beta'| \leq m + Nd} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{N+|\alpha'|} |\partial_\xi^{\alpha'} \partial_x^{\beta'} a(x, \xi)|, \end{aligned}$$

again with estimates uniform in $\varepsilon > 0$ and $k \geq 1$. Since $\varepsilon > 0$ was arbitrary, this results in the desired estimate. ■

4.1. BOUNDED WEYL CALCULUS OF TYPE 0 FOR BANACH SPACE-VALUED STANDARD PAIRS. Let X be a UMD space. On $L^p(\mathbb{R}^d; X)$, $1 < p < \infty$, we consider the vector-valued standard pair $(Q \otimes I_X, P \otimes I_X)$ defined by $Q \otimes I_X = (Q_j \otimes I_X)_{j=1}^d$ and $P \otimes I_X = (P_j \otimes I_X)_{j=1}^d$, where Q_j and P_j are the position and momentum operators as in Example 3.3. Note that $(Q \otimes I_X, P \otimes I_X)$ is a Weyl pair: as in the scalar case, $iQ_j \otimes I_X$ and $iP_j \otimes I_X$ generate multiplication and translation groups on $L^p(\mathbb{R}^d; X)$ given by the same formulas as in the scalar-valued case (Example 3.3). The commutation relations for the vector-valued extensions also follow from their scalar-valued counterparts.

NOTATION. In order to simplify notation we will suppress the tensors with I_X when no confusion is likely to arise.

As an illustration of Definition 4.3 we now prove the following theorem.

THEOREM 4.9. *If X is a UMD Banach space, the standard pair (Q, P) has a bounded Weyl calculus of type 0 on $L^p(\mathbb{R}^d; X)$ for all $1 < p < \infty$.*

To prove this theorem we will use Theorem 6 of [47]. To do so, we need to view $a(Q, P)$ as a pseudo-differential operator acting on $L^2(\mathbb{R}^d; X)$. This is possible thanks to the following lemma.

LEMMA 4.10. *For every $a \in \mathcal{S}(\mathbb{R}^{2d})$ there exists a unique $b \in \mathcal{S}(\mathbb{R}^{2d})$ such that $a(Q, P) = T_b$, where T_b is the pseudo-differential operator on $L^2(\mathbb{R}^d)$ defined by*

$$T_b f(x) = \frac{1}{(2\pi)^{d/2}} \int_{\mathbb{R}^d} b(x, \xi) \widehat{f}(\xi) e^{i\xi x} d\xi.$$

This function is given by

$$(4.8) \quad b(x, \xi) = \sum_{|\alpha|=1} \frac{1}{\alpha!} \frac{1}{i^{|\alpha|}} \partial_\xi^\alpha \partial_y^\alpha p_a(x, \xi, y, \xi')|_{y=x, \xi'=\xi} + r_a(x, \xi),$$

where $r_a \in \mathcal{S}(\mathbb{R}^{2d})$ and $p_a(x, \xi, y, \xi') = a(\frac{x+y}{2}, \xi)$. Moreover, for all $m \in \mathbb{N}$, there exists $\tilde{m} \geq m$, depending only on m and d , such that

$$\max_{|\alpha|, |\beta| \leq m} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha|} |\partial_\xi^\alpha \partial_x^\beta r_a(x, \xi)| \lesssim \max_{|\alpha|, |\beta| \leq \tilde{m}} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha|} |\partial_\xi^\alpha \partial_x^\beta a(x, \xi)|.$$

Proof. The first assertion follows from Proposition 1, p. 554 of [50] (see also formula (58), p. 258 of [50]). As in the proof of Lemma 4.5, the estimate follows from Theorem 3.15 of [1]. ■

Proof of Theorem 4.9. We must show that there exists an integer $m \in \mathbb{N}$ such that for all $a \in \mathcal{S}(\mathbb{R}^{2d})$ we have

$$\|a(Q, P)\|_{\mathcal{L}(L^p(\mathbb{R}^d; X))} \lesssim \max_{|\alpha|, |\beta| \leq m} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha|} |\partial_\xi^\alpha \partial_x^\beta a(x, \xi)|.$$

Let $a \in \mathcal{S}(\mathbb{R}^{2d})$. We first apply Lemma 4.10 to write

$$(4.9) \quad a(Q, P) = T_b = \sum_{|\alpha|=1} \frac{1}{\alpha!} \frac{1}{i^{|\alpha|}} T_{\partial_\xi^\alpha \partial_x^\alpha p_a} + T_{r_a},$$

where $\partial_\xi^\alpha \partial_x^\beta p_a(x, \xi)$ is short-hand for the expression $\partial_\xi^\alpha \partial_y^\alpha p_a(x, \xi, y, \xi')|_{y=x, \xi'=\xi}$ occurring in (4.8). We now estimate the $L^p(\mathbb{R}^d; X)$ -norms of the terms on the right-hand side of (4.9) separately, starting with T_{r_a} . As pointed out in (4.6) and (4.7) we have

$$T_{r_a} f(x) = \int_{\mathbb{R}^d} K_{r_a}(x, y) f(y) dy$$

with

$$\sup_{x \in \mathbb{R}^d} \int_{\mathbb{R}^d} |K_{r_a}(x, y)| dy + \sup_{y \in \mathbb{R}^d} \int_{\mathbb{R}^d} |K_{r_a}(x, y)| dx \lesssim \max_{|\alpha|, |\beta| \leq 2d+1} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha|} |\partial_\xi^\alpha \partial_x^\beta r_a(x, \xi)|.$$

Therefore, by Schur’s lemma (in the formulation of Lemma 4.1 with $p = q, r = 1, \phi = \psi \equiv 1$ of [45], noting that the proof extends without change to the vector-valued case), T_{r_a} extends to a bounded operator on $L^p(\mathbb{R}^d; X)$ of norm at most

$$\begin{aligned} \|T_r\|_{\mathcal{L}(L^p(\mathbb{R}^d; X))} &\lesssim \max_{|\alpha|, |\beta| \leq 2d+1} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha|} |\partial_\xi^\alpha \partial_x^\beta r_a(x, \xi)| \\ (4.10) \qquad \qquad \qquad &\lesssim \max_{|\alpha|, |\beta| \leq \tilde{m}} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha|} |\partial_\xi^\alpha \partial_x^\beta a(x, \xi)|, \end{aligned}$$

for some $\tilde{m} \geq 2d + 1$, the second inequality being a consequence of Lemma 4.10.

Next we estimate the $L^p(\mathbb{R}^d; X)$ -norms of the operators $T_{\partial_\xi^\alpha \partial_x^\beta p_a}$. Let $\alpha, \beta \in \mathbb{N}^d$ be such that $|\alpha|, |\beta| \leq 1$. In order to apply Theorem 6 of [47], we remark that $p_{a, \alpha, \beta} := \partial_\xi^\alpha \partial_x^\beta p_a$ has the following (trivial) properties:

(a) for all $|\gamma| \leq 2d + 5$ and $x \in \mathbb{R}^d$ we have

$$\langle \xi \rangle^{|\gamma|} |\partial_\xi^\gamma p_{a, \alpha, \beta}(x, \xi)| = \langle \xi \rangle^{|\gamma|} |\partial_\xi^{\alpha+\gamma} p_a(x, \xi)| \leq \max_{|\alpha'|, |\beta'| \leq 2d+6} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha'|} |\partial_\xi^{\alpha'} \partial_x^{\beta'} a(x, \xi)|;$$

(b) for all $|\gamma|, |\delta| \leq 2d + 5$ we have

$$|\partial_\xi^\gamma \partial_x^\delta p_{a, \alpha, \beta}(x, \xi)| = |\partial_\xi^{\alpha+\gamma} \partial_x^{\beta+\delta} a(x, \xi)| \leq \max_{|\alpha'|, |\beta'| \leq 2d+6} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha'|} |\partial_\xi^{\alpha'} \partial_x^{\beta'} a(x, \xi)|.$$

This means that each $b_{\alpha, \beta}$ belongs to the class $S_{1,0}^0(2d + 5, X)$ as defined in Definition 3 of [47] (note that the R -boundedness condition in this definition reduces to a uniform boundedness condition in view of the fact that we are considering scalar-valued symbols). Therefore, by Theorem 6 of [47] (and its proof, which shows that the estimates depend linearly on the expressions on the right-hand sides in (a) and (b)), the operators $T_{\partial_\xi^\alpha \partial_x^\beta a}$ are bounded on $L^p(\mathbb{R}^d; X)$, and

$$(4.11) \qquad \|T_{\partial_\xi^\alpha \partial_x^\beta a}\|_{\mathcal{L}(L^p(\mathbb{R}^d; X))} \lesssim \max_{|\alpha|, |\beta| \leq 2d+6} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha|} |\partial_\xi^\alpha \partial_x^\beta a(x, \xi)|.$$

Putting together the estimates (4.10) and (4.11) we obtain the following which concludes the proof:

$$\|a(Q, P)\|_{\mathcal{L}(L^p(\mathbb{R}^d; X))} \lesssim \max_{|\alpha|, |\beta| \leq \max(\tilde{m}, 2d+6)} \sup_{(x, \xi) \in \mathbb{R}^{2d}} \langle \xi \rangle^{|\alpha|} |\partial_\xi^\alpha \partial_x^\beta a(x, \xi)|. \quad \blacksquare$$

5. THE OPERATOR $A^2 + B^2$

In this section we show how the Weyl calculus of the pair (A, B) relates to the functional calculus of the operator $A^2 + B^2$.

For Weyl pairs (A, B) of dimension d we define

$$A^2 := \sum_{j=1}^d A_j^2, \quad B^2 := \sum_{j=1}^d B_j^2$$

with domains $D(A^2) := \bigcap_{j=1}^d D(A_j^2)$ and $D(B^2) := \bigcap_{j=1}^d D(B_j^2)$. The operator $A^2 + B^2$ is understood as being defined on $D(A^2) \cap D(B^2)$. Earlier we have already defined $D(A) := \bigcap_{j=1}^d D(A_j)$ and $D(B) := \bigcap_{j=1}^d D(B_j)$.

The following proposition is an immediate consequence of Lemmas 3.11 and 3.12.

PROPOSITION 5.1. *If (A, B) is a Weyl pair of dimension d on X , then $D(A^2) \cap D(B^2)$ is dense in X and invariant under the groups $(e^{itA_j})_{t \in \mathbb{R}}$ and $(e^{itB_j})_{t \in \mathbb{R}}, 1 \leq j \leq d$.*

The next theorem shows, among other things, that for any Weyl pair (A, B) the operator $-(A^2 + B^2)$ is closable and its closure generates an analytic C_0 -semigroup of angle $\frac{1}{2}\pi$. Up to a scaling, this semigroup can be thought of as an abstract version of the Ornstein–Uhlenbeck semigroup. For the standard pair, such a theorem is well-known to mathematical physicists, going back at least to [52]. It was rediscovered for the Ornstein–Uhlenbeck semigroup in Theorem 3.1 of [45]. Here we prove that it holds for all Weyl pairs.

THEOREM 5.2. *Let (A, B) be a Weyl pair. The operators*

$$P(t) := \left(1 + \frac{1 - e^{-t}}{1 + e^{-t}}\right)^d \exp\left(-\frac{1 - e^{-t}}{1 + e^{-t}}(A^2 + B^2)\right) \quad (t \geq 0)$$

define a uniformly bounded C_0 -semigroup on X . The dense set $D(A^2) \cap D(B^2)$ is a core for its generator $-L$, and, on this core, we have the identity

$$L = \frac{1}{2}(A^2 + B^2) - \frac{1}{2}d.$$

The semigroup $(P(t))_{t \geq 0}$ extends to an analytic semigroup of angle $\frac{1}{2}\pi$ that is uniformly bounded and strongly continuous on every subsector of smaller angle.

In the above formula for $P(t)$, for $t > 0$ the right-hand side is interpreted in terms of the Weyl calculus for the pair (A, B) , i.e., $P(t) = a_t(A, B)$, where

$$(5.1) \quad a_t(x, \zeta) := (1 + \lambda_t)^d e^{-\lambda_t(|x|^2 + |\zeta|^2)}$$

with $\lambda_t = \frac{1 - e^{-t}}{1 + e^{-t}}$. For $t = 0$ we interpret the formula as stating that $P(0) = I$.

Proof. The semigroup property $P(t_1)P(t_2) = P(t_1 + t_2)$ follows from the following identity which is obtained by elementary computation:

$$a_{t_1} \# a_{t_2}(x, \zeta)$$

$$\begin{aligned}
 &= \frac{1}{\pi^{2d}}(1+\lambda_{t_1})^d(1+\lambda_{t_2})^d \times \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} e^{-\lambda_{t_1}((x+u)^2+(\xi+u')^2)} e^{-\lambda_{t_2}((x+v)^2+(\xi+v')^2)} \\
 (5.2) \quad &= (1+\lambda_{t_1+t_2})^d e^{-\lambda_{t_1+t_2}(|x|^2+|\xi|^2)} = a_{t_1+t_2}(x, \xi).
 \end{aligned}$$

Next we prove the strong continuity $\lim_{t \downarrow 0} P(t)f = f$ for all $f \in X$. Fix $t > 0$ for the moment. We have

$$\begin{aligned}
 P(t)f &= a_t(A, B)f = \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{a}_t(u, v) e^{i(uA+vB)} f \, du \, dv \\
 &= \frac{1}{(2\pi)^d} (1+\lambda_t)^d \frac{1}{(2\lambda_t)^d} \int_{\mathbb{R}^{2d}} \exp(-\frac{1}{4\lambda_t}(|u|^2+|v|^2)) e^{i(uA+vB)} f \, du \, dv \\
 (5.3) \quad &= \frac{1}{(2\pi)^d} \frac{1}{(1-e^{-t})^d} \int_{\mathbb{R}^{2d}} \exp(-\frac{1+e^{-t}}{4(1-e^{-t})}(|u|^2+|v|^2)) e^{i(uA+vB)} f \, du \, dv
 \end{aligned}$$

so that

$$(5.4) \quad \|P(t)\| \leq \frac{M_A M_B}{(2\pi)^d} \frac{1}{(1-e^{-t})^d} \int_{\mathbb{R}^{2d}} e^{-(1/4(1-e^{-t}))(|u|^2+|v|^2)} \, du \, dv \lesssim M_A M_B.$$

This proves the uniform boundedness of $P(t)$ for $t > 0$. Strong continuity follows from the fact that $\widehat{a}_t \rightarrow \delta_0$ weakly (in the sense that we have strong convergence against every $f \in C_b(\mathbb{R}; X)$).

Let us denote the generator of the C_0 -semigroup $(P(t))_{t \geq 0}$ by $-L$. We claim that $Lf = \frac{1}{2}df - \frac{1}{2}(A^2f + B^2f)$ for all $f \in D(A^2) \cap D(B^2)$. Our argument will be somewhat formal. The reader will have no difficulty in making it rigorous by proceeding as follows: write

$$e^{i(uA+vB)} f = \psi(u, v) e^{i(uA+vB)} f + (1-\psi(u, v)) e^{i(uA+vB)} f$$

for some compactly supported smooth function $a\psi$ which equals 1 in a neighbourhood of $(0, 0)$. Treating the resulting integrals separately, the ones involving ψ will give the desired convergence while the ones involving $1-\psi$ will vanish as we pass to the limit.

Proceeding to the details, we write $P(t) = (1+\lambda)^d R(\lambda)$, where $\lambda = \lambda_t = \frac{1-e^{-t}}{1+e^{-t}}$. Then,

$$\frac{d}{dt} P(t)f = \frac{d}{d\lambda} [(1+\lambda)^d R(\lambda)f] \frac{d\lambda}{dt} = \frac{1}{2}(1-\lambda^2) \frac{d}{d\lambda} [(1+\lambda)^d R(\lambda)f].$$

In the limit $t \downarrow 0$ we also have $\lambda \downarrow 0$ and $\frac{1}{2}(1-\lambda^2) \rightarrow \frac{1}{2}$. Hence the claim will be proved if we show that $\frac{d}{d\lambda} R(\lambda)f \rightarrow df - (A^2+B^2)f$ for $f \in D(A^2) \cap D(B^2)$. We have

$$\lim_{\lambda \downarrow 0} \frac{d}{d\lambda} [(1+\lambda)^d R(\lambda)f] = \lim_{\lambda \downarrow 0} \frac{d}{d\lambda} \left[(1+\lambda)^d \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} e^{-\lambda(|u|^2+|v|^2)} e^{i(uA+vB)} f \, du \, dv \right]$$

$$\begin{aligned}
 &= \lim_{\lambda \downarrow 0} d(1+\lambda)^{d-1} \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} e^{-\lambda(|u|^2+|v|^2)} e^{i(uA+vB)} f \, d u d v \\
 &\quad + \lim_{\lambda \downarrow 0} \left[(1+\lambda)^d \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} \frac{d}{d\lambda} e^{-\lambda(|u|^2+|v|^2)} e^{i(uA+vB)} f \, d u d v \right] \\
 &= \text{(I)} + \text{(II)}.
 \end{aligned}$$

Now, for any $f \in X$,

$$\begin{aligned}
 \text{(I)} &= \lim_{\lambda \downarrow 0} \frac{d}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} e^{-\lambda(|u|^2+|v|^2)} e^{i(uA+vB)} f \, d u d v \\
 &= d \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} \delta_{(0,0)} e^{i(uA+vB)} f \, d u d v = d f.
 \end{aligned}$$

Similarly, for $f \in D(A^2) \cap D(B^2)$,

$$\begin{aligned}
 \text{(II)} &= \lim_{\lambda \downarrow 0} \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} -(|u|^2 + |v|^2) e^{-\lambda(|u|^2+|v|^2)} e^{i(uA+vB)} f \, d u d v \\
 &= \int_{\mathbb{R}^{2d}} \int_{\mathbb{R}^{2d}} \Delta \delta_{(0,0)} e^{i(uA+vB)} f \, d u d v = -(A^2 + B^2) f.
 \end{aligned}$$

Here, $\Delta \delta_{(0,0)}$ denotes the Laplacian of the Dirac delta function in the sense of distributions.

We will prove next that $D(A^2) \cap D(B^2)$ is a core for L . We have already seen that $D(A^2) \cap D(B^2)$ is contained in $D(L)$. The definition of the operators $P(t)$ together with the commutation relation defining Weyl pairs implies that $D(A^2) \cap D(B^2)$ is invariant under $P(t)$. Since $D(A^2) \cap D(B^2)$ is also dense in X , a standard result in semigroup theory implies that $D(A^2) \cap D(B^2)$ is a core for L .

To complete the proof it remains to show the final assertion. By a standard analytic extension argument, the right-hand side of (5.3) defines an analytic extension of $P(t)$ to the open right half-plane which again satisfies the semigroup property. Estimating as in (5.4) we see that this extension is uniformly bounded on every sector of angle strictly less than $\frac{1}{2}\pi$. A standard semigroup argument (see, e.g., Exercise 9.8 of [27]) gives the strong continuity of the extension on each of these sectors. ■

EXAMPLE 5.3. For the standard pair of momentum and position we recover the standard fact that the harmonic oscillator defined by $-Lf(u) = \frac{1}{2}\Delta f(u) - \frac{1}{2}|u|^2 f(u)$ generates a holomorphic semigroup of angle $\frac{1}{2}\pi$, strongly continuous on each smaller sector, on each of the spaces on $L^p(\mathbb{R}^d)$ with $1 \leq p < \infty$.

For later use we make the following simple observation.

COROLLARY 5.4. For all $t > 0$ we have

$$\|tLP(t)\| \leq 2^{d+2}dM_A M_B(1+t)e^{-t}.$$

Proof. Using the same notation as before, write $\lambda'_t = \frac{2e^{-t}}{(1+e^{-t})^2}$ for the derivative of $t \mapsto \lambda_t = \frac{1-e^{-t}}{1+e^{-t}}$. In view of $LP(t)f = -\frac{d}{dt}P(t)f$, differentiation of the right-hand side of (5.3) (and noting that $\frac{1}{1-e^{-t}} = \frac{1}{2}(1 + \lambda_t^{-1})$) gives

$$\begin{aligned} (4\pi)^d LP(t)f &= -\frac{d}{dt} \left((1 + \lambda_t^{-1})^d \int_{\mathbb{R}^{2d}} \exp\left(-\frac{|u|^2+|v|^2}{4\lambda_t}\right) e^{i(uA+vB)} f \, dudv \right) \\ &= d(1 + \lambda_t^{-1})^{d-1} \frac{\lambda'_t}{\lambda_t^2} \int_{\mathbb{R}^{2d}} \exp\left(-\frac{|u|^2+|v|^2}{4\lambda_t}\right) e^{i(uA+vB)} f \, dudv \\ &\quad - (1 + \lambda_t^{-1})^d \frac{\lambda'_t}{4\lambda_t^2} \int_{\mathbb{R}^{2d}} (|u|^2 + |v|^2) \exp\left(-\frac{|u|^2+|v|^2}{4\lambda_t}\right) e^{i(uA+vB)} f \, dudv \end{aligned}$$

and therefore

$$\begin{aligned} \|LP(t)f\| &\leq d(1 + \lambda_t)^{d-1} \frac{\lambda'_t}{\lambda_t} \frac{M_A M_B \|f\|}{(4\pi\lambda_t)^d} \int_{\mathbb{R}^{2d}} \exp\left(-\frac{|u|^2+|v|^2}{4\lambda_t}\right) dudv \\ &\quad + (1 + \lambda_t)^d \frac{\lambda'_t}{4\lambda_t^2} \frac{M_A M_B \|f\|}{(4\pi\lambda_t)^d} \int_{\mathbb{R}^{2d}} (|u|^2 + |v|^2) \exp\left(-\frac{|u|^2+|v|^2}{4\lambda_t}\right) dudv. \end{aligned}$$

In view of the identities

$$\frac{1}{(4\pi\lambda_t)^d} \int_{\mathbb{R}^{2d}} \exp\left(-\frac{|u|^2+|v|^2}{4\lambda_t}\right) dudv = 1$$

and

$$\begin{aligned} &\frac{1}{(4\pi\lambda_t)^d} \int_{\mathbb{R}^{2d}} (|u|^2 + |v|^2) \exp\left(-\frac{|u|^2+|v|^2}{4\lambda_t}\right) dudv \\ &= \frac{1}{(4\pi\lambda_t)^d} \sum_{j=1}^{2d} \int_{\mathbb{R}^{2d}} w_j^2 \exp\left(-\frac{|w|^2}{4\lambda_t}\right) dw \\ &= \frac{1}{(4\pi\lambda_t)^d} \sum_{j=1}^{2d} \left(\int_{\mathbb{R}} w_j^2 \exp\left(-\frac{w_j^2}{4\lambda_t}\right) dw_j \right) \prod_{1 \leq k \leq 2d, k \neq j} \int_{\mathbb{R}} \exp\left(-\frac{w_k^2}{4\lambda_t}\right) dw_k \\ &= \frac{1}{(4\pi\lambda_t)^{1/2}} \sum_{j=1}^{2d} \int_{\mathbb{R}} w_j^2 \exp\left(-\frac{w_j^2}{4\lambda_t}\right) dw_j = 4d\lambda_t \end{aligned}$$

we obtain

$$\begin{aligned} \|tLP(t)f\| &\leq t\left(d2^{d-1}\frac{\lambda'_t}{\lambda_t} + 2^d\frac{\lambda'_t}{4\lambda_t^2} \cdot 4d\lambda_t\right)M_A M_B\|f\| = 2^{d+1}dt\frac{\lambda'_t}{\lambda_t}M_A M_B\|f\| \\ &= 2^{d+1}dt \cdot \frac{2e^{-t}}{(1+e^{-t})^2} \cdot \frac{1+e^{-t}}{1-e^{-t}}M_A M_B\|f\| \\ &= 2^{d+2}d\frac{t}{1-e^{-2t}}e^{-t}M_A M_B\|f\| \leq 2^{d+2}d(1+t)e^{-t}M_A M_B\|f\|. \quad \blacksquare \end{aligned}$$

5.1. GROUND STATES. Let (A, B) be a Weyl pair on the Banach space X . Upon passing to the limit $t_1, t_2 \rightarrow \infty$ in (5.2) one sees that if the limit

$$a_\infty(A, B) := \lim_{t \rightarrow \infty} a_t(A, B)$$

exists in the weak operator topology of $\mathcal{L}(X)$, then it is a projection. That this limit indeed exists under the assumption that X be reflexive is a consequence of the following lemma.

LEMMA 5.5. *Let $(S(t))_{t \geq 0}$ be a C_0 -semigroup on a reflexive Banach space X and let $(T_t)_{t \geq 0}$ be a uniformly bounded family of operators on X such that $S(s) \circ T_t = T_t \circ S(s) = T_{t+s}$ for all $s, t \geq 0$. Then there exists a bounded operator π on X such that $\lim_{t \rightarrow \infty} T_t x = \pi(x)$ weakly for all $x \in X$.*

Proof. Fix $x \in X$. Since X is reflexive, any sequence $t_n \rightarrow \infty$ has a subsequence $t_{n_k} \rightarrow \infty$ such that $\lim_{k \rightarrow \infty} T_{t_{n_k}} x$ exists weakly. Let $\pi(x)$ be this weak limit. We will show that $\pi(x)$ does not depend on the choice of the sequence $t_n \rightarrow \infty$, nor on the choice of the weakly convergent subsequence $t_{n_k} \downarrow 0$. To this end it suffices to show that if both $r_k \rightarrow \infty$ and $s_k \rightarrow \infty$ are such that the weak limits $y := \lim_{k \rightarrow \infty} T_{r_k} x$ and $y' := \lim_{k \rightarrow \infty} T_{s_k} x$ exist, then $y = y'$. By passing to a further subsequence we may assume that $r_k \leq s_k$ for all k . Then $T_{s_k} x = S(s_k - r_k)T_{r_k} x$ and therefore, for all $x^* \in X^*$,

$$\begin{aligned} |\langle T_{s_k} x - T_{r_k} x, x^* \rangle| &= |\langle S(s_k - r_k)T_{r_k} x - T_{r_k} x, x^* \rangle| \leq \|T_{r_k} x\| \|S^*(s_k - r_k)x^* - x^*\|_{X^*} \\ &\leq M\|x\| \|S^*(s_k - r_k)x^* - x^*\|_{X^*}, \end{aligned}$$

where $M = \sup_{t \geq 0} \|T_t\|$. Since X is reflexive, the adjoint semigroup $(S^*(t))_{t \geq 0}$ is strongly continuous (see Proposition I.5.14 of [21]), it follows that

$$|\langle y - y', x^* \rangle| = \lim_{k \rightarrow \infty} |\langle T_{s_k} x - T_{r_k} x, x^* \rangle| \leq M\|x\| \lim_{k \rightarrow \infty} \|S^*(s_k - r_k)x^* - x^*\|_{X^*} = 0.$$

This being true for all $x^* \in X^*$, we conclude that $y = y'$.

The operator π thus defined is linear and bounded, with norm $\|\pi\| \leq M$. That $\pi(x) = \lim_{t \rightarrow \infty} T_t x$ weakly now follows from a standard subsequence argument. \blacksquare

PROPOSITION 5.6. *If (A, B) is a Weyl pair on a reflexive Banach space X , then the weak operator limit $\pi := \lim_{t \rightarrow \infty} a_t(A, B)$ exists in $\mathcal{L}(X)$. Furthermore, $N(L) = R(\pi)$ and $\overline{R(L)} = N(\pi)$.*

Proof. For all $t_1, t_2 \geq 0$, (5.2) implies

$$e^{-t_1 L} \circ a_{t_2}(A, B) = a_{t_2}(A, B) \circ e^{-t_1 L} = a_{t_1+t_2}(A, B).$$

Hence the first assertion follows from the lemma.

The second assertion is proved by a routine semigroup argument. If $f \in N(L)$, then $a_t(A, B)f = e^{-tL}f = f$ implies $\pi(f) = f$, and conversely if $\pi(f) = f$, then for all $t \geq 0$ and $\phi \in X^*$ we have

$$\begin{aligned} \langle e^{-tL}f, \phi \rangle &= \langle f, e^{-tL^*} \phi \rangle = \lim_{s \rightarrow \infty} \langle a_s(A, B)f, e^{-tL^*} \phi \rangle \\ &= \lim_{s \rightarrow \infty} \langle a_{s+t}(A, B)f, \phi \rangle = \langle \pi(f), \phi \rangle = \langle f, \phi \rangle \end{aligned}$$

and therefore $e^{-tL}f = f$. This implies $f \in D(L)$ and $Lf = 0$. This proves $N(L) = R(\pi)$. The proof that $\overline{R(L)} = N(\pi)$ is equally simple. ■

In particular we see that X admits the direct sum decomposition $N(L) \oplus \overline{R(L)}$. Such a decomposition holds for every sectorial operator on a reflexive Banach space (see Proposition 10.1.9 of [33]); the point of the proposition is to identify the associated projection as being given by π .

6. TRANSFERENCE

In the spectral theory of generators of C_0 -groups on a Banach space X , *transference* is a powerful technique to “transfer” properties of the group of translations to general bounded C_0 -groups. The key result is a theorem of Coifman and Weiss [14] which, in the formulation of [30], may be stated as follows.

For functions $g \in L^1(\mathbb{R})$ we consider the convolution operator $K_g f := f * g$. This operator is bounded on $L^p(\mathbb{R}; X)$, $1 \leq p \leq \infty$, by Young’s inequality.

THEOREM 6.1 (Transference, Coifman–Weiss). *Let iA be the generator of a uniformly bounded C_0 -group $(U(t))_{t \in \mathbb{R}}$ on a Banach space X . Then for all $p \in [1, \infty)$, $g \in L^1(\mathbb{R})$, and $x \in X$ one has*

$$\left\| \int_{-\infty}^{\infty} g(t)U(t)x dt \right\| \leq \|K_g\|_{\mathcal{L}(L^p(\mathbb{R}; X))} \sup_{t \in \mathbb{R}} \|U(t)\|^2 \|x\|.$$

Of course one always has the trivial estimate

$$\left\| \int_{-\infty}^{\infty} g(t)U(t)x dt \right\| \leq \|g\|_1 \sup_{t \in \mathbb{R}} \|U(t)\| \|x\|$$

and moreover Young’s inequality gives the estimate

$$\|K_g\|_{\mathcal{L}(L^p(\mathbb{R};X))} \leq \|g\|_1,$$

but there are important examples where the above theorem provides a better estimate, due to the availability of an improved estimate for the norm of K_g in $L^p(\mathbb{R}; X)$. For a systematic treatment of these matters the reader is referred to [33].

The aim of this section is to prove a variant of the Coifman–Weiss theorem for Weyl pairs. Due to the non-commuting nature of the operators A_j and B_j , the convolution product needs to be adjusted to take care of the commutator structure of these operators. This is achieved by means of the so-called twisted convolution.

DEFINITION 6.2. For functions $a \in \mathcal{S}(\mathbb{R}^{2d})$, the twisted convolution with a function $g \in C_c(\mathbb{R}^{2d}; X)$ is defined by

$$(6.1) \quad C_a g(x, \xi) := \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} e^{(1/2)i(x\eta - y\xi)} a(y, \eta) g(x - y, \xi - \eta) dy d\eta.$$

By the pointwise inequality $\|C_a g\| \leq |a| * \|g\|$ and Young’s inequality, C_a extends to a bounded operator on $L^p(\mathbb{R}^{2d}; X)$ for all $1 \leq p \leq \infty$.

PROPOSITION 6.3 (Transference, for Weyl pairs). Let (A, B) be a Weyl pair of dimension d on a Banach space X . Set $M_A := \sup_{t \in \mathbb{R}} \|e^{itA}\|$ and $M_B := \sup_{t \in \mathbb{R}} \|e^{itB}\|$. Let

$1 \leq p < \infty$.

(i) For all $a \in \mathcal{S}(\mathbb{R}^{2d})$ we have

$$\|a(A, B)\| \leq M_A^2 M_B^2 \|C_{\hat{a}}\|_{\mathcal{L}(L^p(\mathbb{R}^{2d}; X))}.$$

(ii) Let $\{a_j : j \in J\}$ be a family of functions in $\mathcal{S}(\mathbb{R}^{2d})$. If the family of twisted convolutions $\{C_{\hat{a}_j}, j \in J\}$ is R -bounded in $\mathcal{L}(L^p(\mathbb{R}^{2d}; X))$, then $\{a_j(A, B) : j \in J\}$ is R -bounded in $\mathcal{L}(X)$, and in that case

$$\mathcal{R}_p(a_j(A, B) : j \in J) \leq M_A^2 M_B^2 \mathcal{R}_p(C_{\hat{a}_j} : j \in J).$$

(iii) Let $\{a_j : j \in J\}$ be a family of functions in $\mathcal{S}(\mathbb{R}^{2d})$. If the family of twisted convolutions $\{C_{\hat{a}_j}, j \in J\}$ satisfies

$$\mathbb{E} \left\| \sum_{j \in J} \epsilon_j C_{\hat{a}_j} g \right\| \lesssim \|g\| \quad \forall g \in L^p(\mathbb{R}^{2d}; X),$$

then

$$\mathbb{E} \left\| \sum_{j \in J} \epsilon_j a_j(A, B) f \right\| \lesssim \|f\| \quad \forall f \in X.$$

Proof. The proof is an adjustment of the technique of [14], [30]. For $r > 0$ we will use the short-hand notation $[-r, r]^2 = \{(x, \xi) \in \mathbb{R}^{2d} : |x| \leq r, |\xi| \leq r\}$. The

elementary estimate $\|a(A, B)\| \leq M_A M_B \|\widehat{a}\|_1$ shows that, for any given $\varepsilon > 0$, we may choose $N > 0$ so large that the operator

$$a(A, B)_{(N)} f := \frac{1}{(2\pi)^d} \int_{\mathbb{C}_{[-N, N]^2}} \widehat{a}(u, v) e^{i(uA+vB)} f \, d u d v$$

satisfies

$$\|a(A, B)_{(N)}\| \leq \frac{M_A M_B}{(2\pi)^d} \int_{\mathbb{C}_{[-N, N]^2}} |\widehat{a}(u, v)| \, d u d v < \varepsilon.$$

We will therefore concentrate on estimating the norm of the operator

$$a(A, B)^{(N)} f := \frac{1}{(2\pi)^d} \int_{[-N, N]^2} \widehat{a}(u, v) e^{i(uA+vB)} f \, d u d v.$$

Accordingly, set $\widehat{a}^{(N)} := \mathbf{1}_{[-N, N]^2} \widehat{a}$. Choose M so large that $\frac{M+N}{M} \leq 1 + \varepsilon$. Let us write $U(u, v) = e^{i(uA+vB)}$ for brevity. By (4.2) we have $U(u, v) \circ U(-u, -v) = I$ and therefore

$$(6.2) \quad \|f\| \leq M_A M_B \|U(-u, -v) f\|, \quad f \in X.$$

Averaging over $[-M, M]^2$, for all $1 \leq p < \infty$ and $f \in X$ we obtain

$$\begin{aligned} & \|a(A, B)^{(N)} f\|^p \\ & \leq \frac{M_A^p M_B^p}{(2M)^{2d}} \int_{[-M, M]^2} \|U(-u, -v) a(A, B)^{(N)} f\|^p \, d u d v \\ & = \frac{M_A^p M_B^p}{(2M)^{2d}} \int_{[-M, M]^2} \left\| \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{a}^{(N)}(y, \eta) U(-u, -v) U(y, \eta) f \, d y d \eta \right\|^p \, d u d v \\ & = \frac{M_A^p M_B^p}{(2M)^{2d}} \int_{[-M, M]^2} \left\| \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} e^{(1/2)i(u\eta - yv)} \widehat{a}^{(N)}(y, \eta) U(y - u, \eta - v) f \, d y d \eta \right\|^p \, d u d v \end{aligned}$$

using (4.2). Also $\mathbf{1}_{[-M-N, M+N]^2}(y - u, \eta - v) = 1$ if $(u, v) \in [-M, M]^2$ and $(y, \eta) \in [-N, N]^2$, so that with $\chi_{M+N} := \mathbf{1}_{[-M-N, M+N]^2}$ the last expression can be rewritten as

$$\begin{aligned} & = \frac{M_A^p M_B^p}{(2M)^{2d}} \int_{[-M, M]^2} \left\| \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} e^{(1/2)i(u\eta - yv)} \widehat{a}^{(N)}(y, \eta) \right. \\ & \quad \times [\chi_{M+N}(y - u, \eta - v) U(y - u, \eta - v) f] \, d y d \eta \left. \right\|^p \, d u d v \\ & = \frac{M_A^p M_B^p}{(2M)^{2d}} \int_{[-M, M]^2} \|C_{\widehat{a}^{(N)}}[\chi_{M+N}(\cdot, \cdot) U(\cdot, \cdot) f](u, v)\|^p \, d u d v \end{aligned}$$

$$\begin{aligned}
 &\leq \frac{M_A^p M_B^p}{(2M)^{2d}} \|C_{\hat{a}^{(N)}} [\chi_{M+N} Uf]\|_{L^p(\mathbb{R}^{2d}; X)}^p \\
 &\stackrel{(i)}{\leq} \frac{M_A^p M_B^p}{(2M)^{2d}} \|C_{\hat{a}^{(N)}}\|_{\mathcal{L}(L^p(\mathbb{R}^{2d}; X))}^p \int_{[-M-N, M+N]^2} \|U(u, v)f\|^p \, dudv \\
 &\leq \frac{M_A^p M_B^p}{(2M)^{2d}} \|C_{\hat{a}^{(N)}}\|_{\mathcal{L}(L^p(\mathbb{R}^{2d}; X))}^p (2(M+N))^{2d} M_A^p M_B^p \|f\|^p \\
 &\leq (1 + \varepsilon)^{2d} M_A^{2p} M_B^{2p} \|C_{\hat{a}^{(N)}}\|_{\mathcal{L}(L^p(\mathbb{R}^{2d}; X))}^p \|f\|^p.
 \end{aligned}$$

It follows that

$$\begin{aligned}
 \|a(A, B)f\| &\leq \|a(A, B)_{(N)}f\| + \|a(A, B)^{(N)}f\| \\
 &\leq \varepsilon \|f\| + (1 + \varepsilon)^{2d/p} M_A^2 M_B^2 \|C_{\hat{a}^{(N)}}\|_{\mathcal{L}(L^p(\mathbb{R}^{2d}; X))} \|f\|.
 \end{aligned}$$

Letting $N \rightarrow \infty$ in this estimate, and then letting $\varepsilon \downarrow 0$, the desired estimate in (i) is obtained.

Part (ii) is proved in exactly the same way. We replace a by $\sum_{n=1}^N \varepsilon_n a_{j_n}$ (where $j_1, \dots, j_N \in J$ and $(\varepsilon_n)_{n=1}^N$ is a Rademacher sequence) and instead of using one fixed f we use a sequence $(f_n)_{n=1}^N$ to build Rademacher sums; instead of estimating with operator norms in (i), we estimate with R -bounds. The same reasoning applies to part (iii). ■

The next lemma expresses the twisted convolution $C_{\hat{a}}$ in terms of the standard pair on $L^2(\mathbb{R}^{2d})$.

LEMMA 6.4. *Let $((Q_1, Q_2), (P_1, P_2))$ be the standard pair of dimension $2d$ on $L^2(\mathbb{R}^{2d})$, i.e.,*

$$\begin{aligned}
 Q_{1,j}f(x, \xi) &= x_j f(x, \xi), & Q_{2,j}f(x, \xi) &= \xi_j f(x, \xi), \\
 P_{1,j}f(x, \xi) &= \frac{1}{i} \frac{\partial f}{\partial x_j}(x, \xi), & P_{2,j}f(x, \xi) &= \frac{1}{i} \frac{\partial f}{\partial \xi_j}(x, \xi),
 \end{aligned}$$

for $1 \leq j \leq d$. The pair $(-\frac{1}{2}Q_2 - P_1, \frac{1}{2}Q_1 - P_2)$ is a Weyl pair of dimension d on $L^2(\mathbb{R}^{2d})$, and for all $a \in \mathcal{S}(\mathbb{R}^{2d})$ we have

$$C_{\hat{a}} = a\left(-\frac{1}{2}Q_2 - P_1, \frac{1}{2}Q_1 - P_2\right).$$

Proof. The proof of the first assertion is immediate. For all $a \in \mathcal{S}(\mathbb{R}^{2d})$ and $g \in L^2(\mathbb{R}^{2d})$,

$$\begin{aligned}
 &a\left(-\frac{1}{2}Q_2 - P_1, \frac{1}{2}Q_1 - P_2\right)g(x, \xi) \\
 &= \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \hat{a}(u, v) e^{i(u(-(1/2)Q_2 - P_1) + v((1/2)Q_1 - P_2))} g(x, \xi) \, dudv
 \end{aligned}$$

$$\begin{aligned}
 &= \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{a}(u, v) e^{(1/2)iu v} e^{iu(-(1/2)Q_2 - P_1)} e^{iv((1/2)Q_1 - P_2)} g(x, \xi) du dv \\
 &= \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{a}(u, v) e^{(1/2)iu v} e^{-(1/2)iu Q_2} e^{-iu P_1} e^{(1/2)iv Q_1} e^{-iv P_2} g(x, \xi) du dv \\
 &= \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{a}(u, v) e^{(1/2)iv Q_1} e^{-(1/2)iu Q_2} e^{-iu P_1} e^{-iv P_2} g(x, \xi) du dv \\
 &= \frac{1}{(2\pi)^d} \int_{\mathbb{R}^{2d}} \widehat{a}(u, v) e^{(1/2)i(vx - \xi u)} g(x - u, \xi - v) du dv = C_{\widehat{a}} g(x, \xi). \quad \blacksquare
 \end{aligned}$$

In the setting of the lemma, by the Stone–von Neumann theorem (see Theorem 14.8 of [29]), there exist a countable index set L and an orthogonal direct sum decomposition

$$L^2(\mathbb{R}^{2d}) = \bigoplus_{\ell \in L} H_\ell,$$

as well as unitary operators $U_\ell : H_\ell \rightarrow L^2(\mathbb{R}^d)$, such that for all $\ell \in L$ the following assertions hold:

- (i) H_ℓ is invariant under the groups $e^{it(-(1/2)Q_{2,j} - P_{1,j})}$ and $e^{it((1/2)Q_{1,j} - P_{2,j})}$;
- (ii) U_ℓ establishes a unitary equivalence of these groups on H_ℓ with the groups e^{itQ_j} and e^{itP_j} on $L^2(\mathbb{R}^d)$, where (Q, P) is the standard pair on $L^2(\mathbb{R}^d)$.

As a direct consequence we obtain that, for all $\ell \in L$ and $a \in \mathcal{S}(\mathbb{R}^{2d})$:

- (i') H_ℓ is invariant under $a(-\frac{1}{2}Q_2 - P_1, \frac{1}{2}Q_1 - P_2)$;
- (ii') U_ℓ establishes a unitary equivalence of the restriction of $a(-\frac{1}{2}Q_2 - P_1, \frac{1}{2}Q_1 - P_2)$ to H_ℓ and the operator $a(Q, P)$ on $L^2(\mathbb{R}^d)$.

As a result we obtain

$$\begin{aligned}
 \|C_{\widehat{a}}\|_{\mathcal{L}(L^2(\mathbb{R}^{2d}))} &= \left\| a\left(-\frac{1}{2}Q_2 - P_1, \frac{1}{2}Q_1 - P_2\right) \right\|_{\mathcal{L}(L^2(\mathbb{R}^{2d}))} \\
 (6.3) \quad &= \sup_{\ell \in L} \left\| a\left(-\frac{1}{2}Q_2 - P_1, \frac{1}{2}Q_1 - P_2\right) |_{H_\ell} \right\|_{\mathcal{L}(H_\ell)} = \|a(Q, P)\|_{\mathcal{L}(L^2(\mathbb{R}^d))}.
 \end{aligned}$$

REMARK 6.5. In the next section we address the problem of estimating the norm of $C_{\widehat{a}}$ in the vector-valued setting. Here we wish to point out the general fact that the identity (6.3) has a simple, albeit not very useful (cf. the concluding remark at the end of the section), vector-valued extension in terms of spaces of γ -radonifying operators. These are defined as follows (comprehensive treatments are given in [33], [44]). Let H be a Hilbert space with inner product $(\cdot | \cdot)$ and X be a Banach space. Every finite rank operator $T : H \rightarrow X$ can be represented as

$$Th = \sum_{n=1}^N (h | h_n) x_n$$

for some orthonormal sequence $(h_n)_{n=1}^N$ in X , and some sequence $(x_n)_{n=1}^N$ in X . For such operators T we define

$$\|T\|_{\gamma(H,X)}^2 := \mathbb{E} \left\| \sum_{n=1}^N \gamma_n x_n \right\|^2,$$

where $(\gamma_n)_{n=1}^N$ is a sequence of independent standard normal random variables (taken real-valued if the scalar field is \mathbb{R} and complex-valued if the scalar field is \mathbb{C} ; once again, one could insist on using real-valued standard normal variables at the expense of different constants). It is easy to see that this gives a well-defined norm on the space of finite rank operators from H to X . Its completion is denoted by $\gamma(H, X)$.

If X is a Hilbert space, then $\gamma(H, X)$ is isometric in a natural way to the space of Hilbert–Schmidt operators from H to X , and if $X = L^p(M, \mu)$ with $1 \leq p < \infty$, then one has a natural isomorphism of Banach spaces

$$\gamma(H, L^p(M, \mu)) \simeq L^p(M, \mu; H).$$

It is not hard to see (see Theorem 9.6.1 of [33]) that if $S : H \rightarrow H$ is a bounded operator, then the mapping $h \otimes x \mapsto Sh \otimes x$ uniquely extends to a bounded operator $\tilde{S} \in \mathcal{L}(\gamma(H^*, X))$ of the same norm. Here, H^* is the Banach space dual of H . Applying this construction to the twisted convolutions C_a and the operators $a(Q, P)$ with $a \in \mathcal{S}(\mathbb{R}^{2d})$, viewed as bounded operators on the Hilbert spaces $L^2(\mathbb{R}^{2d})$ and $L^2(\mathbb{R}^d)$ respectively, and identifying the duals of these spaces with the spaces themselves via the duality $\langle f, g \rangle = \int fg$ (no conjugation here), we obtain well-defined extensions of these operators to bounded operators on $\gamma(L^2(\mathbb{R}^{2d}), X)$ and $\gamma(L^2(\mathbb{R}^d), X)$ of the same norms. Thus (6.3) self-improves to

$$\|C_{\hat{a}}\|_{\mathcal{L}(\gamma(L^2(\mathbb{R}^{2d}), X))} = \|a(Q, P)\|_{\mathcal{L}(\gamma(L^2(\mathbb{R}^d), X))}.$$

This identity suggests that we could try to bound the Weyl calculus in terms of the $\gamma(L^2(\mathbb{R}^{2d}), X)$ -norm of the twisted convolution. This is possible under a γ -boundedness assumption.

PROPOSITION 6.6 (γ -Transference). *Let (A, B) be a Weyl pair of dimension d on a Banach space X . If the set*

$$\{e^{i(uA+vB)} : (u, v) \in \mathbb{R}^{2d}\}$$

is γ -bounded, with γ -bound Γ , then for all $a \in \mathcal{S}(\mathbb{R}^{2d})$ we have

$$\|a(A, B)\| \leq \Gamma^2 \|C_{\hat{a}}\|_{\mathcal{L}(\gamma(L^2(\mathbb{R}^{2d}), X))}.$$

Similar versions of parts (ii) and (iii) of Proposition 6.3 hold. The proof is a routine adaptation of the proof of Proposition 6.3.

As a corollary we obtain that, if the set $\{e^{i(uA+vB)} : (u, v) \in \mathbb{R}^d\}$ is γ -bounded, with γ -bound Γ , then for all $a \in \mathcal{S}(\mathbb{R}^{2d})$ we have

$$\|a(A, B)\| \leq \Gamma^2 \|a(Q, P)\|_{\mathcal{L}(\gamma(L^2(\mathbb{R}^d), X))}.$$

Admittedly, this result is unlikely to be useful: for the standard pair, the γ -boundedness assumption is satisfied only for $p = 2$ (by Proposition 8.1.16 of [33]).

7. R-SECTORIALITY OF L

To apply the transference theory from Section 6, ideally one needs to bound twisted convolution operators $C_{\hat{a}}$ acting on the Bochner spaces $L^p(\mathbb{R}^{2d}; X)$ in terms of the norm of $a(Q, P)$ for the standard pair, i.e., one needs a vector-valued extension of (6.3). We do not know how to do this in general. The L^p -theory in the scalar-valued case, considered by Mauceri in [41], is already quite subtle and depends on Hilbert space-specific techniques to treat the $p = 2$ case. Extending his theory to UMD-valued functions would be interesting in itself (for the new techniques that need to be developed) and would lead to general estimates for the Weyl calculus of Weyl pairs. Here, we just focus on those twisted convolutions needed to study the semigroup generated by $-L = \frac{1}{2}d - \frac{1}{2}(A^2 + B^2)$. The symbols a involved are such that $C_{\hat{a}}$ can be effectively “untwisted”. The main aim of this section is to prove the following result.

THEOREM 7.1 (R-Sectoriality). *Let (A, B) be a Weyl pair on a UMD Banach lattice X . Then for all $\theta \in (0, \pi)$ the operator $L = \frac{1}{2}(A^2 + B^2) - \frac{1}{2}d$ is R -sectorial of angle θ . Moreover, the set $\{(\frac{\pi}{2} - \theta)^{2d} \exp(-zL) ; |\arg(z)| < \theta\}$ is R -bounded, with R -bound independent of θ .*

The only place in the proof where we use the lattice structure of X is in the following lemma, which reduces the task of proving R -boundedness of twisted convolutions to proving R -boundedness of (standard) convolutions.

LEMMA 7.2. *Let X be a Banach lattice with finite cotype and let $1 \leq p < \infty$. Suppose $(a_j)_{j \in J}$ and $(b_j)_{j \in J}$ are families of functions in $\mathcal{S}(\mathbb{R}^{2d})$ satisfying*

$$|a_j(y, \eta)| \leq |b_j(y, \eta)| \quad \forall y, \eta \in \mathbb{R}^d, j \in J.$$

If the family of (standard) convolution operators $(C_{|b_j|})_{j \in J}$ on $L^p(\mathbb{R}^{2d}; X)$ is R -bounded, with R_p -bound \mathcal{R}_p , then also the family $(C_{a_j})_{j \in J}$ on $L^p(\mathbb{R}^{2d}; X)$ is R -bounded, with R_p -bound $\lesssim_{p,q,X} C^p \mathcal{R}_p^p$.

Proof. Since $L^p(\mathbb{R}^{2d}; X)$ has finite cotype (see Proposition 7.1.4 of [33]), we may use the Khintchine–Maurey theorem (see Theorem 7.2.13 of [33]) to pass from Rademacher sums to square functions.

If $j_1, \dots, j_N \in J$ and $g_1, \dots, g_N \in L^p(\mathbb{R}^{2d}; X)$ are given and $(\epsilon_n)_{n=1}^N$ is a Rademacher sequence, we thus obtain

$$\mathbb{E} \left\| \sum_{n=1}^N \epsilon_n C_{a_{j_n}} g_{j_n} \right\|_{L^p(\mathbb{R}^{2d}; X)}^p$$

$$\begin{aligned}
 &\approx \left\| \left(\sum_{n=1}^N |C_{a_{j_n}} g_{j_n}|^2 \right)^{1/2} \right\|_{L^p(\mathbb{R}^{2d}; X)}^p \\
 &= \int_{\mathbb{R}^{2d}} \left\| \left(\sum_{n=1}^N \left| \int_{\mathbb{R}^{2d}} e^{(1/2)i(x\eta - y\xi)} a_{j_n}(y, \eta) g_{j_n}(x - y, \xi - \eta) dy d\eta \right|^2 \right)^{1/2} \right\|^p dx d\xi \\
 &\leq \int_{\mathbb{R}^{2d}} \left\| \left(\sum_{n=1}^N \left(\int_{\mathbb{R}^{2d}} |b_{j_n}(y, \eta)| |g_{j_n}(x - y, \xi - \eta)| dy d\eta \right)^2 \right)^{1/2} \right\|^p dx d\xi \\
 &= \left\| \left(\sum_{n=1}^N (C_{|b_{j_n}|} |g_{j_n}|)^2 \right)^{1/2} \right\|_{L^p(\mathbb{R}^{2d}, X)}^p \approx \mathbb{E} \left\| \sum_{n=1}^N \epsilon_n C_{|b_{j_n}|} |g_{j_n}| \right\|_{L^p(\mathbb{R}^{2d}, X)}^p \\
 &\leq \mathcal{R}_p^p \mathbb{E} \left\| \sum_{n=1}^N \epsilon_n |g_{j_n}| \right\|_{L^p(\mathbb{R}^{2d}, X)}^p \approx \mathcal{R}_p^p \left\| \left(\sum_{n=1}^N |g_{j_n}|^2 \right)^{1/2} \right\|_{L^p(\mathbb{R}^{2d}, X)}^p \\
 &\approx \mathcal{R}_p^p \mathbb{E} \left\| \sum_{n=1}^N \epsilon_n g_{j_n} \right\|_{L^p(\mathbb{R}^{2d}, X)}^p
 \end{aligned}$$

with constants depending only on p and X . ■

We now consider the kernels relevant to our applications. These are the Fourier transforms of

$$a_z(x, \xi) = (1 + \lambda_z)^d e^{-\lambda_z(|x|^2 + |\xi|^2)}$$

with $\lambda_z = \frac{1 - e^{-z}}{1 + e^{-z}}$ for $z \in \mathbb{C}$ such that $\Re z > 0$. We need an elementary lemma.

LEMMA 7.3. *For all $0 < \theta < \frac{1}{2}\pi$ and non-zero $z \in \mathbb{C}$ satisfying $|\arg z| \leq \theta$ we have*

$$\left(\frac{1}{2}\pi - \theta\right) \lesssim \cos(\arg \lambda_z), \quad |\lambda_z| \lesssim \left(\frac{1}{2}\pi - \theta\right)^{-1},$$

with constants independent of θ and z .

Proof. It suffices to prove the inequalities for non-zero $z \in \mathbb{C}$ satisfying $|\arg z| = \theta$. Writing $z = r(\cos \theta + i \sin \theta)$ and computing the real and imaginary parts of $\frac{1 - e^{-z}}{1 + e^{-z}}$ in terms of r and θ , one readily finds that if $|\arg z| < \theta$, then

$$\tan(\arg(\frac{1 - e^{-z}}{1 + e^{-z}})) < \frac{1}{\cos \theta}$$

and consequently

$$\frac{1}{\cos(\arg(\frac{1 - e^{-z}}{1 + e^{-z}}))} < \left(1 + \frac{1}{\cos^2 \theta}\right)^{1/2} \leq 1 + \frac{1}{\cos \theta} \lesssim \frac{1}{(1/2)\pi - \theta}.$$

Similar elementary estimates show that if $|\arg z| < \theta$, then

$$|\lambda_z| = \left| \frac{1 - e^{-z}}{1 + e^{-z}} \right| \lesssim \frac{1}{\cos \theta} \lesssim \frac{1}{(1/2)\pi - \theta}. \quad \blacksquare$$

This lemma is used to prove the following R -boundedness result.

PROPOSITION 7.4. *Let X be a UMD Banach lattice, and let $1 < p < \infty$. There exists a constant $M \geq 0$, depending only on p and X , such that for all $\theta \in (0, \frac{1}{2}\pi)$ the family $\{C_{\hat{a}_z} ; z \neq 0, |\arg z| < \theta\}$ is R -bounded in $\mathcal{L}(L^p(\mathbb{R}^{2d}; X))$, with constant*

$$\mathcal{R}(\{C_{\hat{a}_z} ; z \neq 0, |\arg z| < \theta\}) \leq \frac{M}{((1/2)\pi - \theta)^{2d}}.$$

Proof. The Banach space X , being UMD, has finite cotype (see Proposition 7.3.15 of [33]). Fix $\theta \in (0, \frac{1}{2}\pi)$ and let $z \in \mathbb{C}$ be a non-zero element such that $|\arg z| < \theta$ for all $k = 1, \dots, N$. Writing $t = \frac{1}{\Re(1/\lambda_z)} = |\lambda_z|/\cos(\arg \lambda_z)$, for all $y, \eta \in \mathbb{R}^d$ we have

$$\begin{aligned} |\hat{a}_z(y, \eta)| &\approx |1 + \lambda_z|^d |\lambda_z|^{-d} e^{-\cos(\arg \lambda_z)(|y|^2 + |\eta|^2)/4} |\lambda_z| \\ &\lesssim |1 + \lambda_z|^d (\cos(\arg \lambda_z))^{-d} t^{-d} e^{-(|y|^2 + |\eta|^2)/4t} \end{aligned}$$

with constants depending only on p and X . Hence by Lemma 7.3,

$$\left(\frac{1}{2}\pi - \theta\right)^{2d} |\hat{a}_z(y, \eta)| \lesssim t^{-d} e^{-(|y|^2 + |\eta|^2)/4t} =: b_t(y, \eta),$$

with constants depending only on p and X . Hence, by Lemma 7.2,

$$\mathcal{R}\left(\left\{\left(\frac{1}{2}\pi - \theta\right)^{2d} C_{\hat{a}_z} ; z \neq 0, |\arg z| < \theta\right\}\right) \lesssim \mathcal{R}(\{C_{b_t} ; t > 0\}),$$

with a constant depending only on p and X . Noting that C_{b_t} is a constant multiple of $\exp(t\Delta \otimes I_X)$, the R -boundedness of the family $\{C_{b_t} ; t > 0\}$ follows from the fact that $-\Delta \otimes I_X$ has a bounded H^∞ -calculus on $L^p(\mathbb{R}^{2d}; X)$ ([33], Theorem 10.2.25). ■

Proof of Theorem 7.1. Let $(P(t))_{t \geq 0}$ be the analytic C_0 -semigroup generated by $-L = \frac{1}{2}d - \frac{1}{2}(A^2 + B^2)$. Fix $\theta \in (0, \frac{1}{2}\pi)$. By Proposition 10.3.3 of [33] it suffices to show that the set

$$V_\theta := \{P(z) : z \neq 0, |\arg z| < \theta\}$$

is R -bounded with

$$\mathcal{R}(V_\theta) \lesssim_{p,X} \frac{1}{((1/2)\pi - \theta)^{2d}}.$$

By Theorem 5.2 and Proposition 6.3, for this it suffices to show that the set

$$V'_\theta := \{C_{\hat{a}_z} : z \neq 0, |\arg z| < \theta\}$$

is R -bounded with

$$\mathcal{R}(V'_\theta) \lesssim_{p,X} \frac{1}{((1/2)\pi - \theta)^{2d}}.$$

This has been done in Proposition 7.4. ■

8. FUNCTIONAL CALCULUS OF L

In this section we prove that boundedness of the Weyl calculus of a Weyl pair (A, B) implies a spectral multiplier theorem for the operator $L = \frac{1}{2}(A^2 + B^2) - \frac{1}{2}d$, acting on a UMD lattice X . This is done by applying the theory developed in [35] to obtain a holomorphic functional calculus of angle zero from square function estimates and appropriate R -sectoriality bounds. The precise form of the latter then allows us to apply the theory developed in [38] to extend this holomorphic functional calculus to a full Hörmander type spectral multiplier theorem.

THEOREM 8.1 (Bounded H^∞ -calculus). *If (A, B) is a Weyl pair of dimension d on a UMD lattice X with a bounded Weyl calculus of type 0, then $L := \frac{1}{2}(A^2 + B^2) - \frac{1}{2}d$ has a bounded $H^\infty(\Sigma_\theta)$ -calculus for all $\theta \in (0, \pi)$.*

For the proof of the theorem we need two lemmas. The first provides an expression for derivatives of the exponentials in the Weyl calculus representation formula of the operators e^{-tL} of Theorem 5.2.

LEMMA 8.2. *For all multi-indices $\alpha, \beta \in \mathbb{N}^d$ there is a polynomial $p_{\alpha, \beta}$ of degree (α, β) in the variables $(x, \xi) \in \mathbb{R}^{2d}$ such that for all $\lambda > 0$ we have*

$$\partial_\xi^\alpha \partial_x^\beta e^{-\lambda(|x|^2 + |\xi|^2)} = \sqrt{\lambda}^{|\alpha| + |\beta|} p_{\alpha, \beta}(\sqrt{\lambda}x, \sqrt{\lambda}\xi) e^{-\lambda(|x|^2 + |\xi|^2)}.$$

Proof. If p is a polynomial in $2d$ variables $x = (x_1, \dots, x_d)$, $\xi = (\xi_1, \dots, \xi_d)$, of degree $\gamma = (\gamma_1, \dots, \gamma_d)$ in x and $\delta = (\delta_1, \dots, \delta_d)$ in ξ , then for any $\lambda > 0$,

$$\begin{aligned} & \partial_{x_j} [p(\sqrt{\lambda}x, \sqrt{\lambda}\xi) e^{-\lambda(|x|^2 + |\xi|^2)}] \\ &= [\sqrt{\lambda}(\partial_{x_j} p)(\sqrt{\lambda}x, \sqrt{\lambda}\xi) - 2\lambda x_j p(\sqrt{\lambda}x, \sqrt{\lambda}\xi)] e^{-\lambda(|x|^2 + |\xi|^2)} \\ &= \sqrt{\lambda} q(\sqrt{\lambda}x, \sqrt{\lambda}\xi) e^{-\lambda(|x|^2 + |\xi|^2)}, \end{aligned}$$

where q is polynomial of degree $(\gamma_1, \dots, \gamma_{j-1}, \gamma_j + 1, \gamma_{j+1}, \dots, \gamma_d)$ in x and of degree δ in ξ . A similar identity holds for the partial derivatives with respect to ξ_j , which add one to the degree in the variable ξ_j . The lemma now follows by induction on α and β . ■

As an application of the preceding lemma, the next lemma provides a uniform bound on the derivatives of certain signed sums of exponentials which will be used later to prove that certain related sums belong to the symbol class S^0 uniformly.

As before, let $\lambda_t = \frac{1 - e^{-t}}{1 + e^{-t}}$ for $t > 0$.

LEMMA 8.3. *For all multi-indices $\alpha, \beta \in \mathbb{N}^d$ such that $|\alpha| + |\beta| \neq 0$ the functions*

$$\kappa_{k, \epsilon, s}(x, \xi) := \sum_{j=1}^k \epsilon_j \exp(-\lambda_{2^{-j}s}(|x|^2 + |\xi|^2))$$

satisfy

$$\sup_{(x,\zeta) \in \mathbb{R}^{2d}} \langle \zeta \rangle^{|\alpha|} |\partial_{\zeta}^{\alpha} \partial_x^{\beta} \kappa_{k,\epsilon,s}(x, \zeta)| < \infty$$

uniformly with respect to $k \geq 1$, $\epsilon = (\epsilon_j)_{j=1}^k \in \{\pm 1\}^k$, and $s \in [1, 2]$.

Proof. Let us set $\mu_{j,s} = \lambda_{2^{-j}s}$ for brevity. Given any two multi-indices $\alpha, \beta \in \mathbb{N}^d$ such that $|\alpha| + |\beta| \neq 0$ we may estimate, using Lemma 8.2,

$$\begin{aligned} & \langle \zeta \rangle^{|\alpha|} |\partial_{\zeta}^{\alpha} \partial_x^{\beta} \kappa_{k,\epsilon,s}(x, \zeta)| \\ & \leq \langle \zeta \rangle^{|\alpha|} \sum_{j=1}^k \sqrt{\mu_{j,s}}^{|\alpha|+|\beta|} |p_{\alpha,\beta}(\sqrt{\mu_{j,s}}x, \sqrt{\mu_{j,s}}\zeta)| \exp(-\mu_{j,s}(|x|^2 + |\zeta|^2)). \end{aligned}$$

For $|\zeta| \leq 1$ we estimate the right-hand side by

$$\begin{aligned} & \lesssim_{\alpha} \sum_{j=1}^k \sqrt{\mu_{j,s}}^{|\alpha|+|\beta|} |p_{\alpha,\beta}(\sqrt{\mu_{j,s}}x, \sqrt{\mu_{j,s}}\zeta)| \exp(-\mu_{j,s}(|x|^2 + |\zeta|^2)) \\ & \lesssim_{\alpha,\beta} \sum_{j=1}^k \sqrt{\mu_{j,s}}^{|\alpha|+|\beta|} \lesssim_{\alpha,\beta} \sum_{j=1}^k 2^{-(1/2)j(|\alpha|+|\beta|)} \lesssim_{\alpha,\beta} 1 \end{aligned}$$

where we used that $\sup_{(x',\zeta') \in \mathbb{R}^{2d}} |p_{\alpha,\beta}(x', \zeta')| \exp(-(|x'|^2 + |\zeta'|^2)) < \infty$, and that

$\mu_{j,s} \lesssim 2^{-j}$; while for $|\zeta| > 1$ we may estimate it by

$$\lesssim_{\alpha} \sum_{j=1}^k (\sqrt{\mu_{j,s}}|\zeta|)^{|\alpha|+|\beta|} |p_{\alpha,\beta}(\sqrt{\mu_{j,s}}x, \sqrt{\mu_{j,s}}\zeta)| \exp(-\mu_{j,s}(|x|^2 + |\zeta|^2)) \lesssim_{\alpha,\beta} 1,$$

where the last step follows by an application of Proposition H.2.3 in [33]. In all these estimates, the constants are uniform in k, ϵ , and s . ■

Proof of Theorem 8.1. By Theorem 7.1 L is R -sectorial of angle θ for any $\theta \in (0, \pi)$. Hence by Theorem 10.4.9 of [33] it suffices to show that

$$\|f\|^2 \sim \sup_{s \in [1,2]} \sup_{N \in \mathbb{N}} \mathbb{E} \left\| \sum_{|j| \leq N} \epsilon_j (\exp(-2^{j+1}sL) - \exp(-2^j sL)) f \right\|^2 \quad \forall f \in D(L) \cap R(L),$$

where $(\epsilon_j)_{j \in \mathbb{Z}}$ is a Rademacher sequence, noting that the function $z \mapsto \exp(-2z) - \exp(-z)$ belongs to $H^1(\Sigma_{\theta}) \cap H^{\infty}(\Sigma_{\theta})$ for each $\theta \in (0, \frac{1}{2}\pi)$ using the notation of Chapter 10 in [33]. It actually suffices to prove the one-sided inequality

$$(8.1) \quad \sup_{s \in [1,2]} \sup_{N \in \mathbb{N}} \mathbb{E} \left\| \sum_{|j| \leq N} \epsilon_j (\exp(-2^{j+1}sL) - \exp(-2^j sL)) f \right\|^2 \lesssim \|f\|^2 \quad \forall f \in X,$$

since the reverse inequality (for $f \in \overline{R(L)} = \overline{D(L) \cap R(L)}$) will then follow by duality as in the proof of Theorem 10.4.4(3) in [33], noting that the pair of adjoint operators (B^*, A^*) is a Weyl pair in the dual lattice X^* (which is UMD by Proposition 4.2.17 of [32] and Proposition 7.5.15 of [33]).

Referring to the direct sum decomposition $X = N(L) \oplus \overline{R(L)}$ provided by Proposition 5.6, we will prove (8.1) separately for $f \in N(L)$ and $f \in \overline{R(L)}$. For $f \in N(L)$, (8.1) is immediate from the fact that $\exp(-tL)f = f$. For $f \in \overline{R(L)}$ we consider indices $j \in \mathbb{N}$ and $j \in \mathbb{Z} \setminus \mathbb{N}$ separately. For $j \in \mathbb{N}$ and $f \in R(L)$, say $f = Lg$, we use that $\exp(-tL)f = L\exp(-tL)g$ decays to 0 exponentially as $t \rightarrow \infty$ by Corollary 5.4. In combination with the triangle inequality and the contraction principle, this gives

$$\begin{aligned} & \sup_{s \in [1,2]} \sup_{N \in \mathbb{N}} \left(\mathbb{E} \left\| \sum_{j=0}^N \epsilon_j (\exp(-2^{j+1}sL) - \exp(-2^j sL)) f \right\|^2 \right)^{1/2} \\ & \leq 2 \sup_{s \in [1,2]} \sup_{N \in \mathbb{N}} \left(\mathbb{E} \left\| \sum_{j=0}^N \epsilon_j \exp(-2^j sL) f \right\|^2 \right)^{1/2} \\ & \leq 2 \sup_{s \in [1,2]} \sup_{N \in \mathbb{N}} \sum_{j=0}^N \|\exp(-2^j sL) f\| \lesssim \|f\|. \end{aligned}$$

By continuity, this estimate extends to $f \in \overline{R(L)}$.

Let

$$\tilde{a}_t(x, \xi) := a_{2t}(x, \xi) - a_t(x, \xi),$$

where as always $a_t(x, \xi) = (1 + \lambda_t)^d e^{-\lambda_t(|x|^2 + |\xi|^2)}$ with $\lambda_t := \frac{1 - e^{-t}}{1 + e^{-t}}$. In view of the cases already dealt with, the proof will be complete once we have shown that, for all $f \in \overline{R(L)}$,

$$(8.2) \quad \sup_{s \in [1,2]} \sup_{N \geq 1} \mathbb{E} \left\| \sum_{j=1}^N \epsilon_j \tilde{a}_{2^{-j}s}(A, B) f \right\|^2 \lesssim \|f\|^2$$

with a constant independent of f . Set

$$\tilde{b}_t := (1 + \lambda_{2t})^{-d} a_{2t} - (1 + \lambda_t)^{-d} a_t$$

so that

$$\tilde{a}_t = \tilde{b}_t - ((1 + \lambda_{2t})^{-d} - 1) a_{2t} + ((1 + \lambda_t)^{-d} a_t - 1) a_t.$$

We now take $t = 2^{-j}s$ and estimate each of the resulting three sums separately. Fix an integer $N \geq 1$. We first estimate

$$\begin{aligned} & \sup_{s \in [1,2]} \sup_{N \geq 1} \mathbb{E} \left\| \sum_{j=1}^N \epsilon_j ((1 + \lambda_{2^{-j+1}s})^{-d} - 1) a_{2^{-j+1}s}(A, B) f \right\|^2 \\ & \lesssim \sup_{s \in [1,2]} \left(\sum_{j=1}^{\infty} |(1 + \lambda_{2^{-j+1}s})^{-d} - 1| \|\exp(-2^{-j+1}sL) f\| \right)^2 \end{aligned}$$

$$\lesssim_d \left(\sum_{j=0}^{\infty} 2^{-j} \|f\| \right)^2 \lesssim \|f\|^2,$$

where we used the bound $(1 + \lambda_{2^{-j}})^{-d} - 1 \lesssim_d 2^{-j}$ together with the uniform boundedness of the operators $\exp(-tL) = P(t)$. Similarly,

$$\sup_{s \in [1,2]} \sup_{N \geq 1} \mathbb{E} \left\| \sum_{j=1}^N \epsilon_j ((1 + \lambda_{2^{-j}})^{-d} - 1) a_{2^{-j}s}(A, B) f \right\|^2 \lesssim \|f\|^2.$$

To prove (8.2), it therefore remains to show that

$$\sup_{s \in [1,2]} \sup_{N \geq 1} \mathbb{E} \left\| \sum_{j=1}^N \epsilon_j \tilde{b}_{2^{-j}s}(A, B) f \right\|^2 \lesssim \|f\|^2.$$

To this end we claim that the functions

$$\tilde{\kappa}_{N,\epsilon,s} := \sum_{j=1}^N \epsilon_j \tilde{b}_{2^{-j}s}$$

belong to the symbol class S^0 , uniformly in $N, \epsilon \in \{\pm 1\}^N$, and $s \in [1, 2]$. Since by assumption (A, B) has a bounded Weyl calculus of type 0, this claim, once it has been proved, will prove the theorem.

We have

$$\begin{aligned} |\tilde{\kappa}_{N,\epsilon,s}(x, \zeta)| &\leq \sum_{j=1}^N |\tilde{b}_{2^{-j}s}(x, \zeta)| \\ &= \sum_{j=1}^N |\exp(-\lambda_{2^{-j+1}s}(|x|^2 + |\zeta|^2)) - \exp(-\lambda_{2^{-j}s}(|x|^2 + |\zeta|^2))| \\ &= \sum_{j=1}^N (\exp(-\lambda_{2^{-j+1}s}(|x|^2 + |\zeta|^2)) - \exp(-\lambda_{2^{-j}s}(|x|^2 + |\zeta|^2))) \leq 1 \end{aligned}$$

using a telescoping argument in the last step. In combination with Lemma 8.3 (which remains true if we replace the summation $\sum_{j=1}^N$ by $\sum_{j=0}^{N-1}$), this proves the claim. ■

REMARK 8.4. For the standard pair (Q, P) on $L^p(\mathbb{R}^d; X)$ with $1 < p < \infty$ and X any UMD space, the operator $\frac{1}{2}(Q^2 + P^2) - \frac{1}{2}d$ is R -sectorial and has a bounded H^∞ -calculus for any angle $\theta \in (0, \pi)$. Following the lines of [40], this follows from the results of [9] (see also [2]).

Using the theory developed in [38], we can extend the functional calculus of L from $H^\infty(\Sigma_\theta)$ to an appropriate Hörmander class. We thus obtain a calculus in one of their \mathcal{H}_2^β classes. As pointed out in Remark 3.3 of [38], these classes are

slightly larger (but more complicated to define) than the standard Hörmander classes of functions $f \in C^m[0, \infty)$ satisfying

$$\sup_{R>0} R^{2k} \int_{(1/2)R}^{2R} |f^{(k)}(t)|^2 \frac{dt}{R} < \infty,$$

for all $k = 0, \dots, m$. Note that the latter class contains all smooth functions with compact support in $(0, \infty)$.

THEOREM 8.5. *If (A, B) is a Weyl pair on a UMD lattice X with a bounded Weyl calculus of type 0, then $L = \frac{1}{2}(A^2 + B^2) - \frac{1}{2}d$ has an R -bounded $\mathcal{H}_2^{2d+(1/2)}$ -Hörmander calculus.*

Proof. By Theorems 7.1, 8.1, the assumptions of Theorem 7.1 in [38] are satisfied (note that UMD lattices have the required property (α) by Theorem 7.5.20; see the Notes to this section for the terminology of [33]). ■

9. OPEN PROBLEMS

As explained in the introduction, this paper is mostly meant as a foundation for the development of pseudo-differential calculi in “rough” settings. We nonetheless think that the general theory of Weyl pairs presented here is also worth developing further in its own right. This would include solving the following problems:

(1) To extend Mauceri’s results on twisted convolutions [41] to $L^p(\mathbb{R}^{2d}; X)$, where X is UMD and $p \in (1, \infty)$.

An affirmative answer would automatically solve the next problem.

(2) To extend Theorem 4.9 to general Weyl pairs (A, B) .

As observed in Remark 8.4, for standard pairs on $L^p(\mathbb{R}^d; X)$ with $1 < p < \infty$, the conclusions of Theorems 7.1, 8.1, and 8.5 hold for arbitrary UMD Banach spaces X . In the three theorems for general Weyl pairs, the lattice structure of X was only used through the proof of Lemma 7.2. Thus one may pose the following problem:

(3) To decide whether Theorems 7.1, 8.1, and 8.5 hold for arbitrary UMD spaces X .

An affirmative answer to the first problem could possibly also solve the third problem, since it might pave the way for an alternative proof via transference. For these three problems, studying the particular case where X is a non-commutative L^p -space would be particularly interesting, yet potentially much simpler than the general case (thanks to the availability of both domination and extrapolation techniques).

(4) To prove an analogue of the Stone–von Neumann uniqueness theorem for Weyl pairs.

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